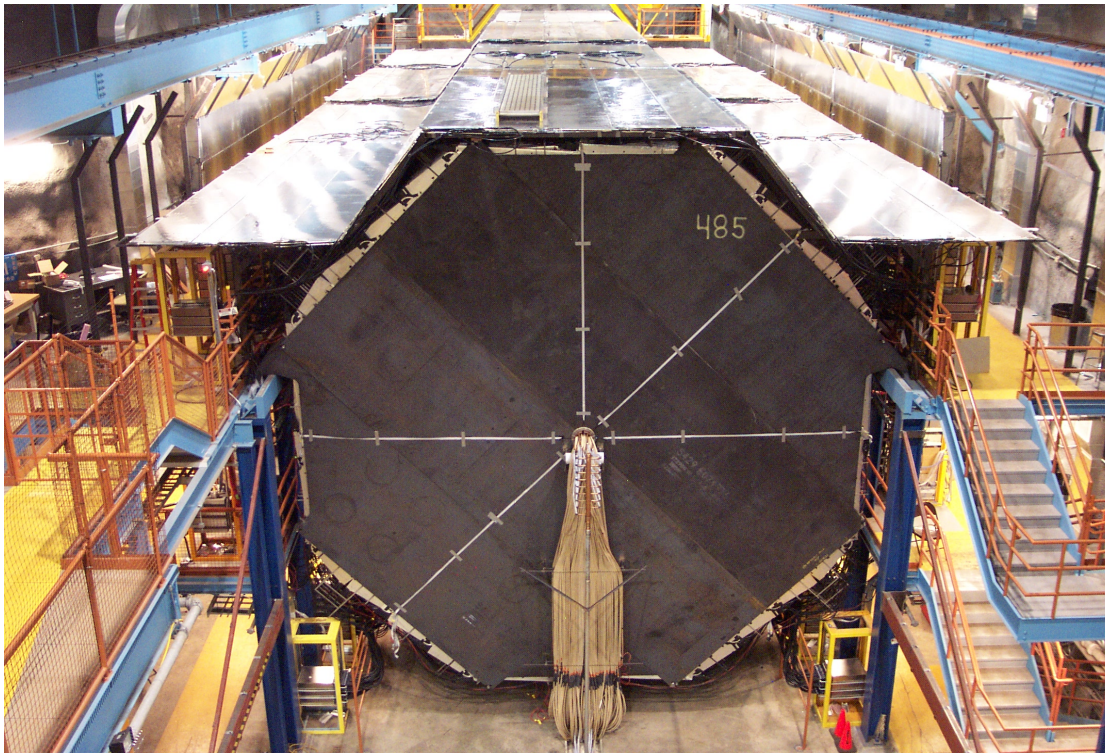




Recent Results from MINOS



Mark Thomson
University of Cambridge



This talk:

- ① **Neutrinos: 1998 – 2008**
- ② **The MINOS Experiment**
- ③ **The NuMI Beam**
- ④ **The MINOS Detectors**
- ⑤ **ν_{μ} Disappearance Results**
- ⑥ **Other Results**
- ⑦ **Future Prospects**
- ⑧ **Summary**



1 Neutrinos: 1998 – 2008



10 years ago (PDG1998):

- ★ Standard Model : assumed massless ν
- ★ Fundamental states : ν_e, ν_μ, ν_τ
- ★ $m(\nu_e) < 3 \text{ eV}, \dots$

+ hints of Neutrino Oscillations:

- ★ Atmospheric neutrino oscillations
 - Statistically marginal + positive & negative results
- ★ Solar neutrino oscillations
 - Required faith in Astrophysics/Astrophysicists....!

January 2008:

- ★ Standard Model : massive ν
- ★ Fundamental (mass) eigenstates : ν_1, ν_2, ν_3
- ★ Atmospheric neutrino oscillations
 - Compelling evidence : Super-Kamiokande
- ★ Solar neutrino oscillations
 - Compelling evidence : SNO + Super-Kamiokande
- ★ Reactor neutrino oscillations
 - Compelling evidence : KamLAND
- ★ Beam neutrino oscillations
 - **Moving into the era of precision measurements: MINOS**



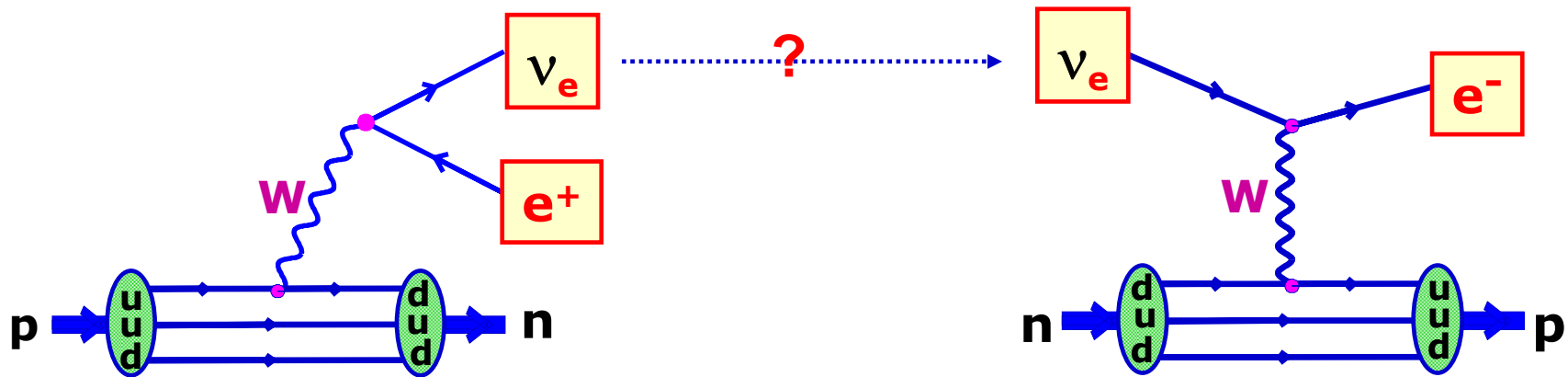
Neutrino Flavours Revisited



- ★ Never **directly** observe neutrinos – can only detect them by their weak interactions hence by **definition** ν_e is the neutrino state produced along with an electron. Similarly, charged current weak interactions of the state ν_e produce an electron

$$\nu_e, \nu_\mu, \nu_\tau = \text{weak eigenstates}$$

- For many years, assumed that ν_e, ν_μ, ν_τ were massless fundamental particles
- **Experimental evidence:** **at short distances** neutrinos produced along with an electron always produced an electron in CC Weak interactions, etc.



- ★ Now know that the weak eigenstates, ν_e, ν_μ, ν_τ , are linear combinations of the mass eigenstates

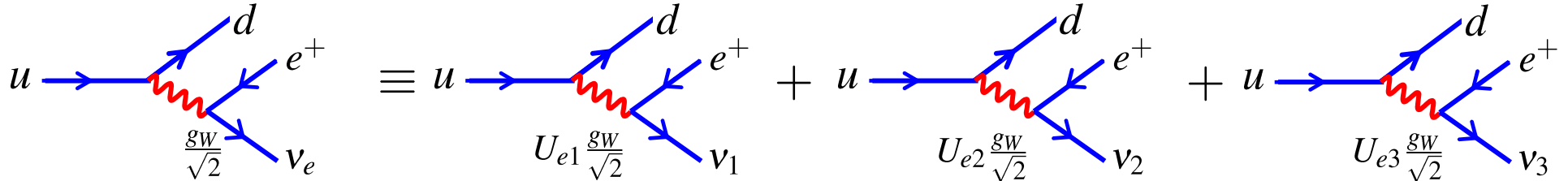


Neutrino Oscillations for Three Flavours



- ★ Relate the weak eigenstates to the mass eigenstates via the Unitary **PMNS** Pontecorvo-Maki-Nakagawa-Sakata matrix.

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu1} & U_{\mu2} & U_{\mu3} \\ U_{\tau1} & U_{\tau2} & U_{\tau3} \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$



- ★ To calculate the oscillation probability, consider a state which is produced at $t = 0$ as a $|\nu_e\rangle$

$$|\psi(t = 0)\rangle = |\nu_e\rangle = U_{e1}|\nu_1\rangle + U_{e2}|\nu_2\rangle + U_{e3}|\nu_3\rangle$$

i.e. a coherent linear combination of the mass eigenstates

- From which we can calculate the **oscillation probability**

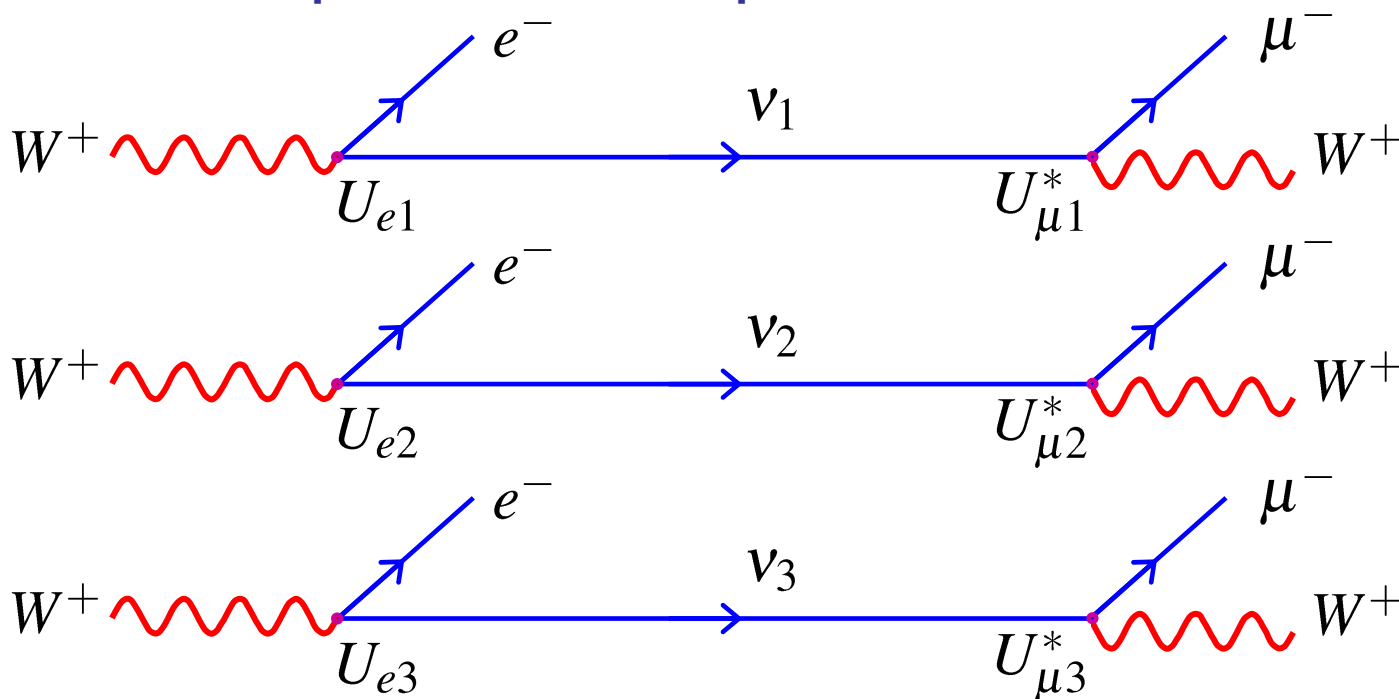
Phases from time evolution of mass eigenstates

$$\begin{aligned} P(\nu_e \rightarrow \nu_\mu) &= |\langle \nu_\mu | \psi(L) \rangle|^2 \\ &= |U_{e1}U_{\mu1}^*e^{-i\phi_1} + U_{e2}U_{\mu2}^*e^{-i\phi_2} + U_{e3}U_{\mu3}^*e^{-i\phi_3}|^2 \end{aligned}$$



$$P(\nu_e \rightarrow \nu_\mu) = |U_{e1}U_{\mu 1}^*e^{-i\phi_1} + U_{e2}U_{\mu 2}^*e^{-i\phi_2} + U_{e3}U_{\mu 3}^*e^{-i\phi_3}|^2$$

- The terms in this expression can be represented as:



- Because of the unitarity of the PMNS matrix:

$$U_{e1}U_{\mu 1}^* + U_{e2}U_{\mu 2}^* + U_{e3}U_{\mu 3}^* = 0$$

Consequently, unless the phases of the different components are different, the sum of these three diagrams is zero, i.e., require **different neutrino masses** for osc.



PMNS Matrix



- ★ The PMNS matrix is usually expressed in terms of 3 rotation angles $\theta_{12}, \theta_{23}, \theta_{13}$ and a complex phase δ , using the notation $s_{ij} = \sin \theta_{ij}$, $c_{ij} = \cos \theta_{ij}$

$$\begin{pmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu 1} & U_{\mu 2} & U_{\mu 3} \\ U_{\tau 1} & U_{\tau 2} & U_{\tau 3} \end{pmatrix} = \underbrace{\begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix}}_{\text{“Atmospheric”}} \times \begin{pmatrix} c_{13} & 0 & s_{13}e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{-i\delta} & 0 & c_{13} \end{pmatrix} \times \underbrace{\begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix}}_{\text{“Solar”}}$$

Dominates:

“Atmospheric”

“Solar”

- Writing this out in full:

$$U = \begin{pmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta} \\ -s_{12}c_{23} - c_{12}s_{23}e^{i\delta} & c_{12}c_{23} - s_{12}s_{23}s_{13}e^{i\delta} & s_{23}c_{13} \\ s_{12}s_{23} - c_{12}c_{23}s_{13}e^{i\delta} & -c_{12}s_{23} - s_{12}c_{23}s_{13}e^{i\delta} & c_{23}c_{13} \end{pmatrix}$$

- ★ There are **six SM parameters** that can be measured in ν oscillation experiments

$ \Delta m_{21} ^2 = m_2^2 - m_1^2 $	θ_{12}	Solar and reactor neutrino experiments
$ \Delta m_{32} ^2 = m_3^2 - m_2^2 $	θ_{23}	Atmospheric and beam neutrino experiments
	θ_{13}	Reactor neutrino experiments + future beam
	δ	Future beam experiments (CP violation)



Neutrinos Today: What do we/don't we know ?

“KNOWN”

- ★ Solar neutrino oscillations (mainly Super-K and SNO and KamLAND):

$$|\Delta m_{21}^2| = 7.9_{-0.5}^{+0.6} \times 10^{-5} \text{ eV}^2$$

$$\tan^2 \theta_{12} = 0.40_{-0.07}^{+0.1}$$

- ★ Atmospheric neutrino oscillations (mainly Super-K): + MINOS

$$|\Delta m_{32}^2| \approx (2.5 \pm 0.5) \times 10^{-3} \text{ eV}^2$$

$$\sin^2 \theta_{23} > 0.92$$

Near maximal mixing

- ★ Reactor neutrino non-oscillations (CHOOZ):

$$\sin^2 \theta_{13} < 0.06$$

Small

UNKNOWN:

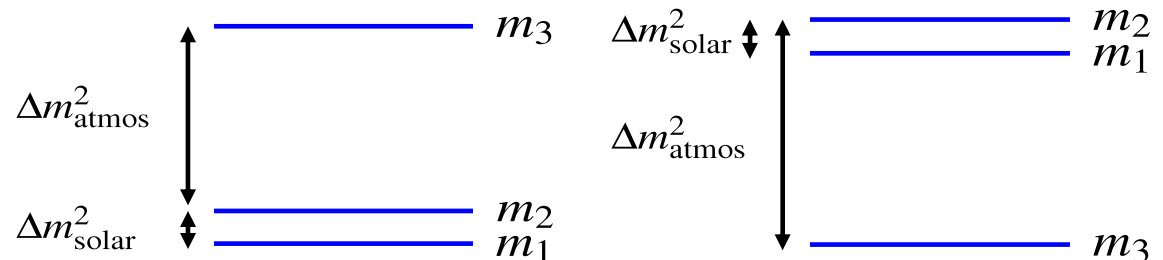
Mass

$$m_{\nu_e} < 2 \text{ eV}$$

CP-Phase

δ

Mass Hierarchy

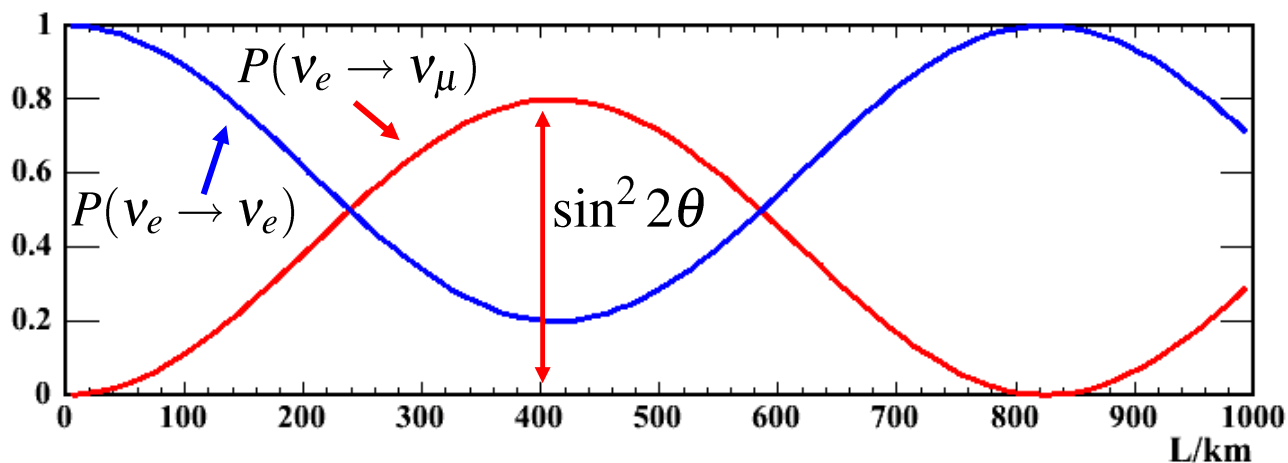




★ Because θ_{13} is small, in many circumstances the two oscillation scales, $|\Delta m_{12}^2|$, $|\Delta m_{32}^2|$ decouple and can use the two flavour oscillation formula

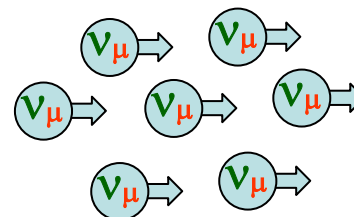
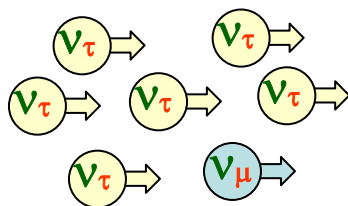
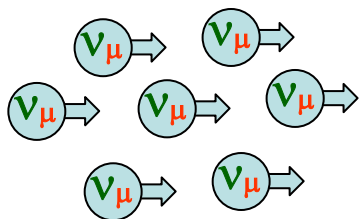
$$P(\nu_\mu \rightarrow \nu_\tau) = \sin^2 2\theta_{23} \sin^2 \left(\frac{\Delta m_{32}^2 L}{4E} \right) \quad \text{with} \quad \Delta m_{32}^2 = m_3^2 - m_2^2$$

• e.g. $\Delta m^2 = 0.003 \text{ eV}^2$, $\sin^2 2\theta = 0.8$, $E_\nu = 1 \text{ GeV}$



• wavelength

$$\lambda_{\text{osc}} = \frac{4\pi E}{\Delta m^2}$$

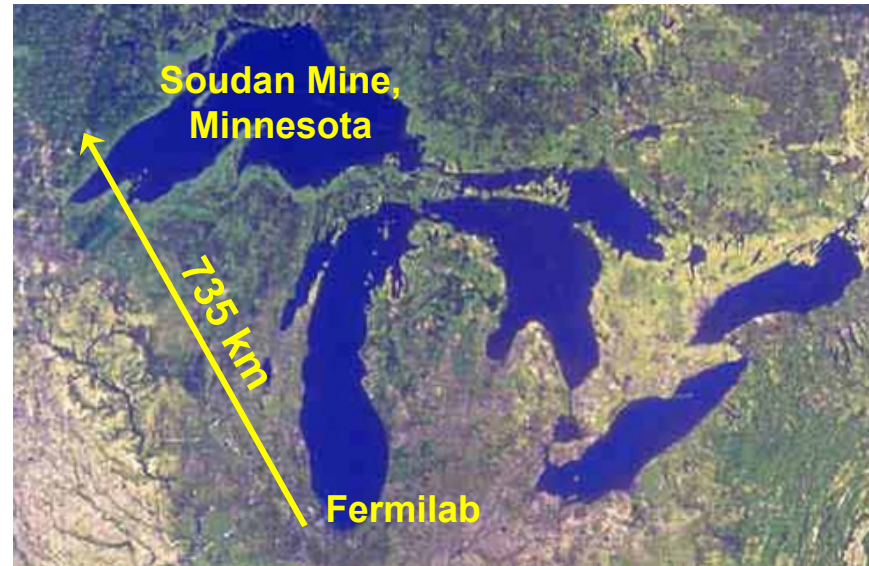




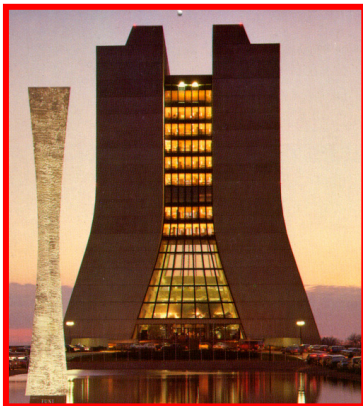
② The MINOS Experiment



- **120 GeV** protons extracted from the MAIN INJECTOR at Fermilab
- **2.5×10^{13}** protons per pulse hit target → very intense beam - **0.3 MW** on target



Two detectors:



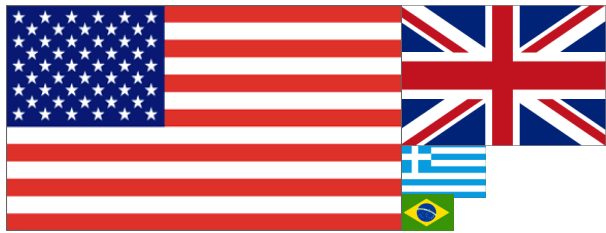
★ **1000 ton, NEAR Detector at Fermilab: 1 km from beam**

★ **5400 ton FAR Detector, 720 m underground in Soudan mine, N. Minnesota: 735 km from beam**





The MINOS Collaboration



**27 institutions
175 scientists**



Argonne • Athens • Benedictine • Brookhaven • Caltech • Cambridge • Campinas
Fermilab • Harvard • IIT • Indiana
Minnesota-Twin Cities • Minnesota-Duluth • Oxford • Pittsburgh • Rutherford
Sao Paulo • South Carolina • Stanford • Sussex • Texas A&M
Texas-Austin • Tufts • UCL • William & Mary



MINOS Physics Goals



- ★ **Precise Measurement of “atmospheric” osc. parameters**
 - $< 10\%$ measurement of Δm_{32}^2
 - Test whether $\sin^2 2\theta_{23}$ is maximal
- ★ **Demonstrate oscillation behaviour:**
 - observe oscillatory dip **and rise**
 - confirm neutrino flavour oscillations describe data
 - test alternative scenarios (sterile, decay, decoherence)
- ★ **Search for sub-dominant $\nu_{\mu} \rightarrow \nu_e$ oscillations**
 - first measurements of θ_{13} ?

+ MINOS is the 1st large deep underground detector with a B-field

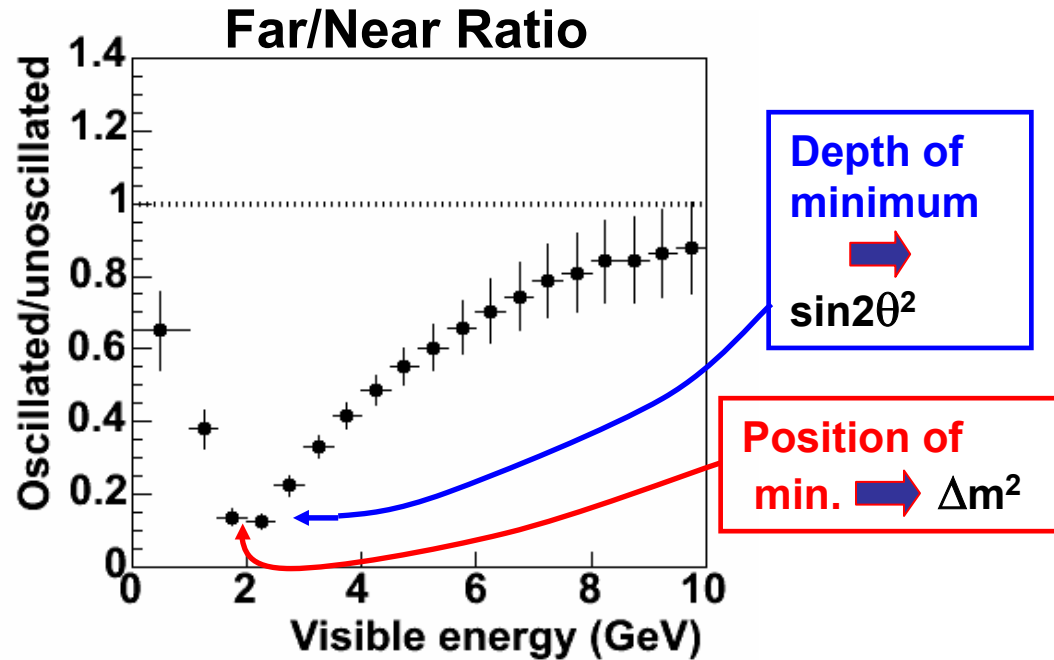
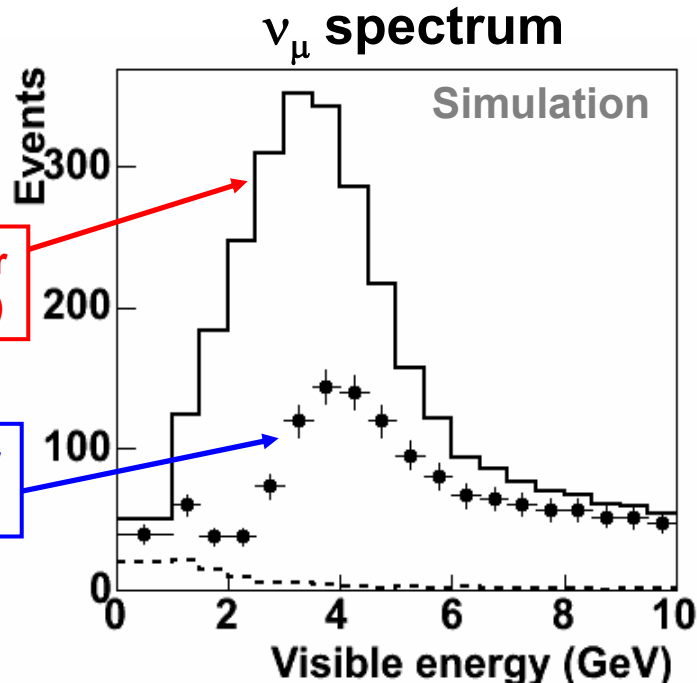
- ★ **Test of CPT invariance in neutrino sector**
 - first direct measurements of ν vs $\bar{\nu}$ oscillations from
 - atmospheric neutrino events
 - possibility of “reverse horn current” beam anti- ν run
- ★ **Cosmic-ray physics**



MINOS in a Nutshell



- ★ Intense neutrino beam: 0.2 MW
- ★ Two detectors: one close to beam the other 735 km away
- ★ Measure ratio of the neutrino energy spectrum in far detector (**oscillated**) to that in the near detector (**unoscillated**)



- ★ Two detectors vital to understand beam \Rightarrow precise measurements
- ★ Leads to a significant **cancellation** of systematic biases

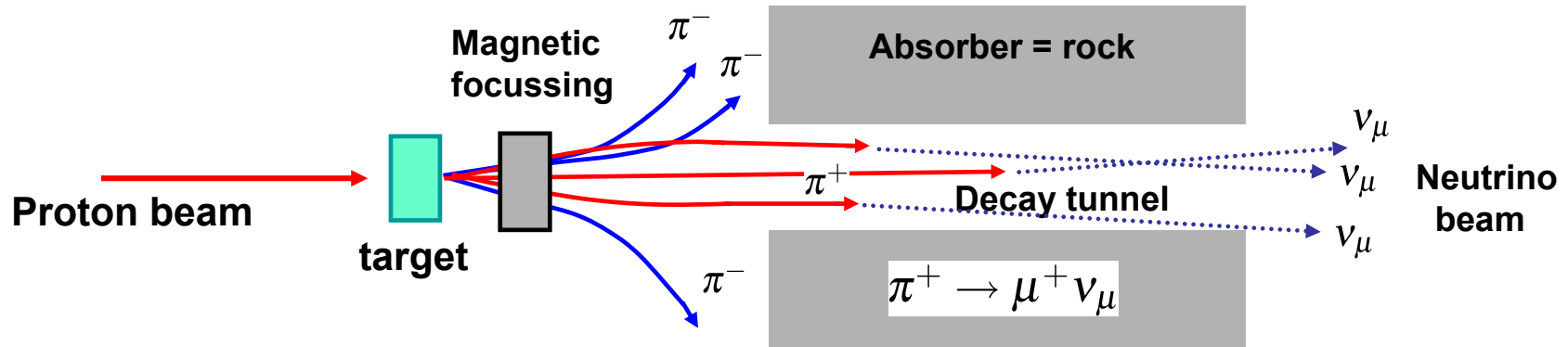


③ The NuMI Neutrino Beam



★ Neutrino Beams for beginners

- Smash high energy protons into a fixed target → hadrons
- Focus positive pions/kaons
- Allow them to decay $\pi^+ \rightarrow \mu^+ \nu_\mu$ + $K^+ \rightarrow \mu^+ \nu_\mu$ ($BR \approx 64\%$)
- Gives a beam of “collimated” ν_μ
- Could focus negative pions/kaons to give beam of $\bar{\nu}_\mu$





The NuMI beam in more detail



FERMILAB #98-765D

Basic Design:

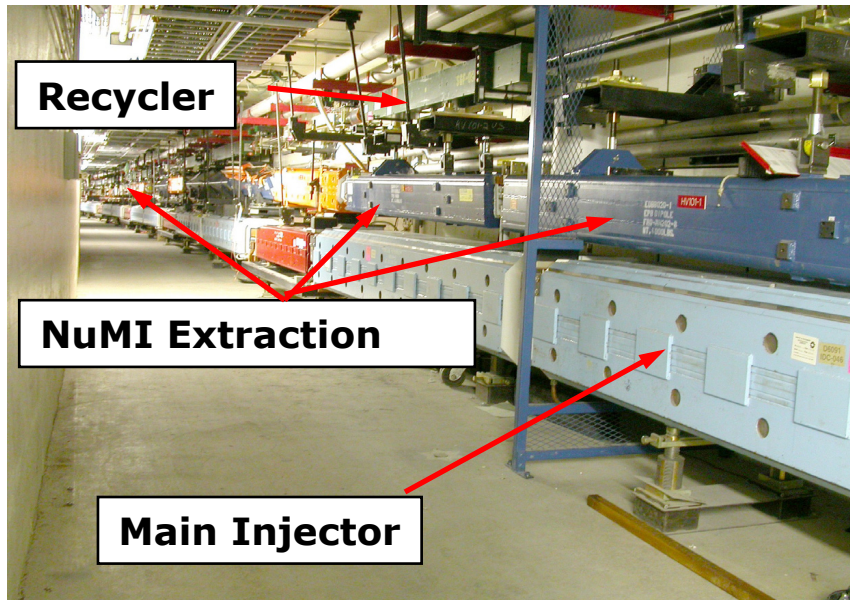
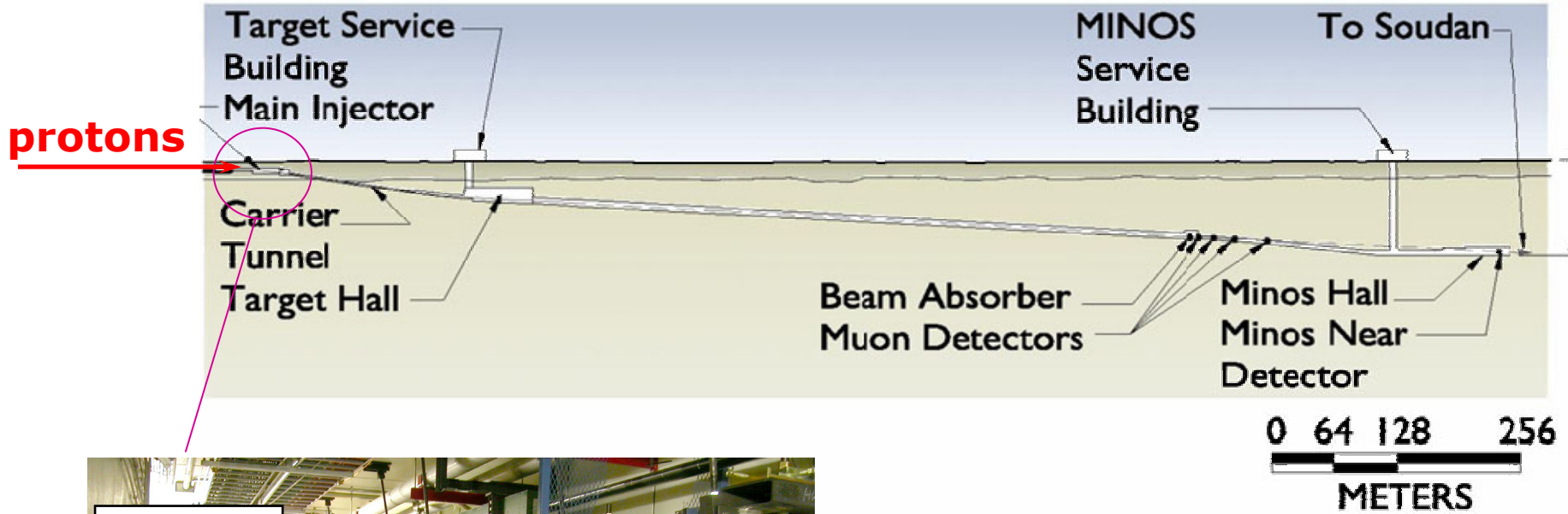
- ★ 120 GeV protons extracted from the **MAIN INJECTOR** in a single turn ($8.7 \mu\text{s}$)
- ★ 2.4 second cycle time
- ★ *i.e.* ν beam “on” for $8.7 \mu\text{s}$ every 2.4 seconds

Beam Performance (2007):

- ★ 2.4×10^{13} protons/pulse
- ★ 0.2 MW on target !
- ★ **Integrated intensity**
 2.5×10^{20} protons/year

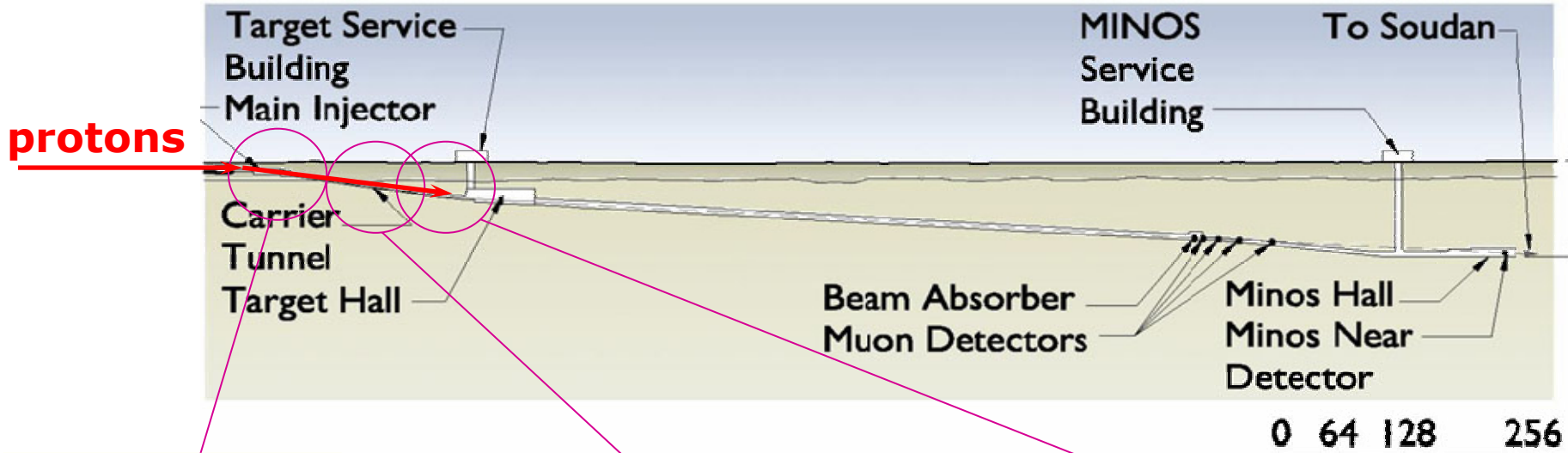


The NuMI ν beam : I





The NuMI ν beam : II



Steep incline



Carrier tunnel



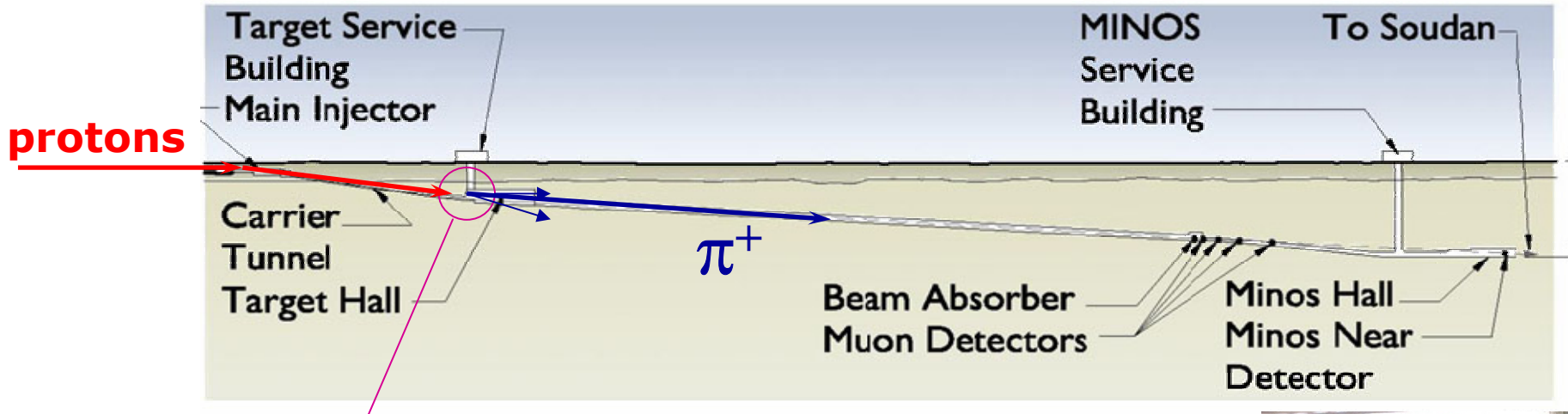
Pre-target



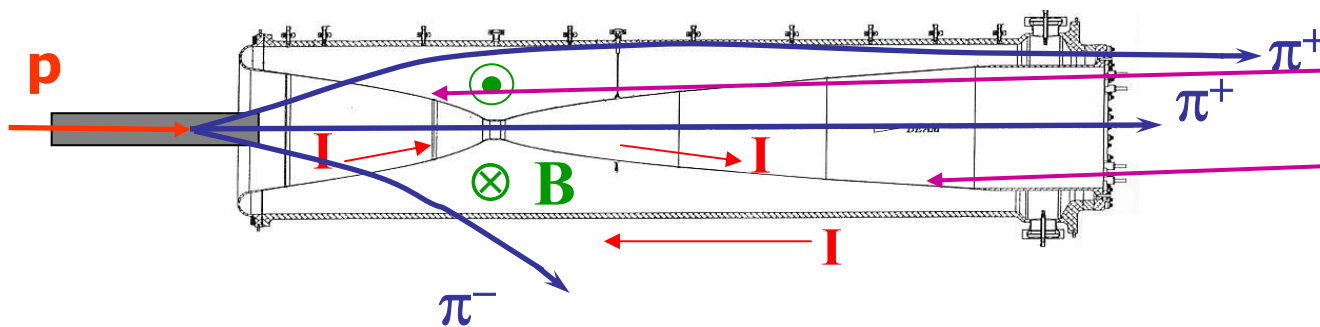
NuMI segmented graphite target with water cooling lines



The NuMI ν beam : III

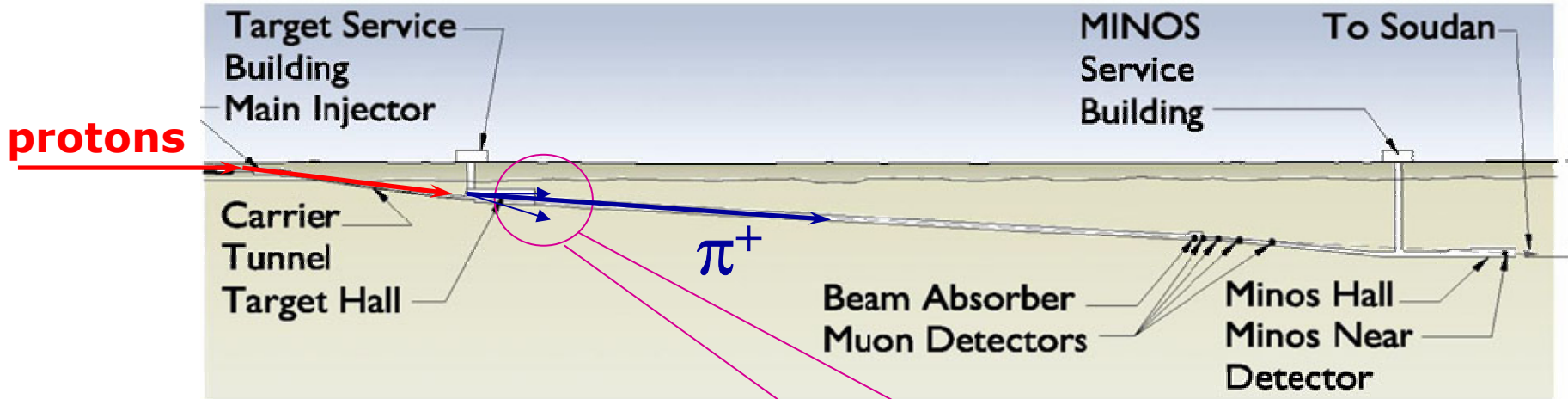


- Two focusing horns **pulsed with 200 kA**
- Toroidal Magnetic field $B \sim I/r$ between the inner and outer conductors
- Maximum field – 3 T



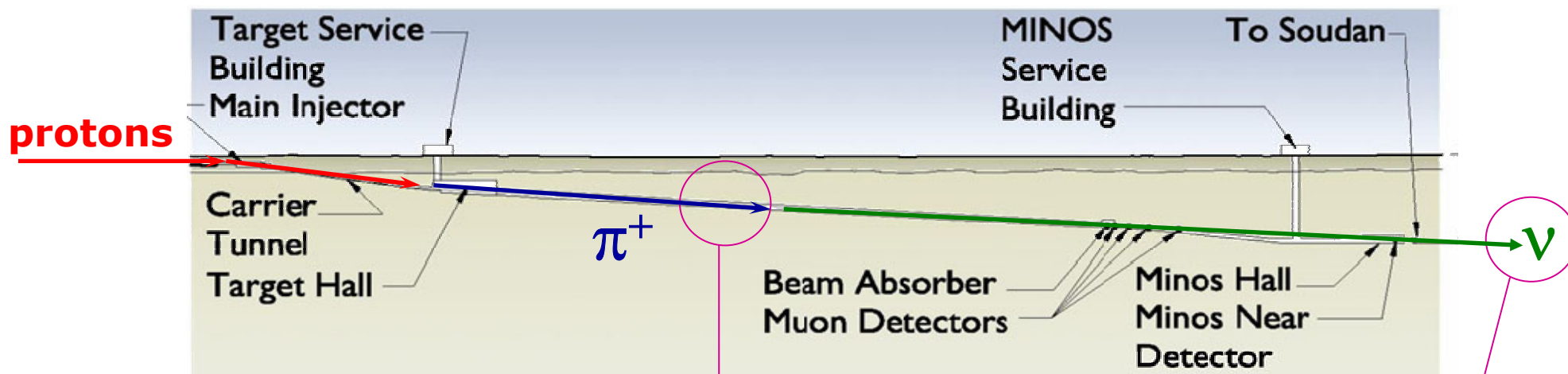


The NuMI ν beam : IV





The NuMI ν beam : ν



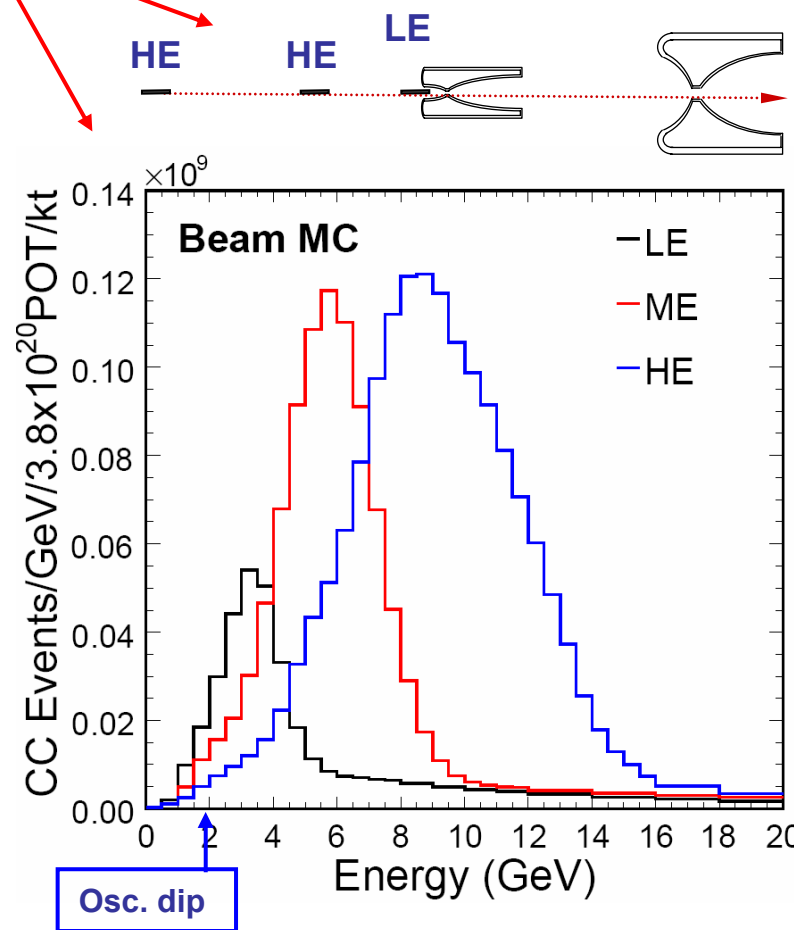
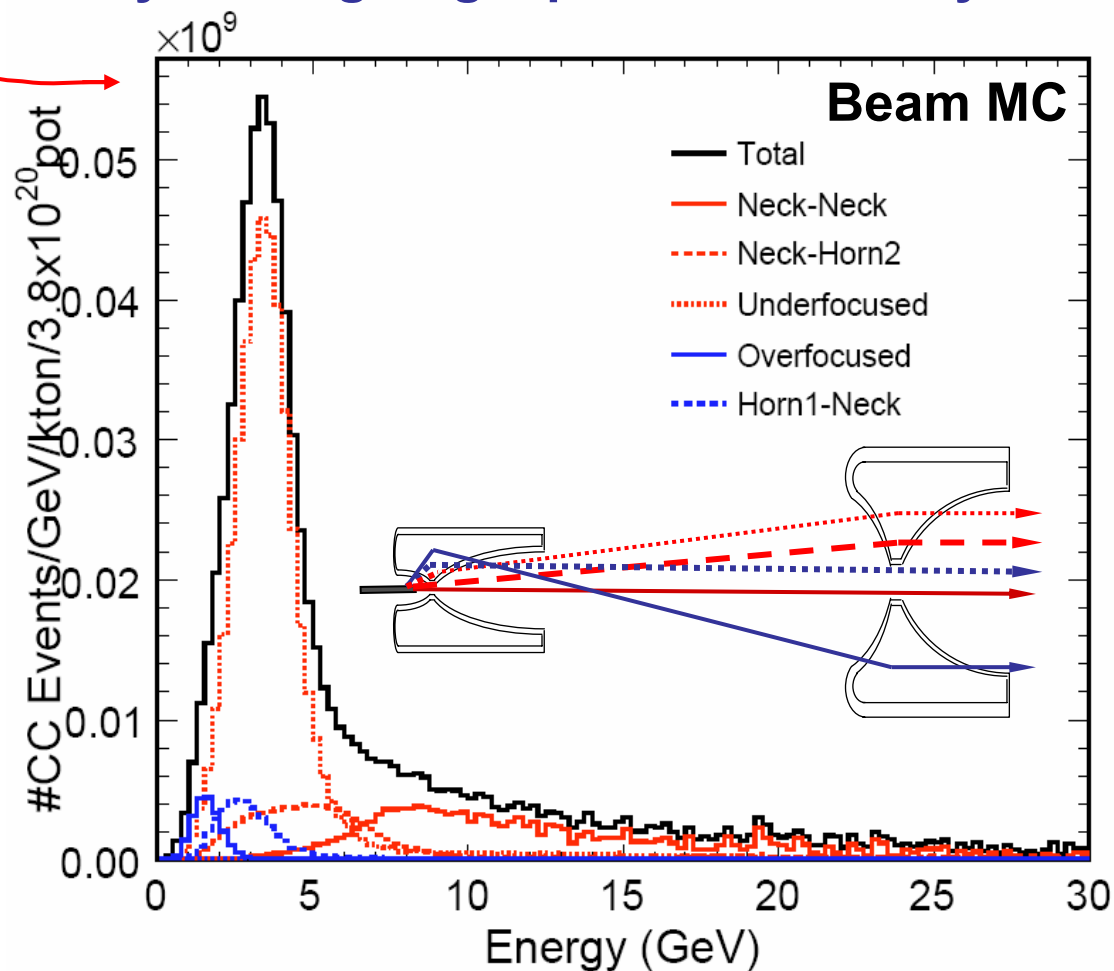
675 m long decay pipe

- ★ Need long decay pipe:
for a 5 GeV π^+
 $\gamma_{ct} \sim 200$ m
- ★ Evacuated to 1.5 Torr
- ★ Steel decay pipe installed
and encased in 2-3 m of
concrete to protect
ground water





- ★ Focussing is achieved with a **two horn** system
- ★ Behaves like a pair of (achromatic) lenses
- ★ By moving target position can vary energy spectrum (default = LE)

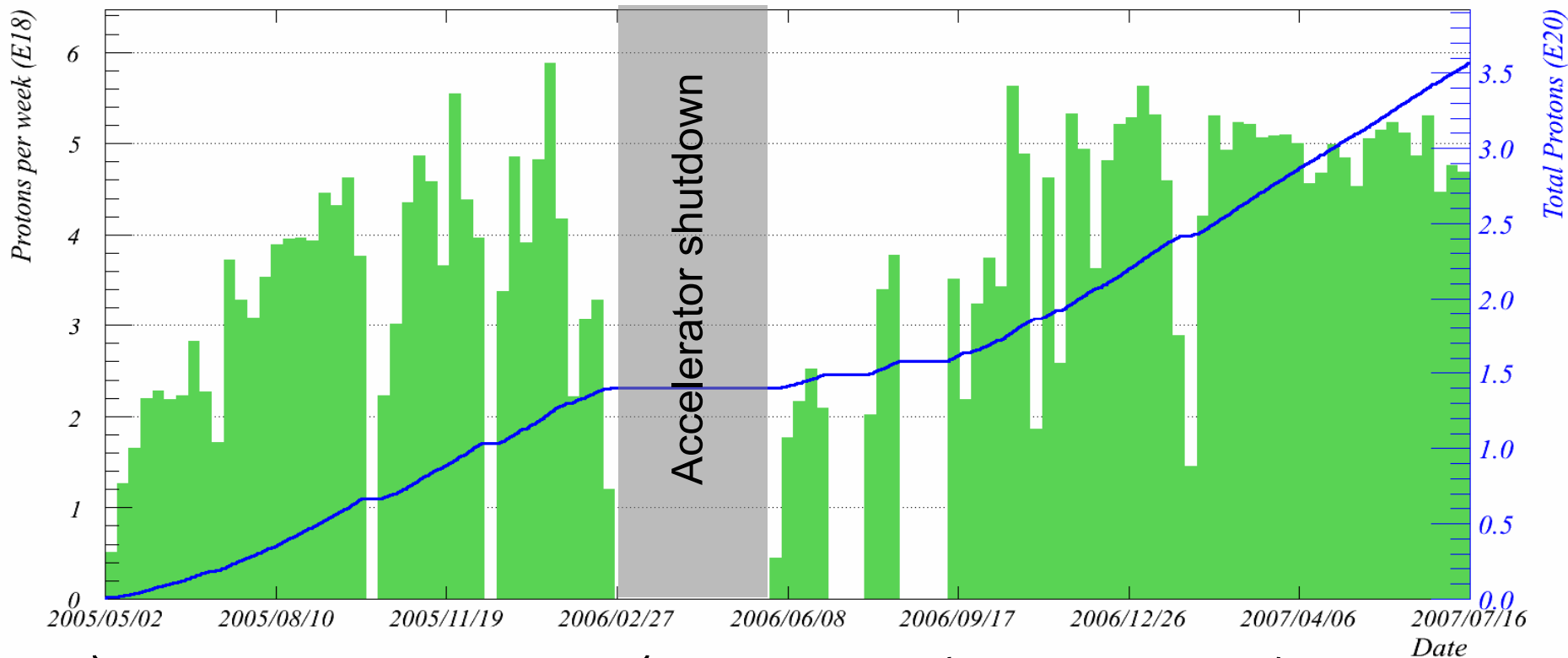




NuMI Performance



Total NuMI protons to 00:00 Monday 16 July 2007



RUN I -
1.27x10²⁰ POT
 (published in PRL)

Higher
 energy
 beam
 running

RUN IIa -
1.23x10²⁰ POT
 (arXiv:0711.0769)

RUN IIb - 0.8x10²⁰
POT
 (Analysis in
 progress)

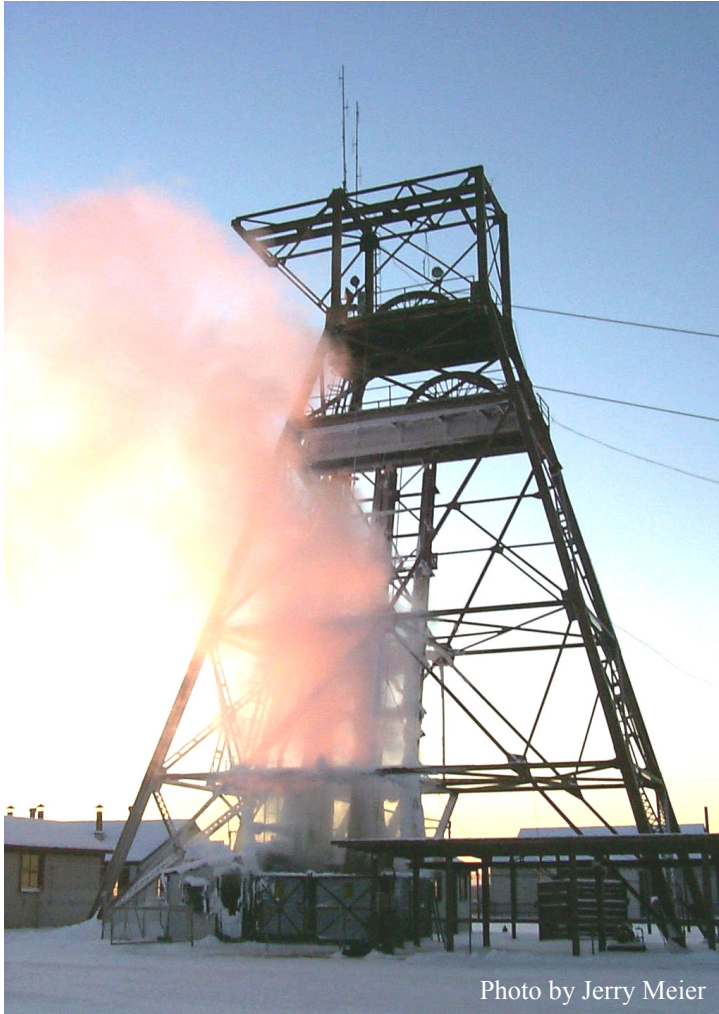
Results presented here: **Run I + Run IIa - 2.5x10²⁰ POT**



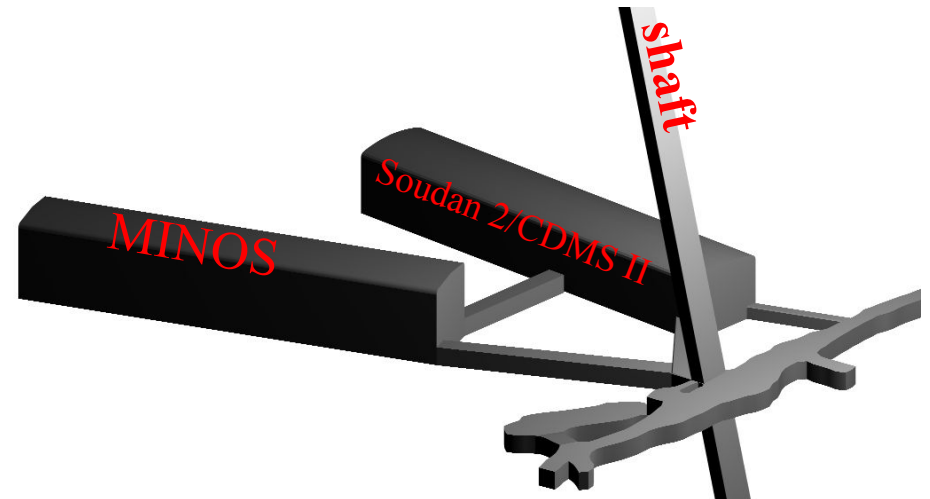
4 The MINOS Detectors



Far Detector in deepest darkest Minnesota – nearest main town **Ely, MN**



2070 mwe

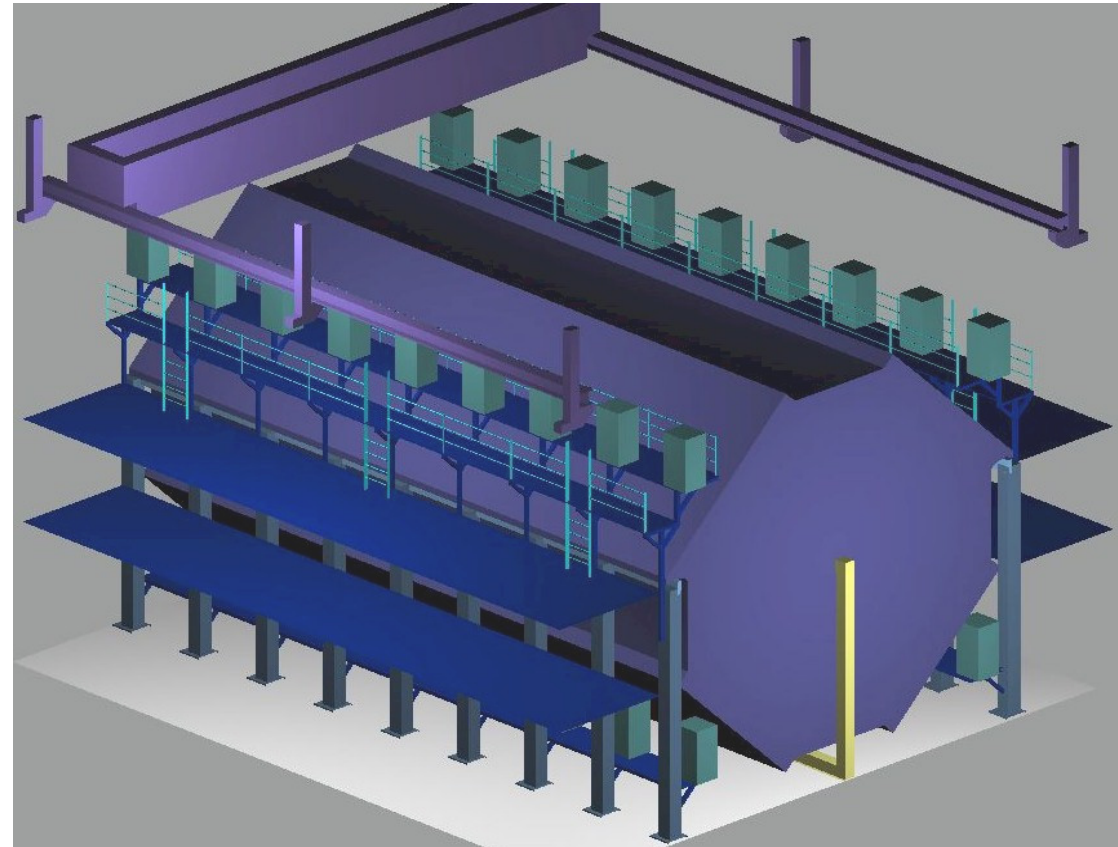




MINOS Far Detector



- **8m octagonal steel & scintillator tracking calorimeter**
 - 2 sections, 15m each
 - 5.4 kton total mass
 - $55\%/\sqrt{E}$ for hadrons
 - $23\%/\sqrt{E}$ for electrons
- **Magnetized Iron ($B \sim 1.2$ T)**
- **484 planes of scintillator**



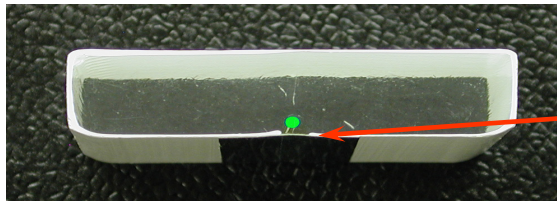
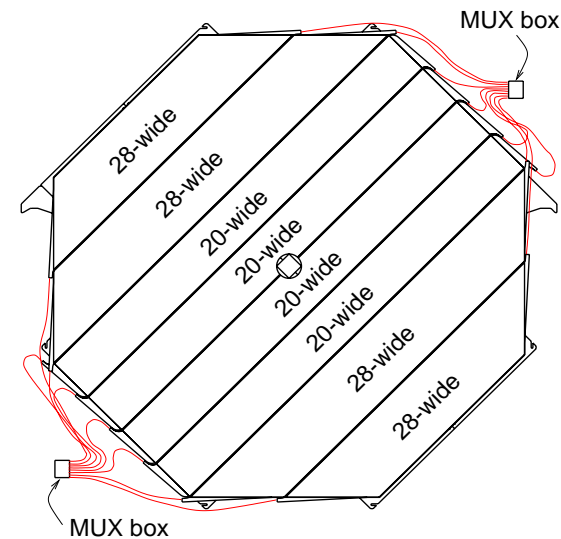
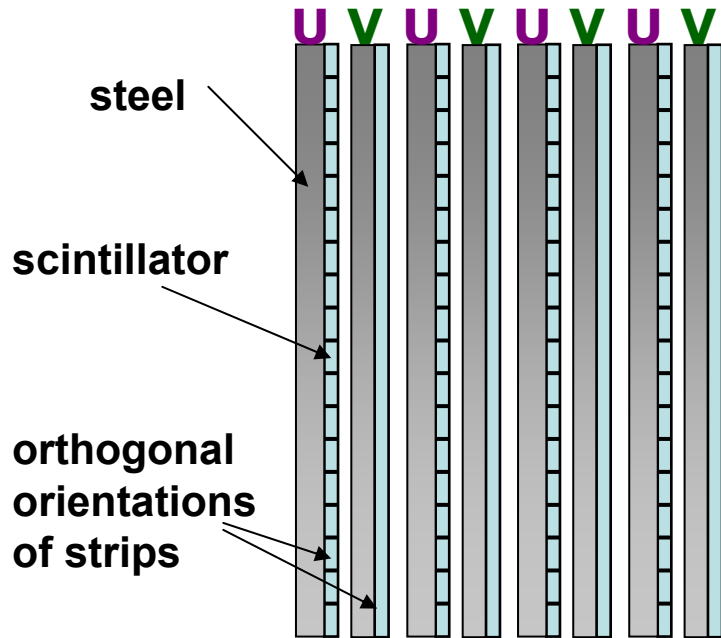
**One Supermodule of the Far Detector...
Two Supermodules total.**



Basic Detector Elements



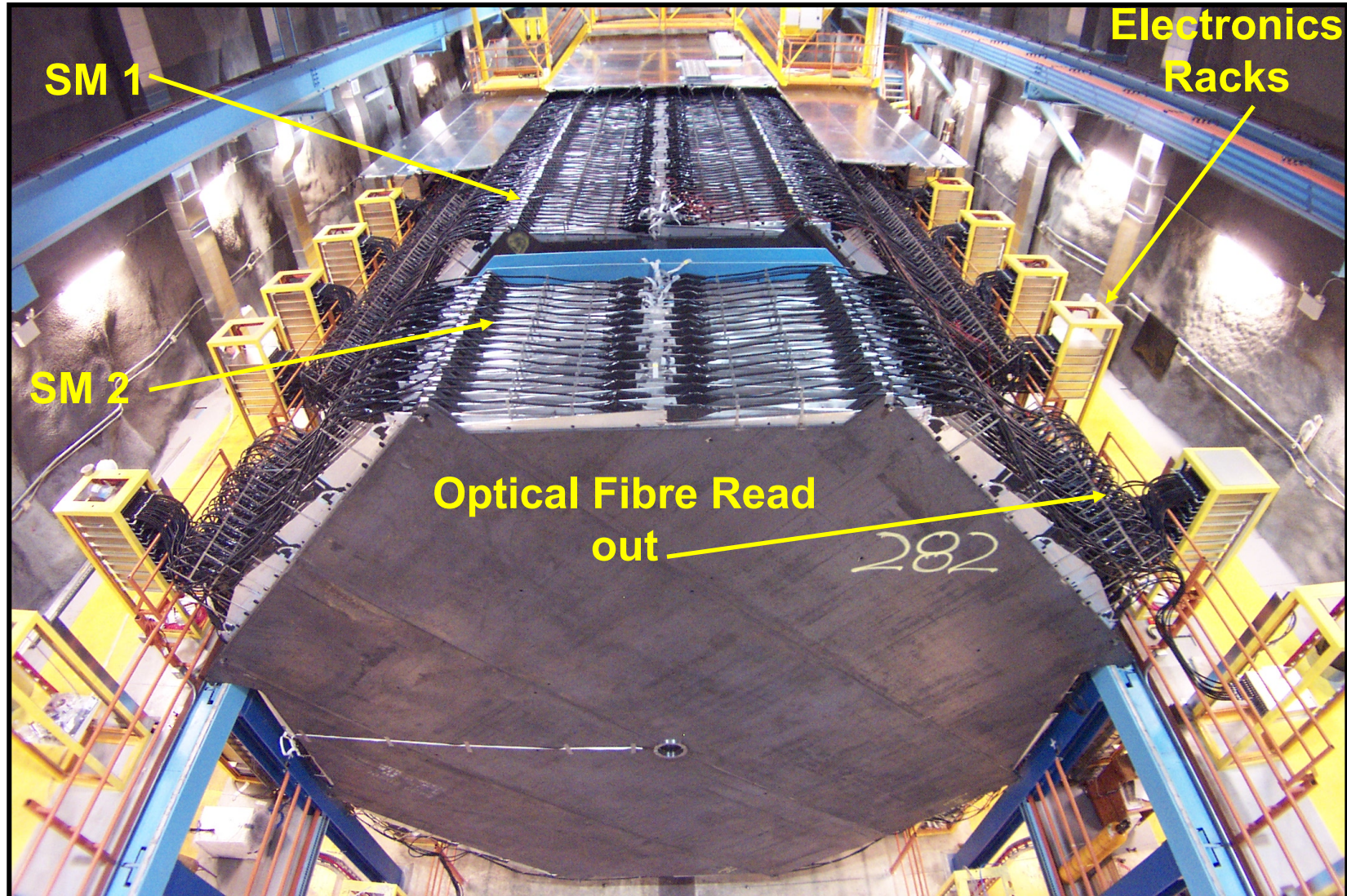
- ★ **Steel-Scintillator sandwich : SAMPLING CALORIMETER**
- ★ **Each plane consists of a 2.54 cm steel +1 cm scintillator**
- ★ **Each scintillator plane divided into 192 x 4cm wide strips**
- ★ **Alternate planes have orthogonal strip orientations (U and V)**



- ★ **Scintillation light collected by WLS fibre glued into groove**
- ★ **Readout by multi-pixel PMTs**

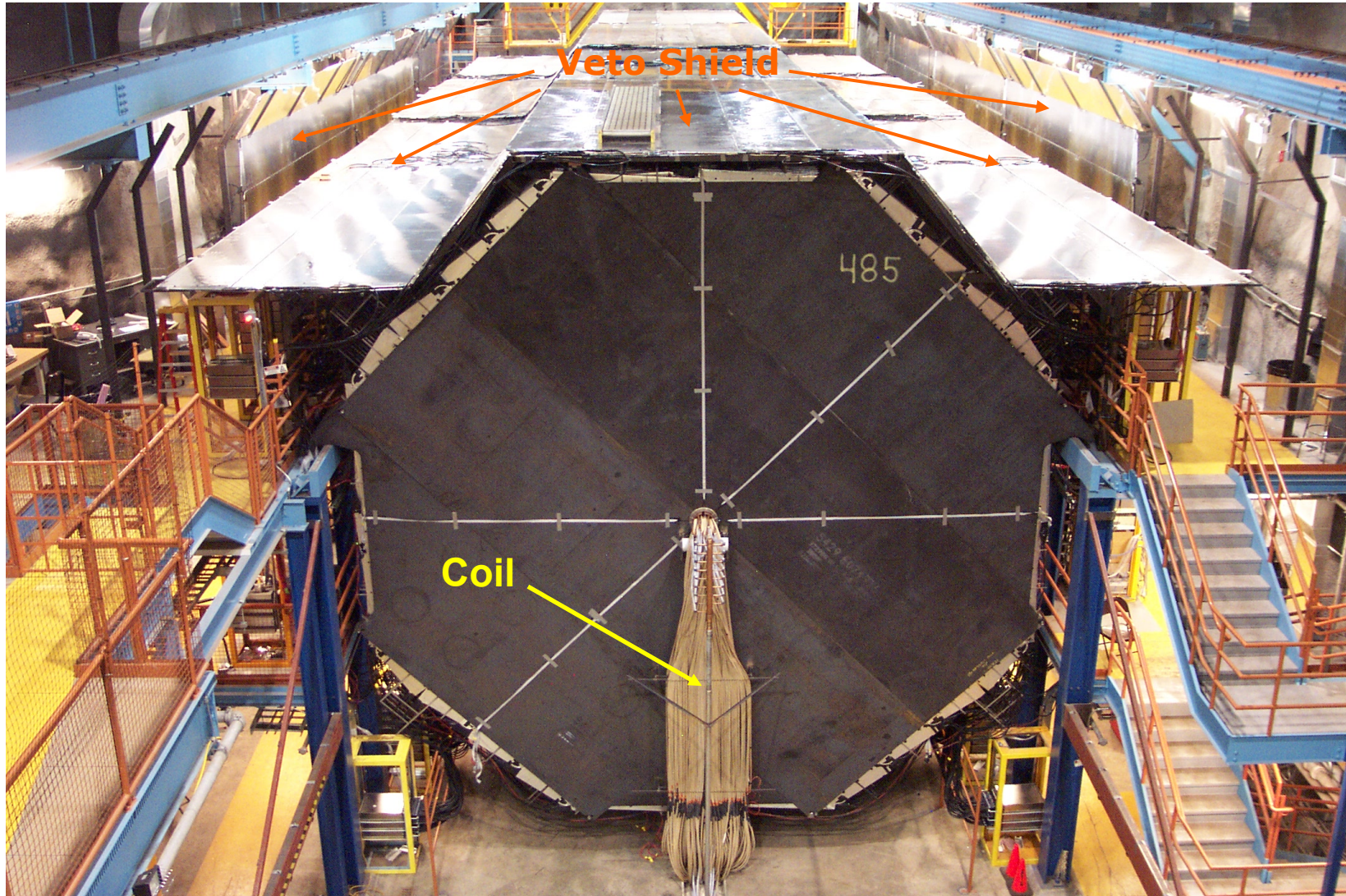


MINOS Far Detector during installation





Far Detector : fully operational since July 2003

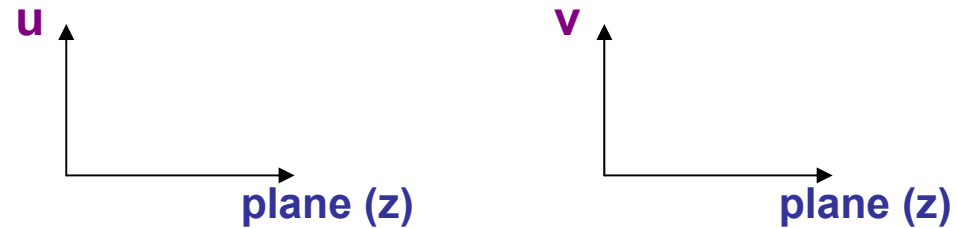




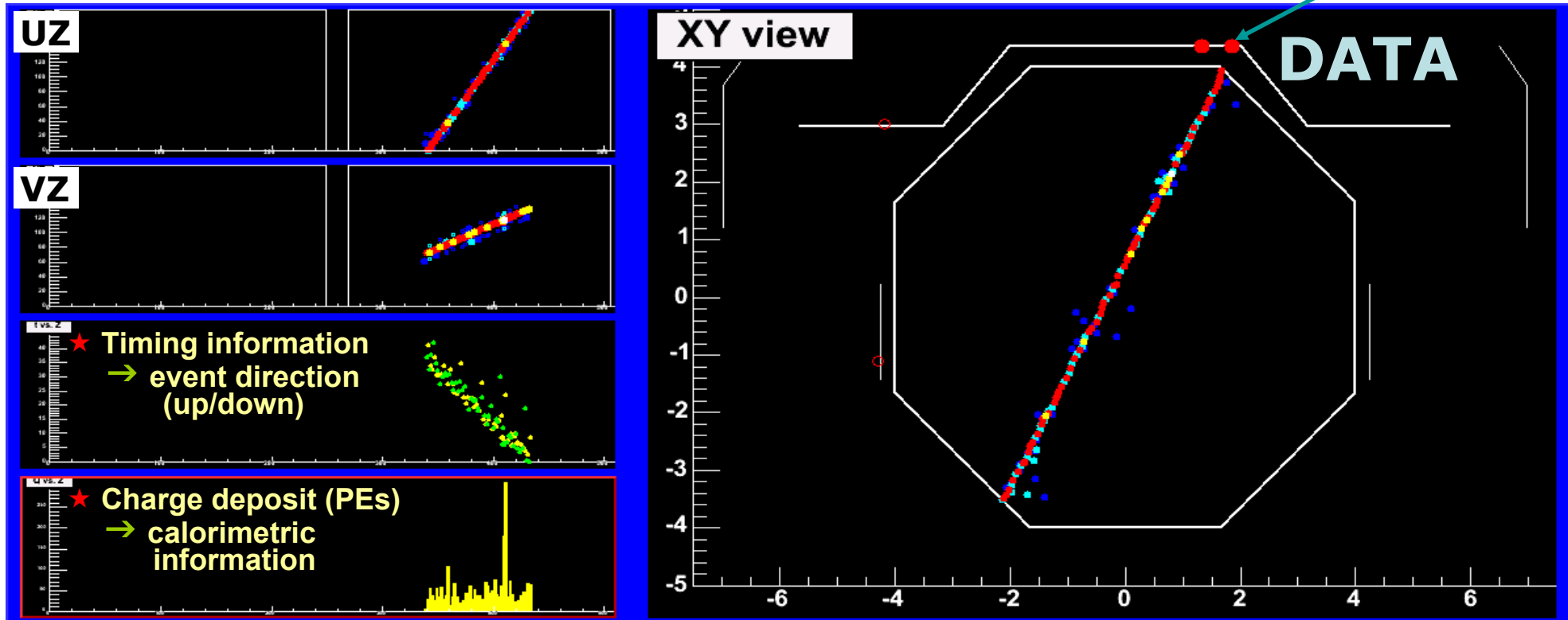
Event Information



★ Two 2D views of event



★ Software combination to get '3D' event



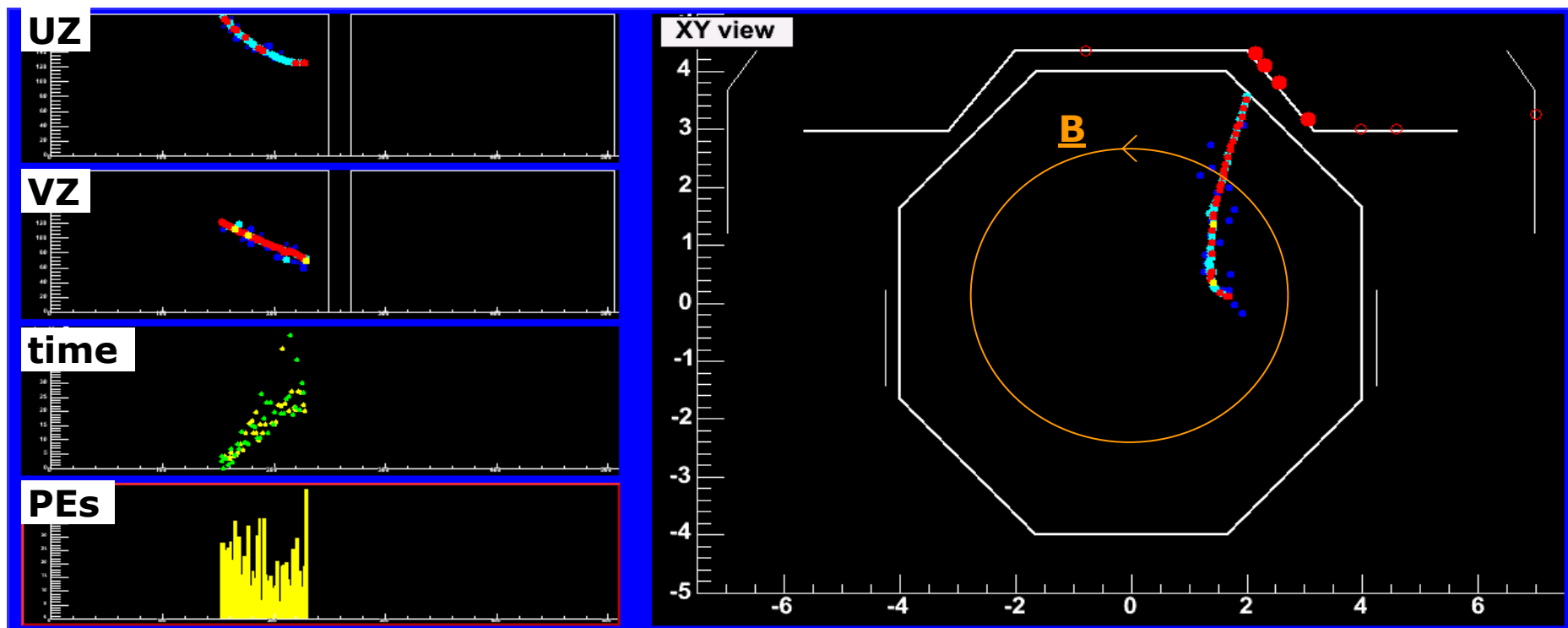


B-Field



~1.3 T Magnetic Field

- ★ Charge separation
- ★ Momentum measurement from curvature



Single hit timing res. : 2.5 ns

Stopping cosmic-ray muon:

$$P_{\text{range}} = 3.86 \text{ GeV/c}$$

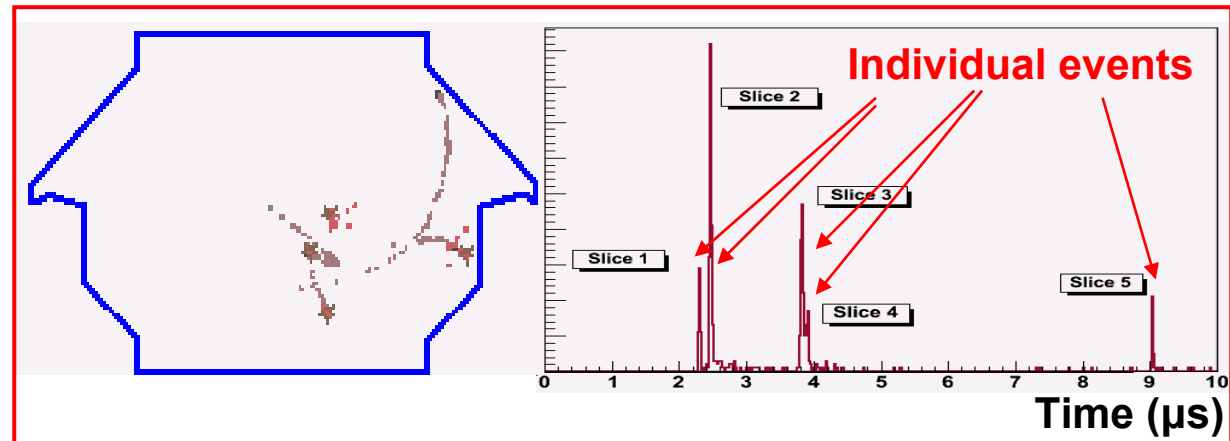
$$P_{\text{curvature}} = 4.03 \text{ GeV/c}$$



MINOS Near Detector



- ★ 1 km from beam
- ★ 1 kton total mass
- ★ Same basic design as Far Detector
steel, scintillator, etc
- ★ But some differences:
 - ◆ Faster electronics
 - ◆ Different PMTs (M64 vs M16)
 - ◆ Different triggering
 - ◆ Only partially instrumented
 - ◆ 282 planes of steel
 - ◆ 153 planes of scintillator
 - ◆ (Rear part only used to track muons)
- ★ But the main difference is
EVENT RATE
- ★ Multiple event interactions per beam spill
- ★ Separated using timing + spatial information



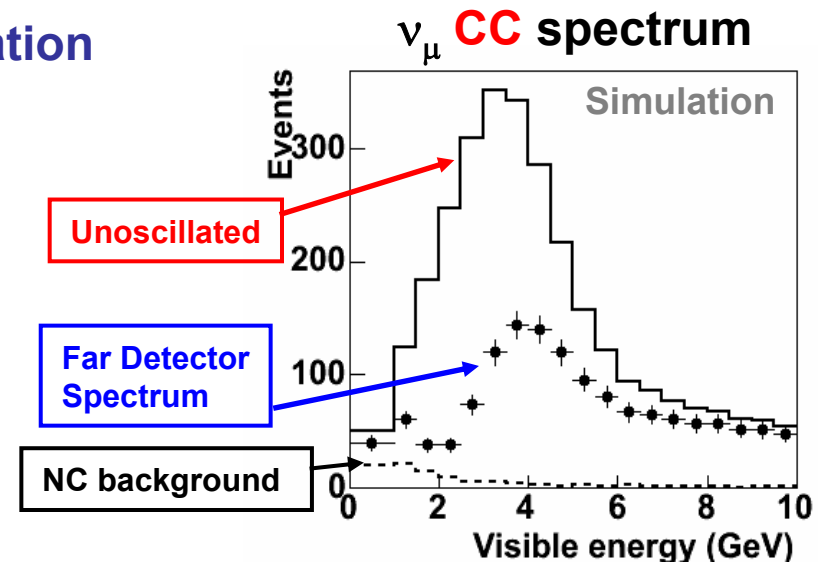
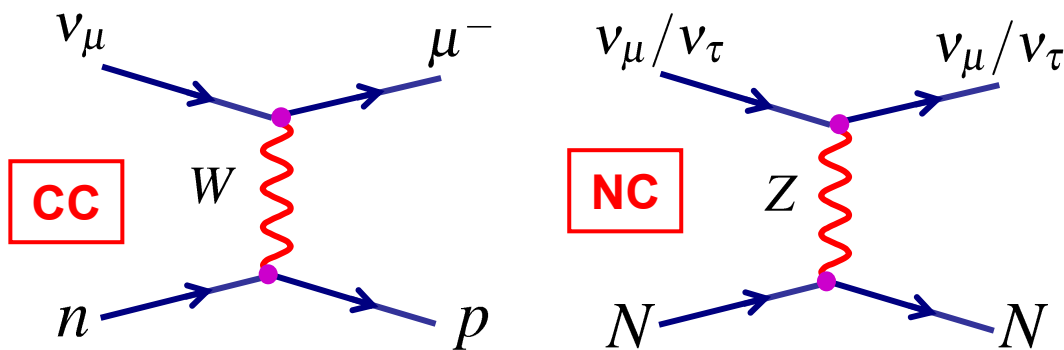


5 ν_μ Oscillation Analysis



- ★ MINOS neutrino beam is **93 % ν_μ** , **6 % $\bar{\nu}_\mu$** , **1 % ν_e** , and **0.1 % $\bar{\nu}_e$**
- ★ For values of Δm_{32}^2 from atmospheric neutrinos oscillation minimum at **~ 2 GeV**
- ★ In **region where oscillations occur** – predominantly ν_μ
- ★ Oscillations expected to be predominantly $\nu_\mu \rightarrow \nu_\tau$ (see later for ν_e)
- ★ However threshold for **Charged Current (CC) ν_τ interactions** is

$$E_{\nu_\tau} > 2m_N m_\tau \sim 3.5 \text{ GeV}$$
- ★ Oscillated $\nu_\mu \rightarrow \nu_\tau$ mostly below/close to CC threshold – effectively disappear
- ★ Analysis strategy:
 - Identify CC ν_μ interactions (i.e. reject NC interactions)
 - Reconstruct neutrino energy
 - Compare Far Detector Spectrum to expectation

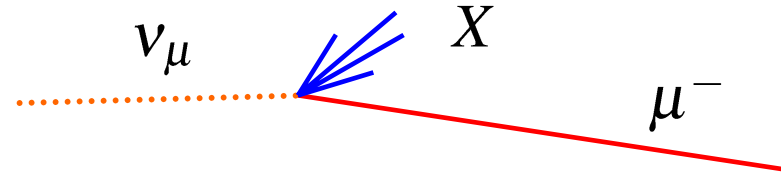
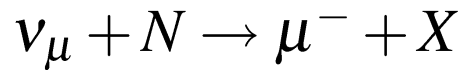




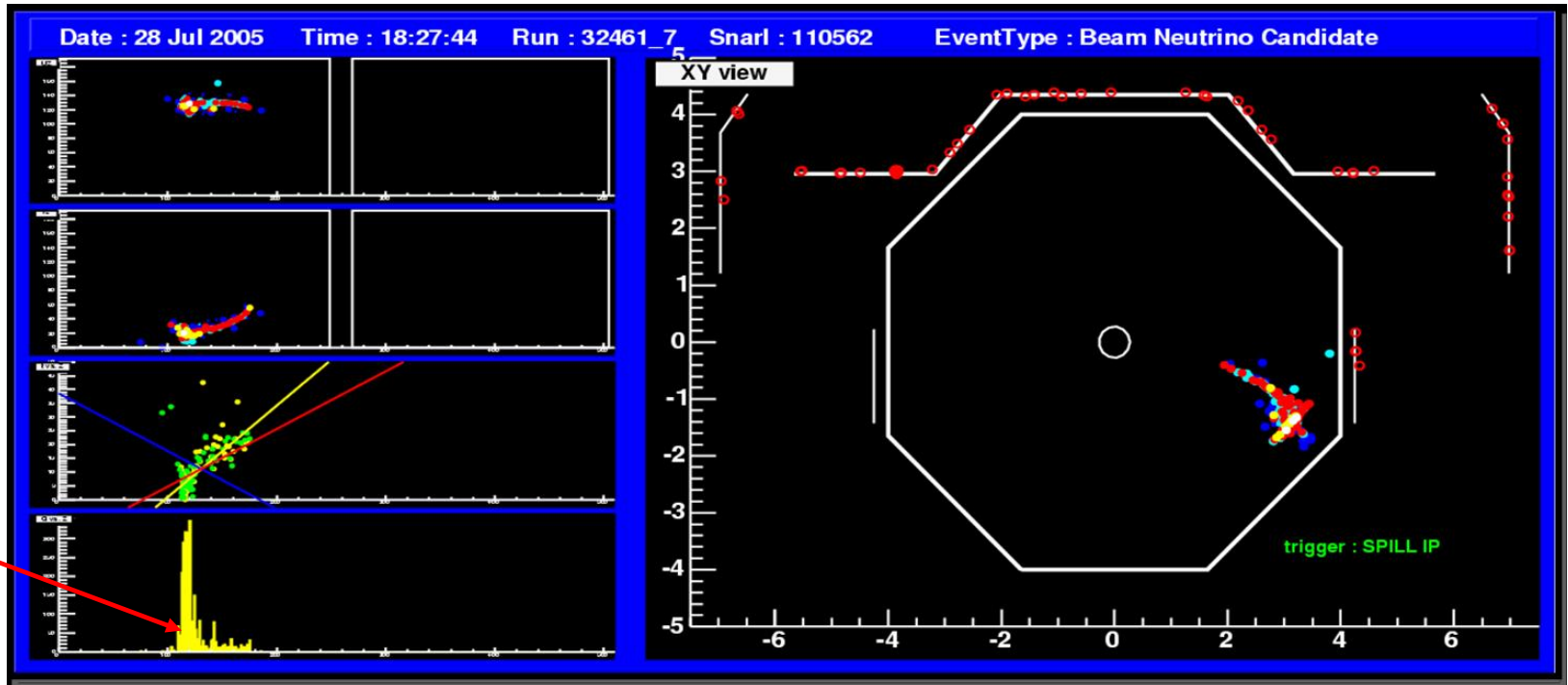
Event reconstruction



- Neutrino detection via CC interactions on nucleon ($\sim 5/\text{day}$ in FD)



Example event:



- Reconstruct muon momentum + energy of hadronic system

$$E_{\nu} = E_{\mu} + E_X$$

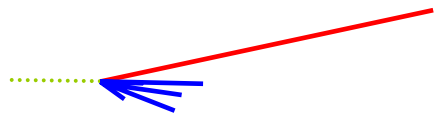
$$y = E_X / (E_{\mu} + E_X)$$



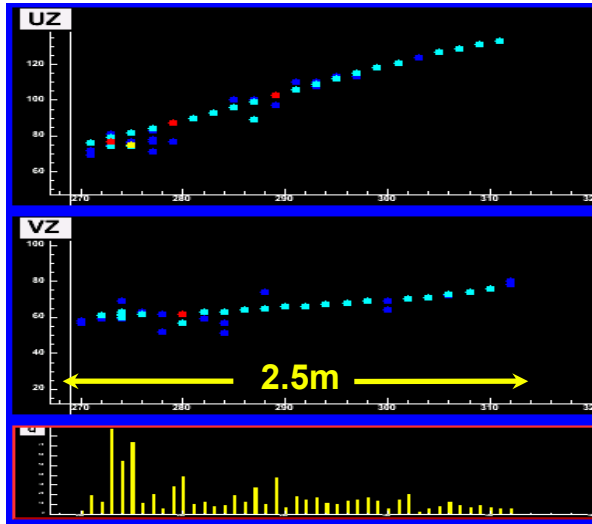
Event Identification



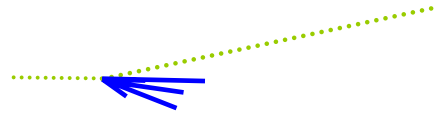
★ Different Neutrino interactions have very different event topologies



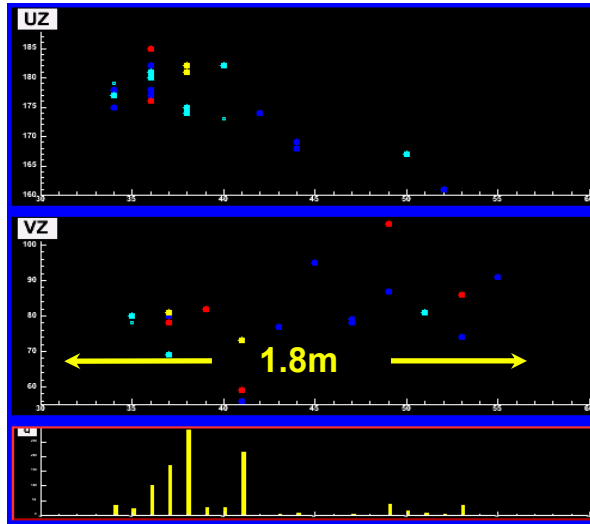
$\nu_\mu CC$



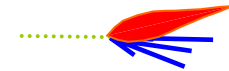
- ◆ Clear muon track
- ◆ Hadronic activity at interaction vertex



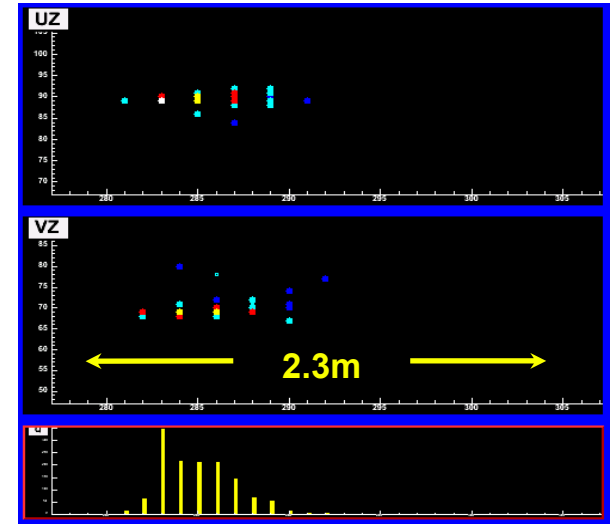
NC



- ◆ Short
- ◆ Diffuse



$\nu_e CC$



- ◆ Compact EM shower
- ◆ +Hadronic activity

★ Use multivariate likelihood method to select $\nu_\mu CC$ events in NEAR and FAR detectors

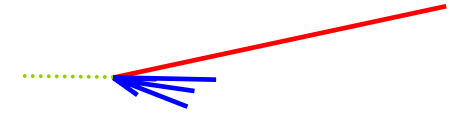
Monte Carlo



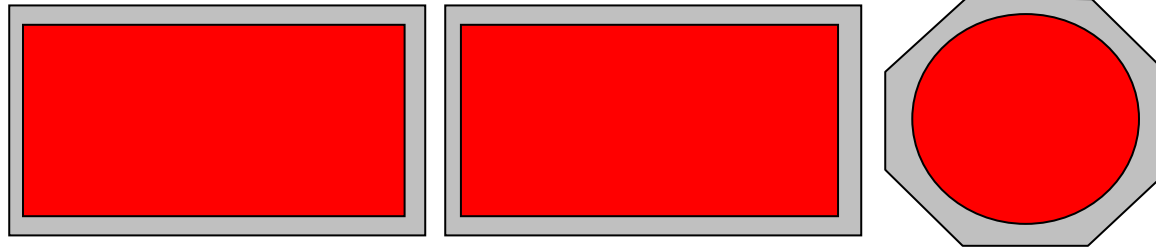
Event selection cuts : Near and Far



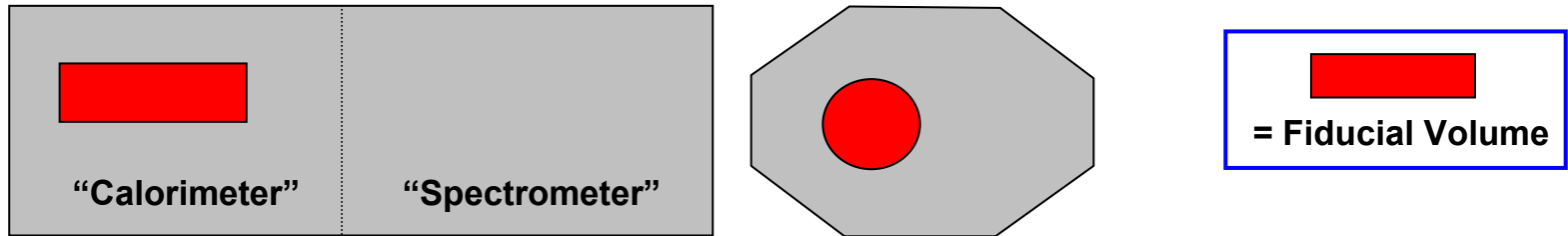
- ★ ν_μ CC candidate events are selected by requiring:
 - The event must have a good reconstructed track
 - The reconstructed track vertex must lie in the detector fiducial volume (avoid edges and less well understood regions of detector)



FAR DETECTOR



NEAR DETECTOR



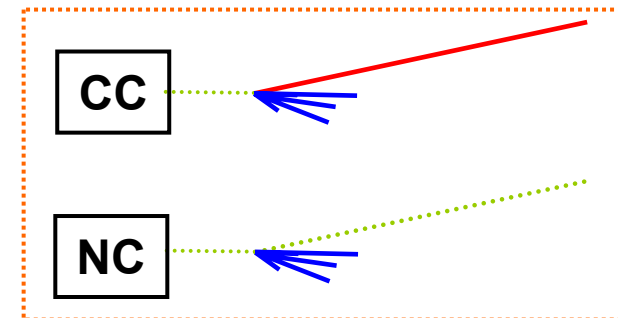
- ★ Likelihood selection to separate CC and NC events using 7 reconstructed quantities

Track Topology Variables:

- ◆ Track Pulse Height Per Plane
- ◆ Number of Track-Like Planes
- ◆ Number of Planes
- ◆ Goodness of Muon Track Fit
- ◆ Reconstructed Track Charge

Event Variables:

- ◆ Neutrino Energy
- ◆ y

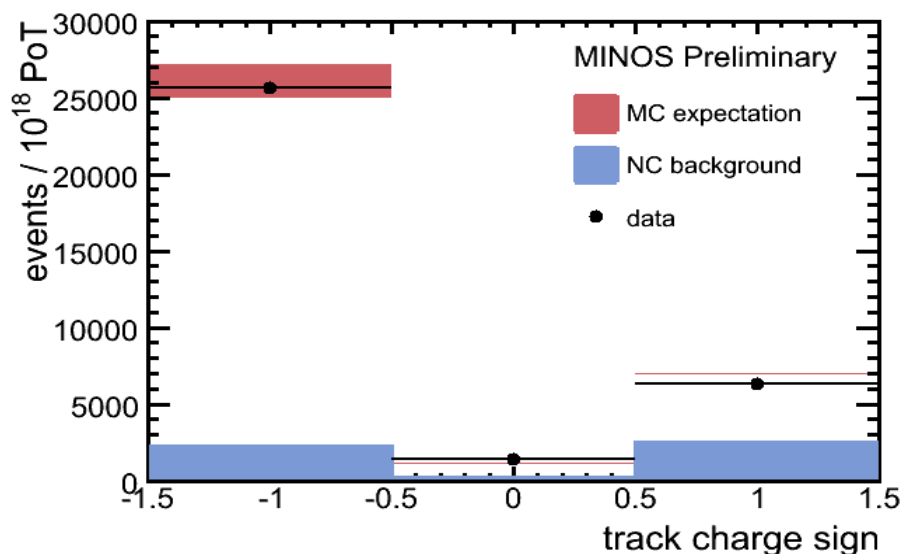
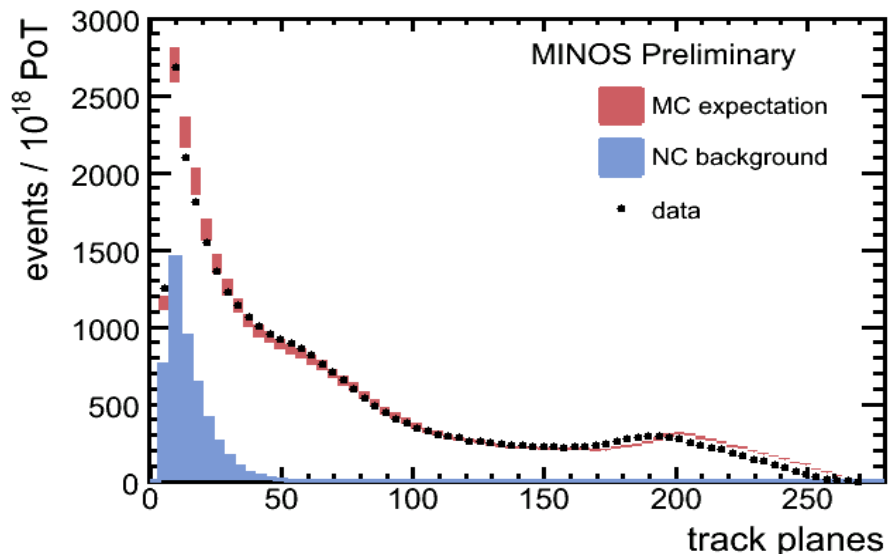




Near detector Data/MC comparisons: PID inputs



e.g.

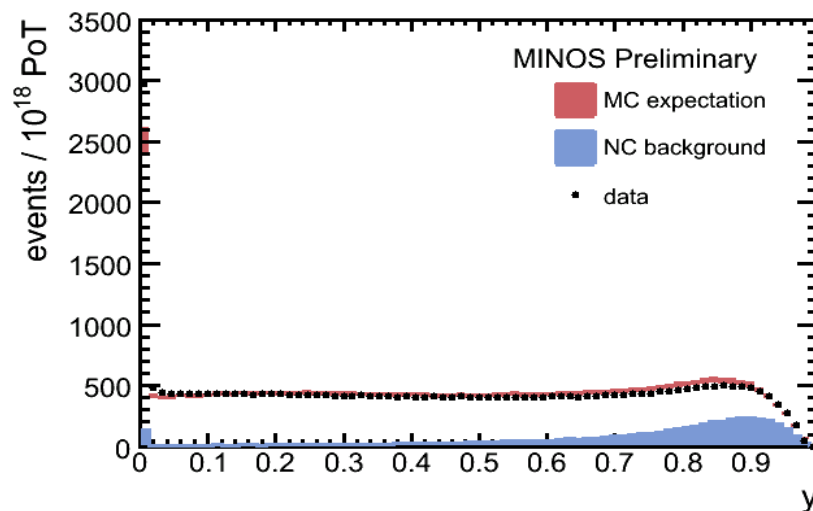
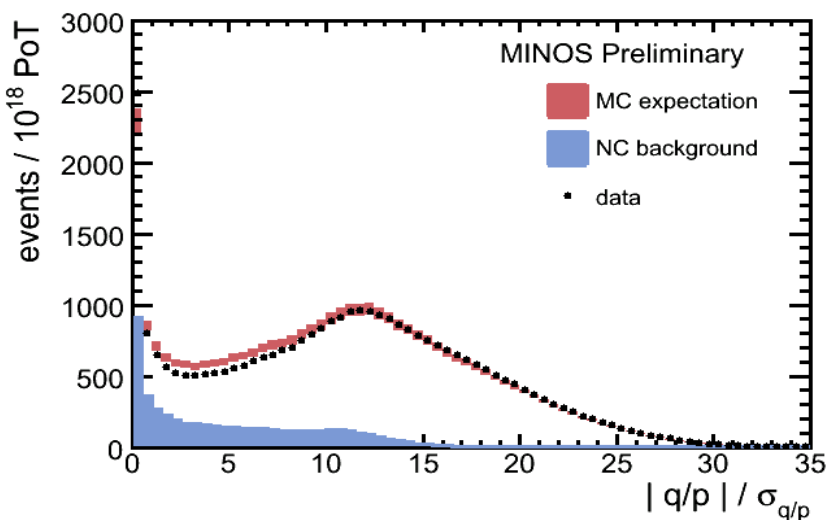
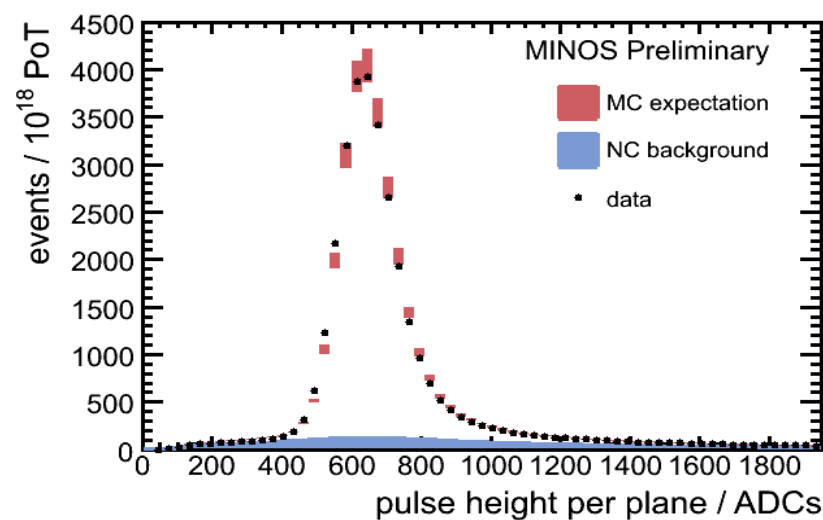
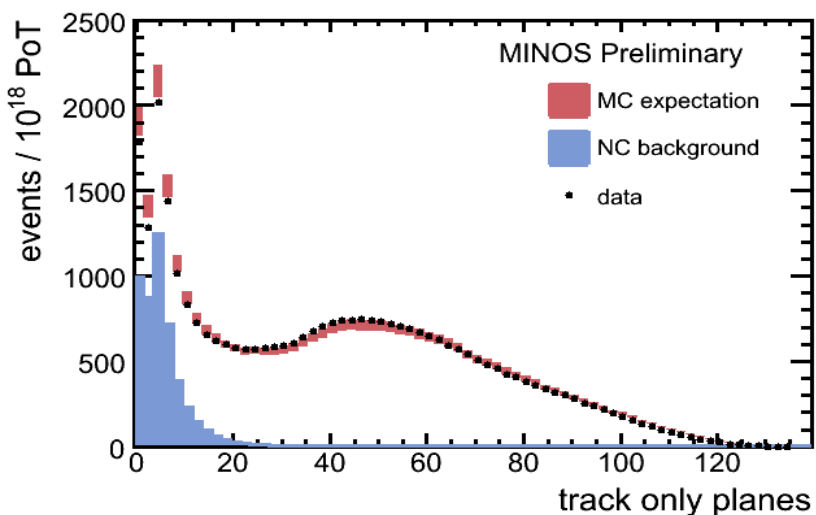


- ★ **Track Planes** (essentially event length)
 - CC events much longer due to presence of muon track

- ★ **Track charge sign** (from curvature)
 - ν_μ CC events produce a μ^-
 - Track in NC events usually a π^\pm from hadronic system or a fake track
 - In either case equally likely to be \pm



Near detector Data/MC comparisons: PID inputs, cont.



★ High statistics Near Detector data demonstrates that all variables are reasonably well modelled !



Combine into Likelihood discriminant



- ★ Use MC to create **NC** and **CC** probability density functions (pdfs) for each variable
- ★ Using pdfs calculate probability that an event is consistent with being **NC** and **CC**

$$P_{CC} = \prod_{i=1,7} P_i(x_i|CC) = P(x_1|CC) \cdot P(x_2|CC) \dots$$

$$P_{NC} = \prod_{i=1,7} P_i(x_i|NC)$$

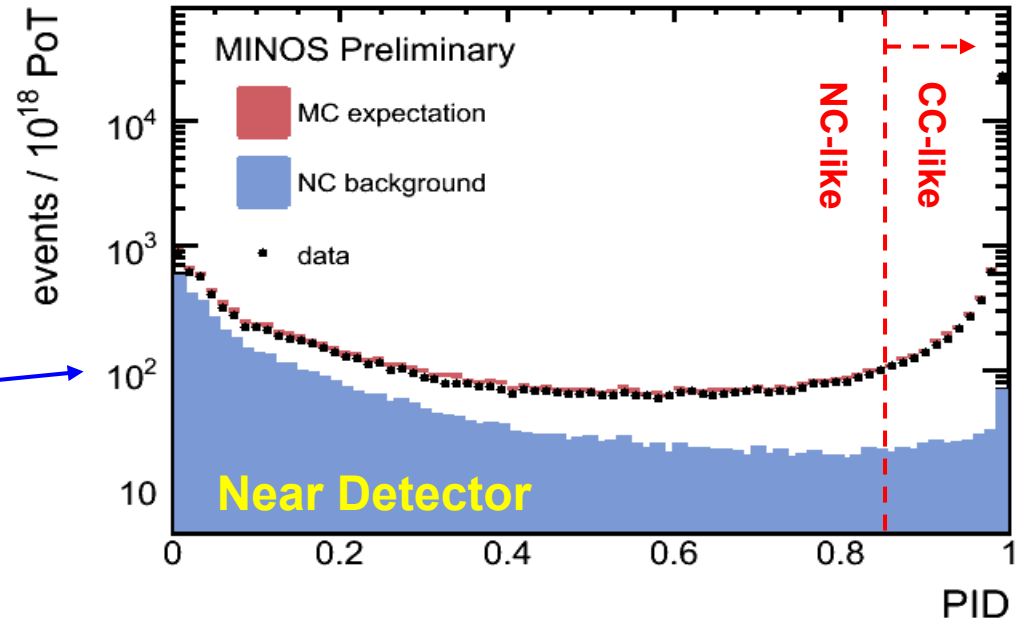
from CC PDFs for individual variables

- ★ Combine in to event “particle” identification variable “PID”

$$PID = \frac{P_{CC}}{P_{CC} + P_{NC}}$$

- ★ Require $PID > 0.85$

- ★ PID variable well-modelled in **Near Detector data**

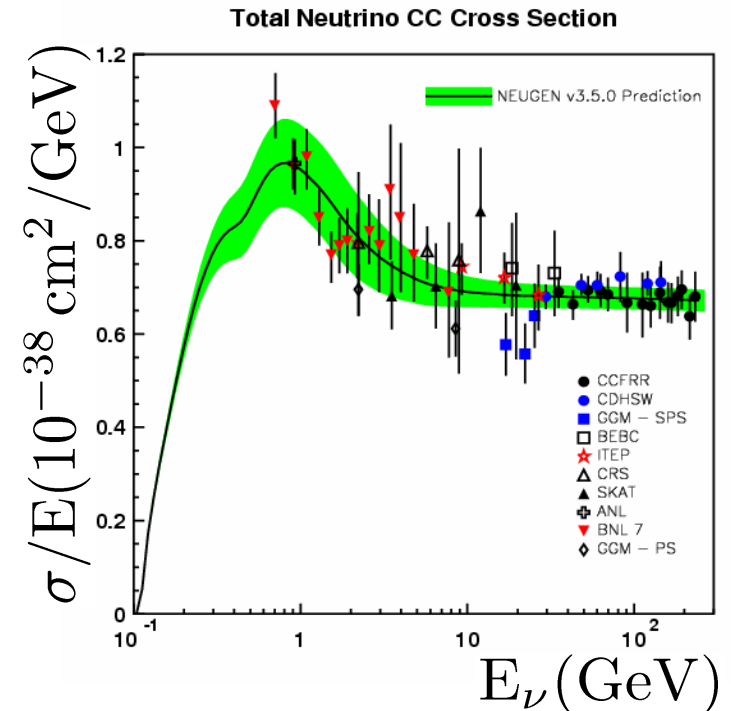
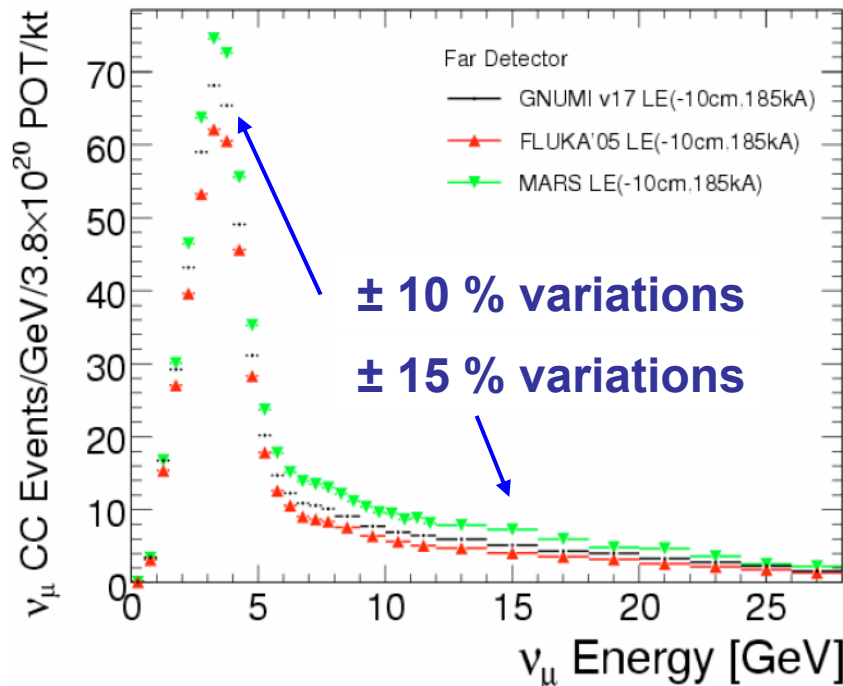




Neutrino Energy Spectrum



- ★ We've covered the easy part – i.e. selecting CC like neutrino interactions
- ★ Now want to compare CC neutrino energy spectrum in Far Detector to Monte Carlo expectation with and without oscillations, and fit etc.
- ★ To do this need to be able to accurately predict expected event rate
- ★ Require:
 - accurate simulation of neutrino flux from 120 GeV protons hitting target
 - accurate simulation of (low energy) neutrino cross sections
- ★ NEITHER EXIST – due to lack of appropriate data !

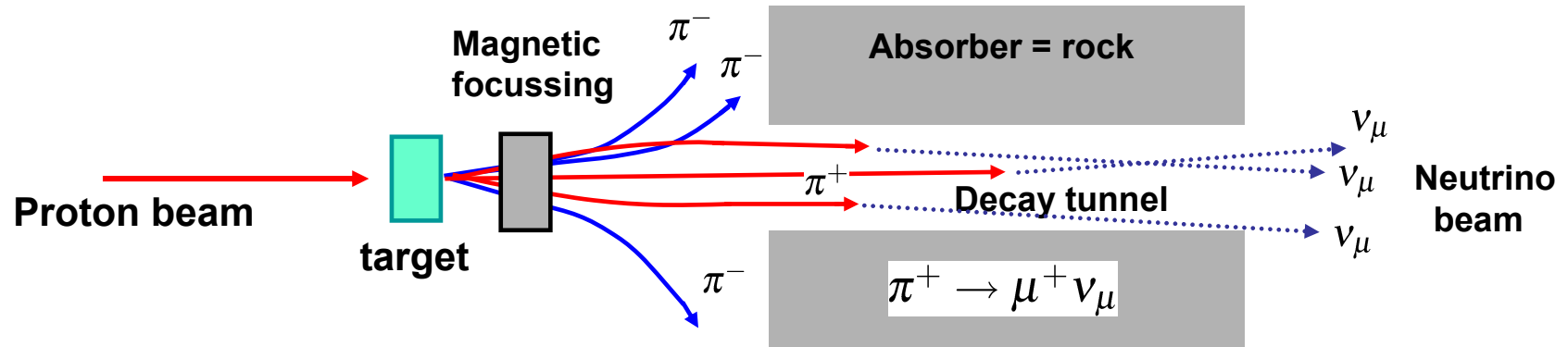




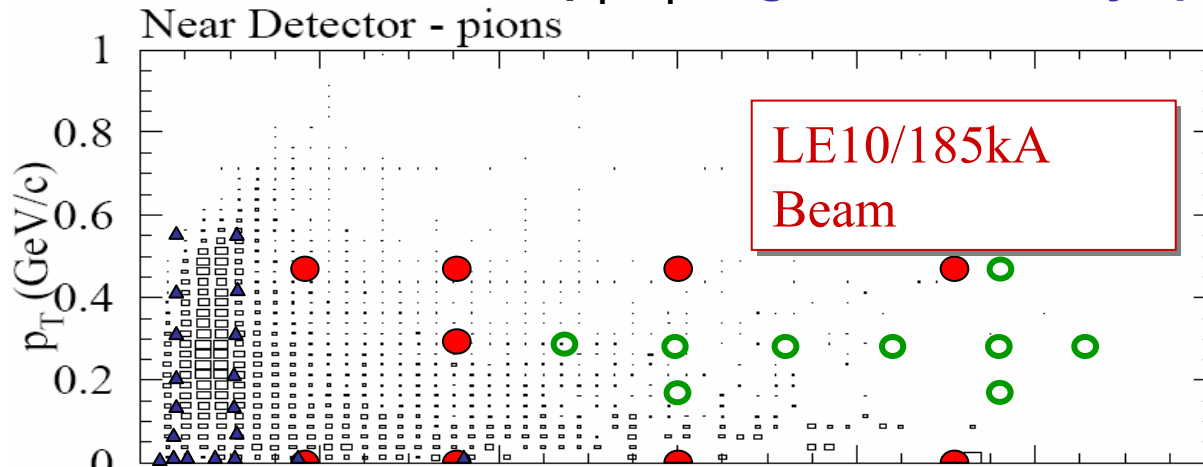
Hadron Production...



- ★ Perhaps the hardest part is predicting the neutrino flux
- ★ To do this need to know energy and p_T spectrum of meson from target



- ★ Hadron cascade models (e.g. GEANT, Fluka, MARS) all tuned to data
- ★ But data in relevant p_T, x_F region is relatively sparse



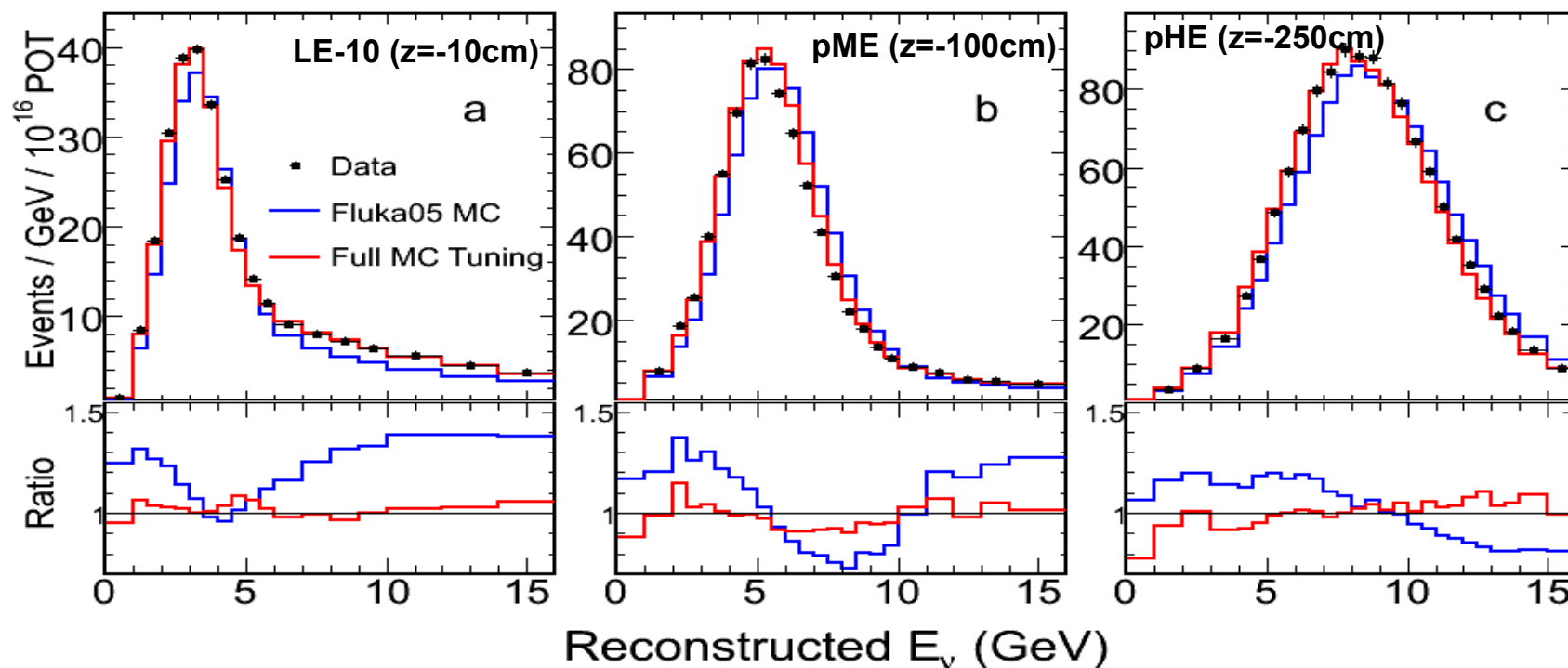
- ★ Situation will improve with data from MIPP experiment at Fermilab



The Near Detector to the Rescue



- ★ Want the expected Far Detector (FD) energy spectrum for selected CC events
- ★ Use the measured Near Detector (ND) energy spectrum
- ★ First “tune” Monte Carlo using ND data recorded in 7 different beam settings, e.g.



- Discrepancy between data and nominal (FLUKA05) MC changes with beam setting
- Suggestive that discrepancy is mainly due to flux rather than cross-section model
- Tune MC to ND data using a function that varies smoothly with hadronic x_F and p_T
- Tuned MC gives better agreement with data in all beam configurations

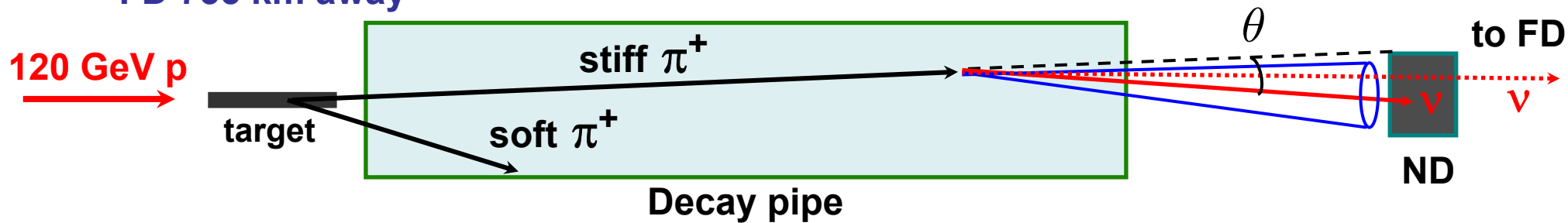


Extrapolating to the Far Detector

- ★ BUT: even in the absence of oscillations the NEAR and FAR detector neutrino spectra are different !

Easy to understand...

- ★ Consider a pion decaying in the decay pipe
- ★ Neutrino can intersect the ND for a relatively wide range of decay angles
- ★ For far detector only decays in a very small range of angles will cross the FD 735 km away



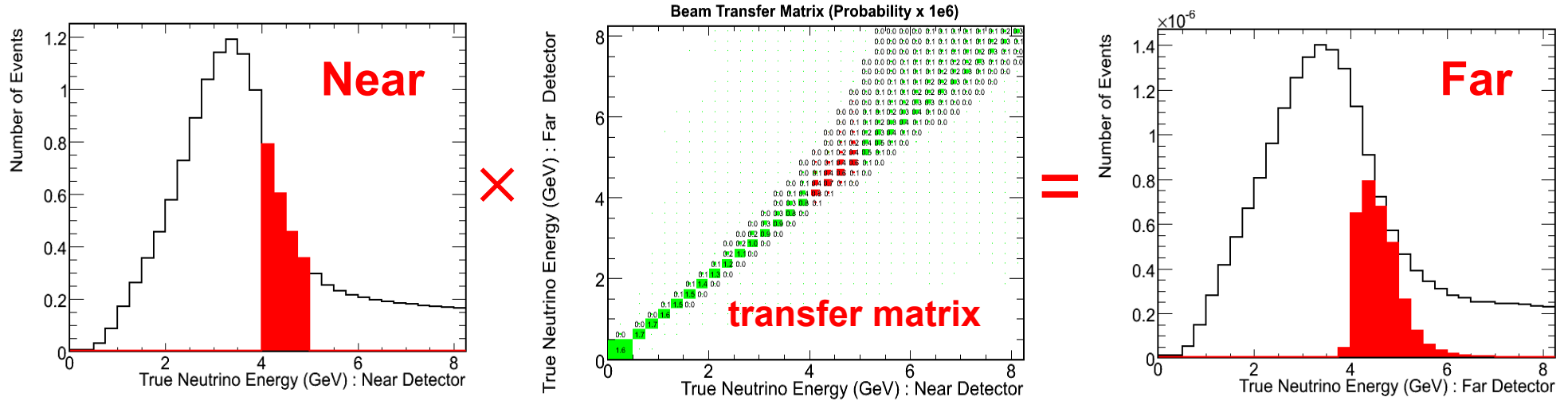
- ★ From simple relativistic kinematics for pion decay – neutrino energy depends on decay angle relative to pion line of flight

$$E_{\nu} = \frac{0.43E_{\pi}}{1 + \gamma^2\theta^2}$$

- ★ Decays with neutrinos pointing towards the FD tend to have smaller θ and hence have slightly higher energy
- ★ Difference is just kinematics, i.e. well understood !



The Beam Transfer Matrix

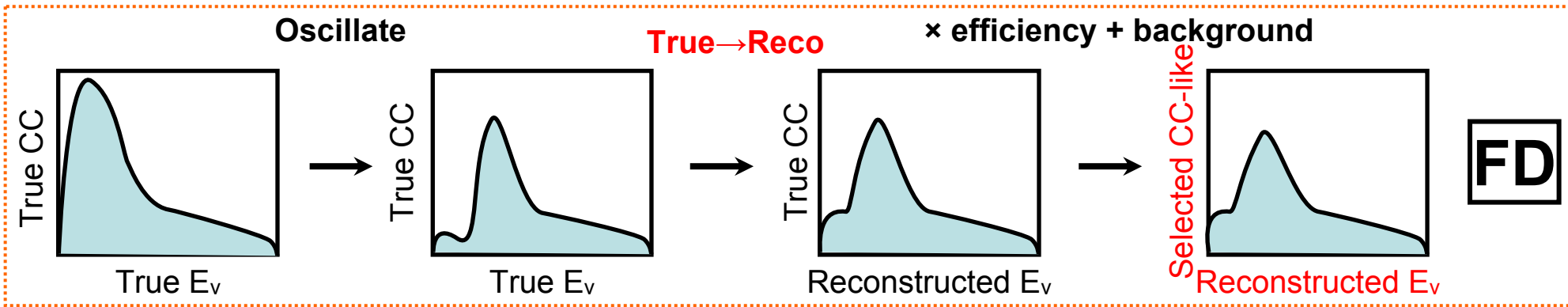
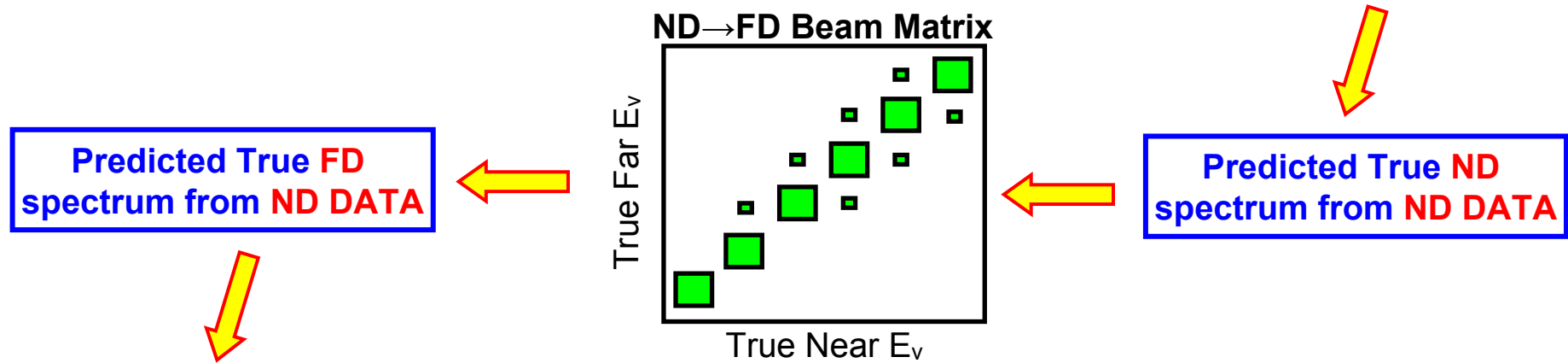
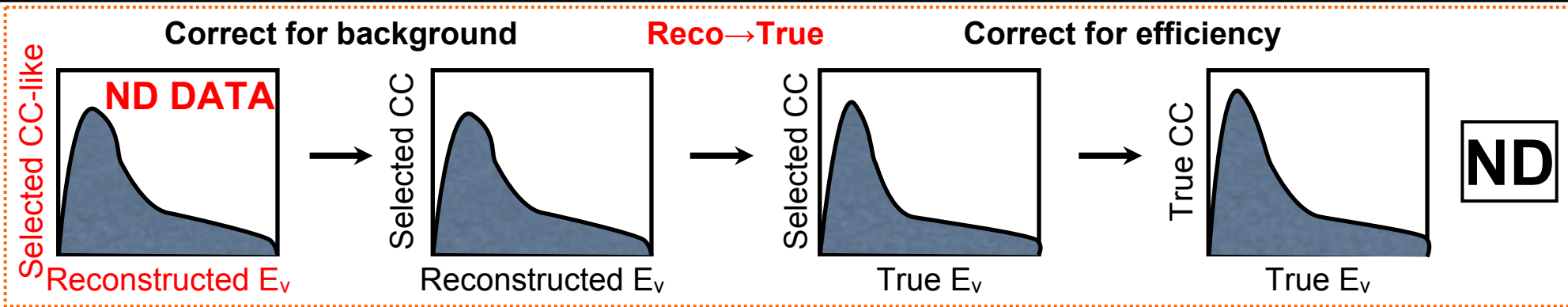


Beam Transfer Matrix:

- Encapsulates knowledge of 2-body pion decay and geometry
- Provides a simple way of relating near and far detector energy spectra
- Beam matrix determined from MC but does not depend strongly on details; kinematics & geometry dominate
- Near detector data **“directly”** determines predicted Far Detector spectrum



Details of matrix Near \rightarrow Far beam extrapolation

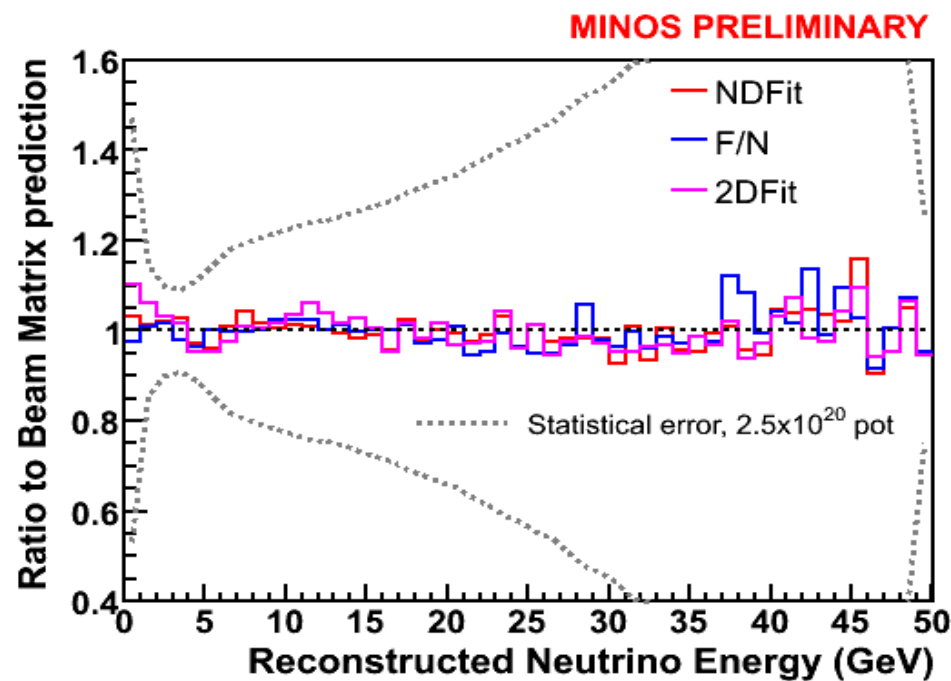
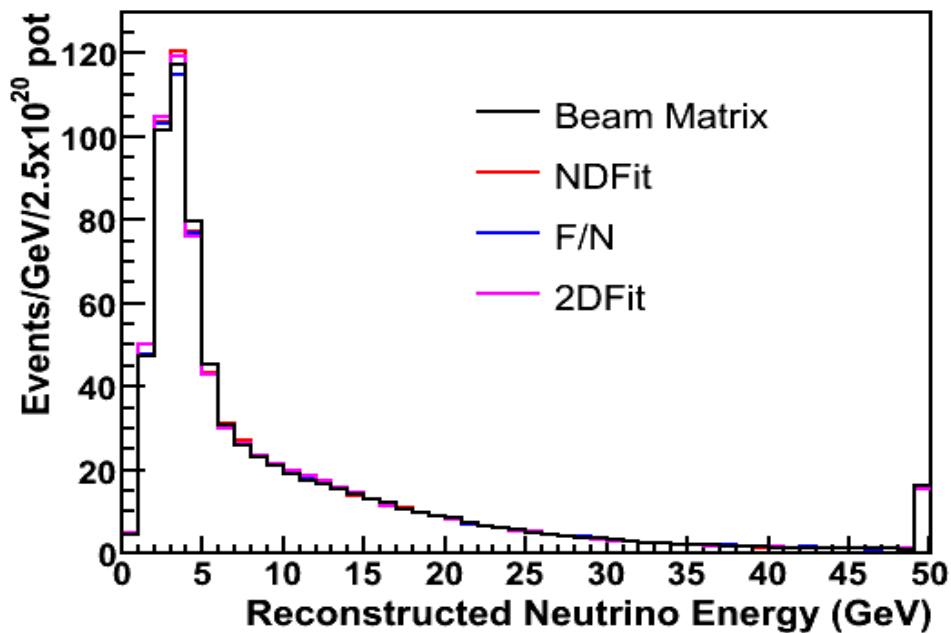




Cross-checks of the extrapolated spectrum



- ★ In addition to Beam Matrix Method have 3 cross-check methods to extrapolate ND energy spectrum:
 - Data-driven : **Far/Near ratio** “simple ID version of beam matrix”
 - Fit-based Methods : **NDFIT** and **2DFit**



- ★ Predicted Far Detector energy spectra agree with $\pm 4\%$
- ★ Much better than expected statistical error
- ★ Confident in far detector expectation...
- ★ LOOK AT FAR DETECTOR DATA (**blind analysis**)



Far Detector beam events



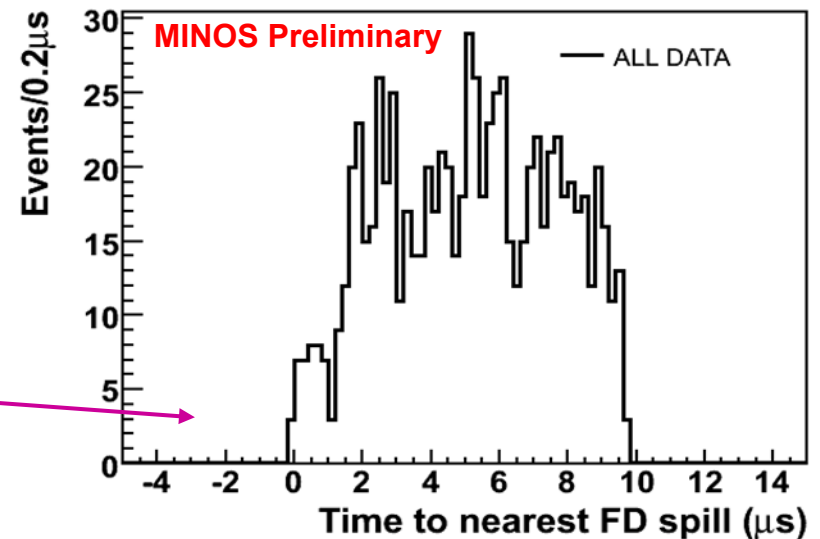
Cut	Number of Events
Track in fiducial volume	847
Data quality cuts	830
Timing cuts	828
Beam quality cuts	812
Track quality cuts	811
Track charge ≤ 0	672
CC PID parameter > 0.85	564
Reco $E_\nu < 200$ GeV	563 (Final sample)

Reject non-beam background, cosmic ray muons, etc.

Reject CC $\bar{\nu}_\mu$ Events

Reject NC Events

- ★ Very clean event sample
- ★ Non-beam background < 0.5 events (no candidates in $50\mu\text{s}$ window around spill time)

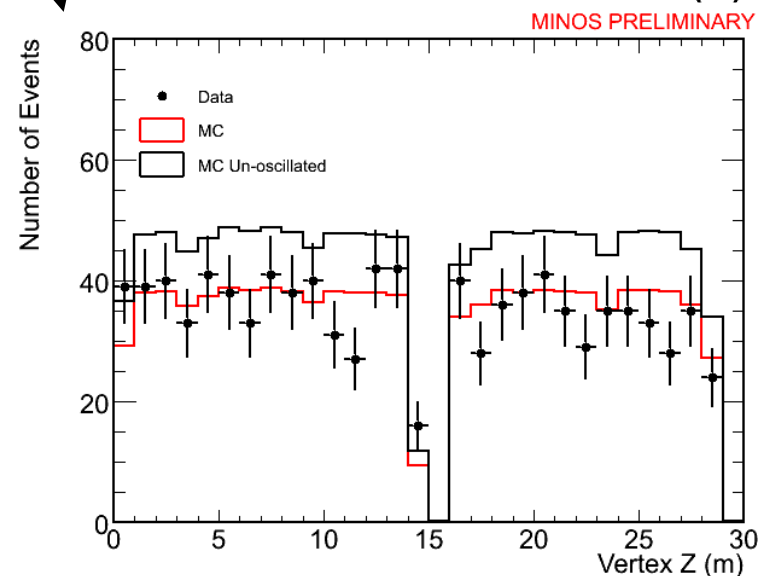
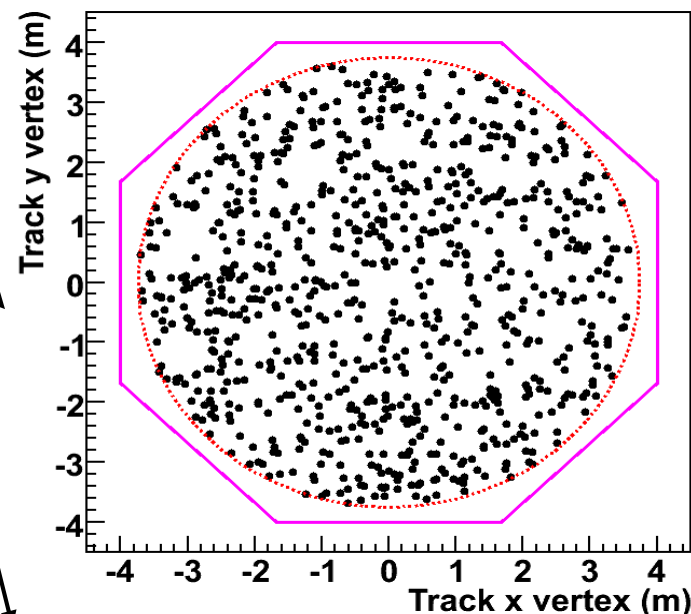
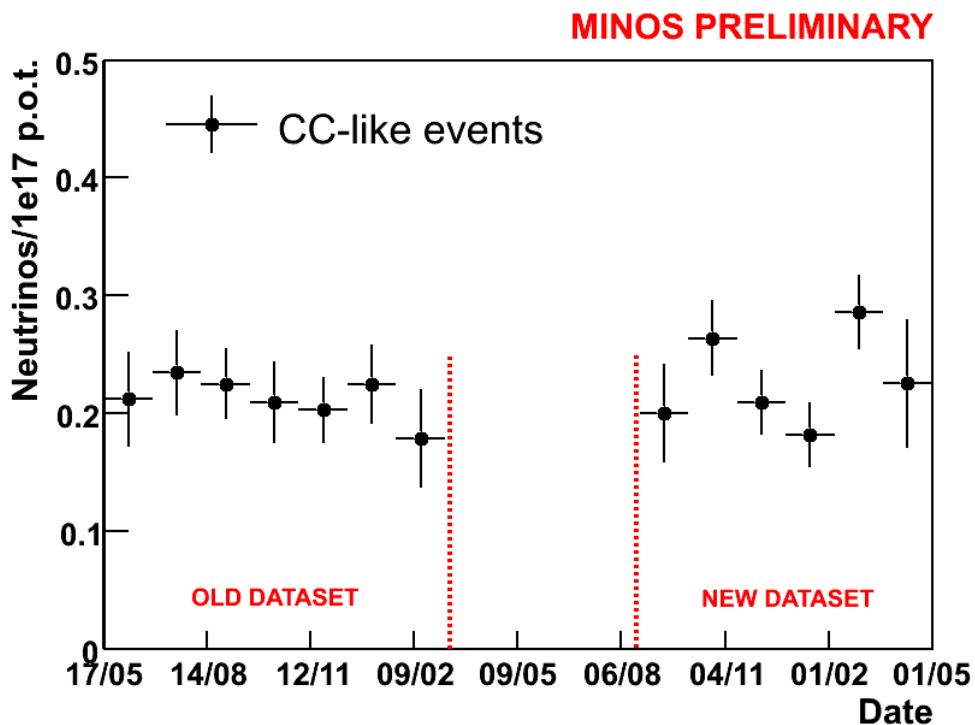




FD Events



- Far detector data well modelled by MC
- No indication of any unexpected background
- Events distributed uniformly in time and space

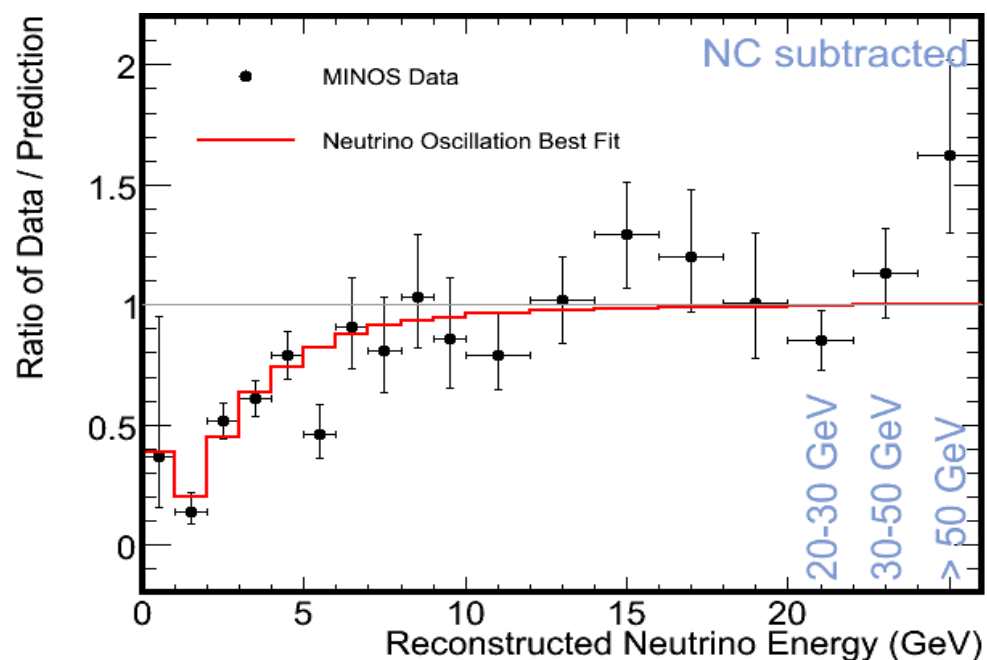
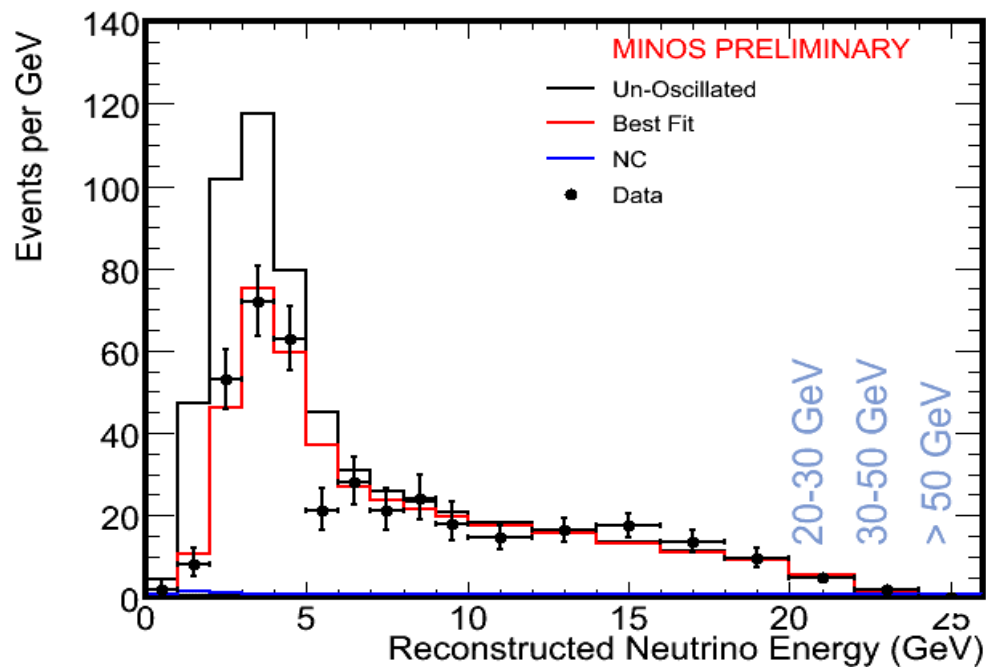




Oscillation Analysis



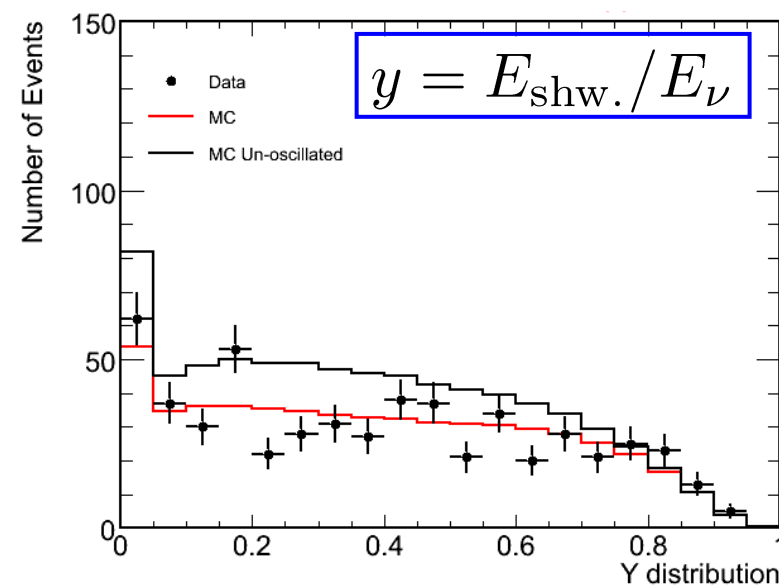
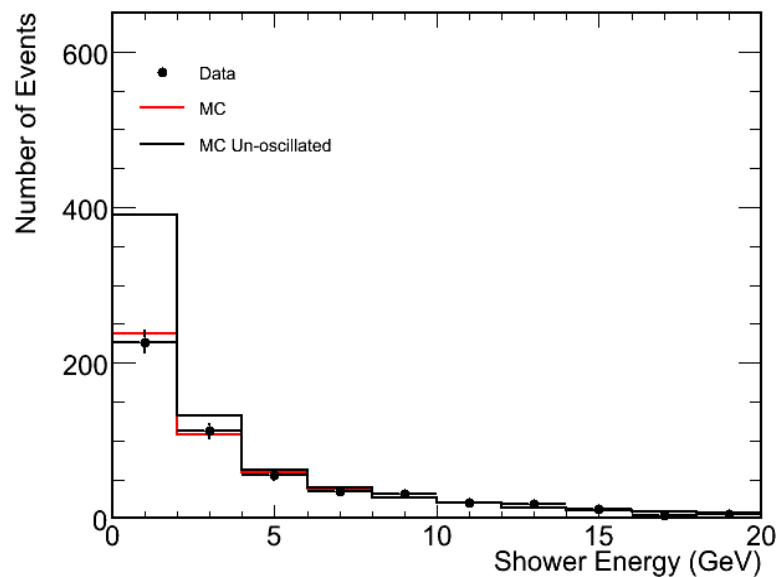
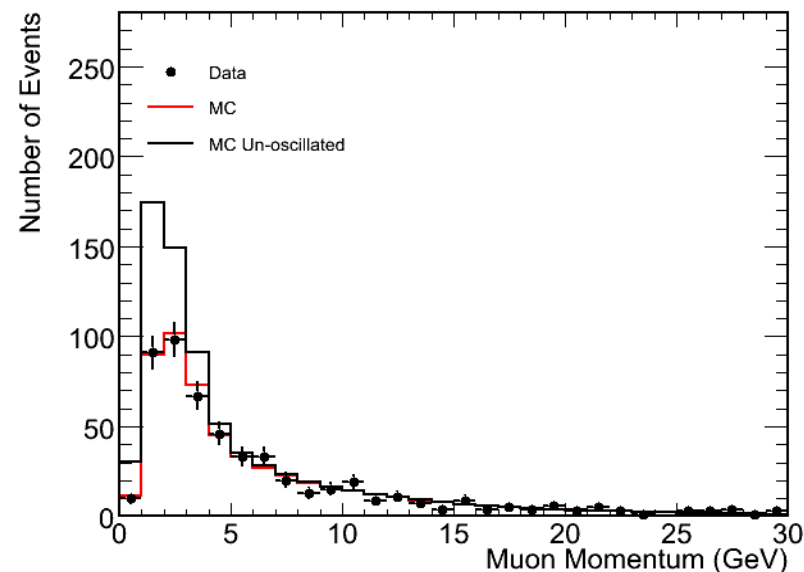
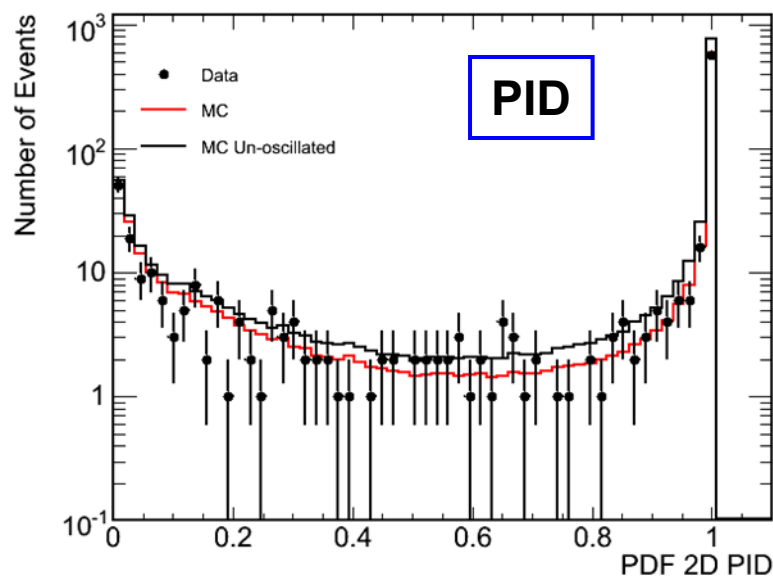
PRELIMINARY OSCILLATION RESULTS FOR 2.5×10^{20} POTs DATA.



Data sample	Observed	Expected (no osc.)	Observed / Expected
ν_μ (all E)	563	738 ± 30	0.74 (4.4σ)
ν_μ (<10 GeV)	310	496 ± 20	0.62 (6.2σ)
ν_μ (<5 GeV)	198	350 ± 14	0.57 (6.5σ)



Far detector distributions





Oscillation Fit/Systematic Uncertainties



- ◆ Oscillation parameters extracted from likelihood fit to reconstructed energy distribution of 563 selected Far Detector events

$$\chi^2(\Delta m^2, \sin^2 2\theta, \alpha_j, \dots) = \sum_{i=1}^{n_{bins}} \underbrace{2(e_i - o_i) + 2o_i \ln(o_i/e_i)}_{\text{statistical error}} + \underbrace{\sum_{j=1}^{n_{syst}} \frac{\Delta \alpha_j^2}{\sigma_{\alpha_j^2}}}_{\text{systematic errors}}$$

- ◆ The three largest uncertainties identified from this study are included as nuisance parameters in the oscillation analysis.

Uncertainty	Δm^2 (10^{-3} eV ²)	$\sin^2 2\theta$
Near/far normalization (4%)	0.065	<0.005
Abs. shower energy scale (10%)	0.075	<0.005
NC normalization (50%)	0.010	0.008
All other systematics	0.040	<0.005
Total uncertainty (quad. sum)	0.11	0.008
Statistical uncertainty	0.17	0.080

★ Currently statistical uncertainties dominate !



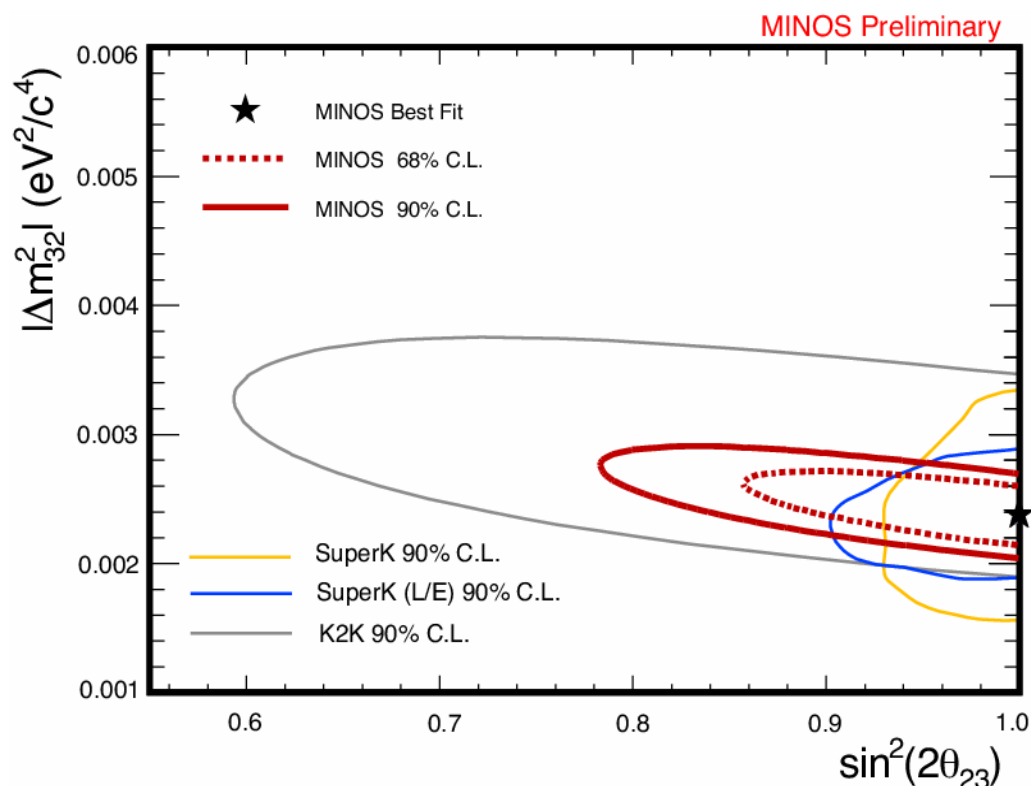
Results



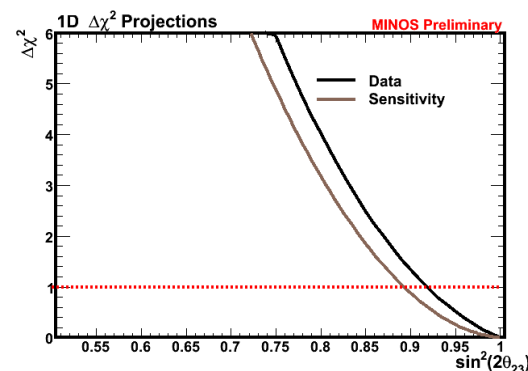
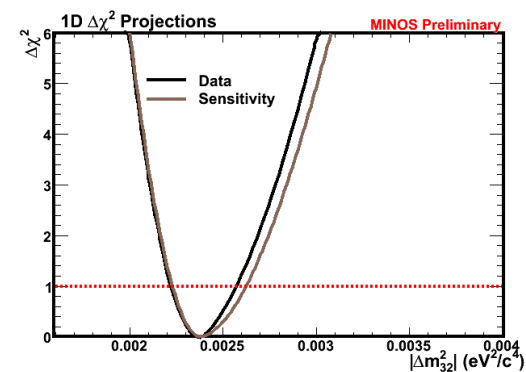
Best fit values:

$$|\Delta m_{32}^2| = 2.38_{-0.16}^{+0.20} \times 10^{-3} \text{ eV}^2$$
$$\sin^2 2\theta_{23} = 1.0 \quad (> 0.92 \text{ @68 \% C.L.})$$

$$\chi^2/n_{d.o.f} = 41.2/32$$



1D $\Delta\chi^2$ Projections



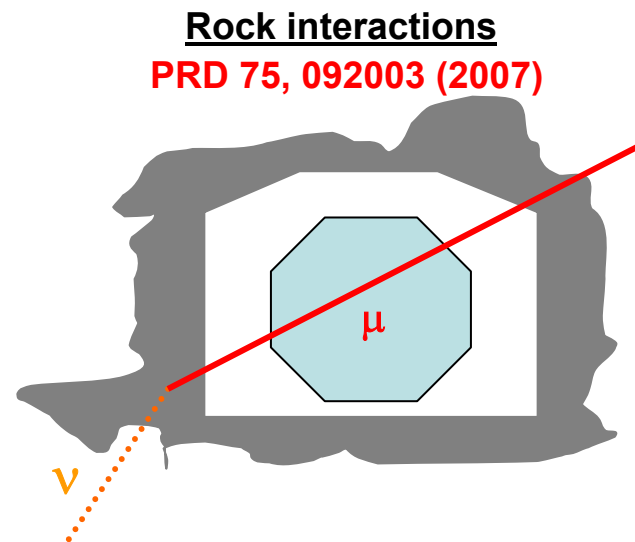
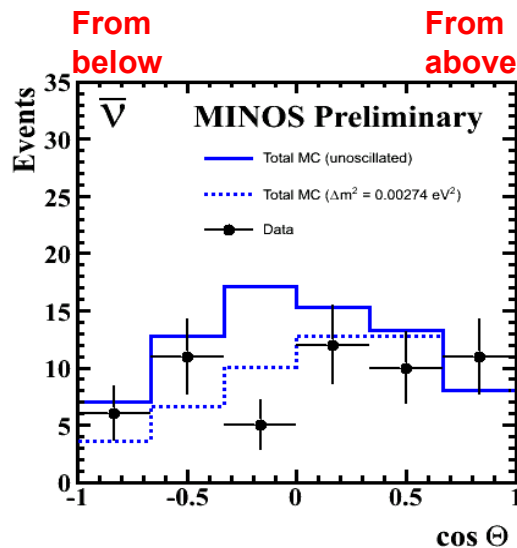
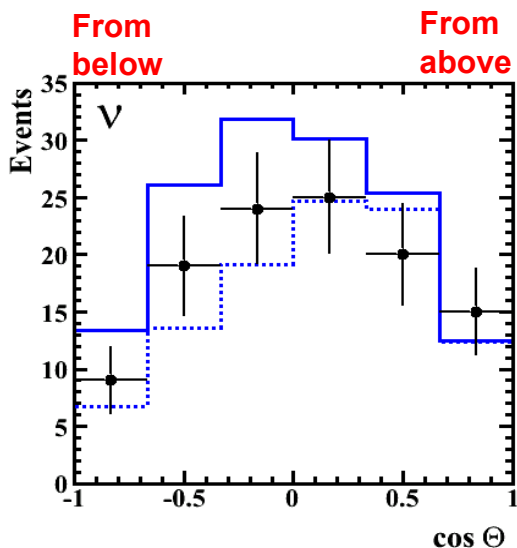
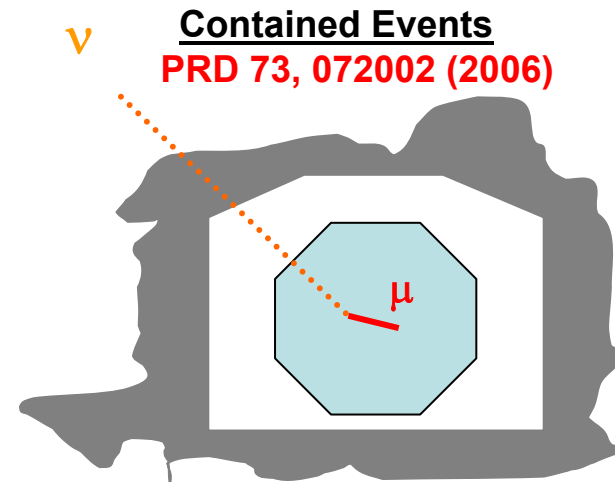


6 Other Results: Atmospheric Neutrinos



- ★ Event rate for 5.4 kton FD: 200 per year
- ★ 700m depth provides shielding from cosmic-ray μ
 - Event rate ~ 0.5 Hz
- ★ Magnetic field enables separation of ν_μ and $\bar{\nu}_\mu$
 - ◆ **MINOS is unique in this capability**
- ★ Start to test oscillations separately for $\nu_\mu/\bar{\nu}_\mu$
 - ◆ Test of CPT in neutrino sector
- ★ Currently cleanly identify 112 ν_μ and 55 $\bar{\nu}_\mu$

$$R_{\bar{\nu}/\nu}^{\text{data}} / R_{\bar{\nu}/\nu}^{\text{MC}} = 0.93_{-0.15}^{+0.19} \pm 0.12(\text{sys.})$$



▪ Higher statistics results out later this year.

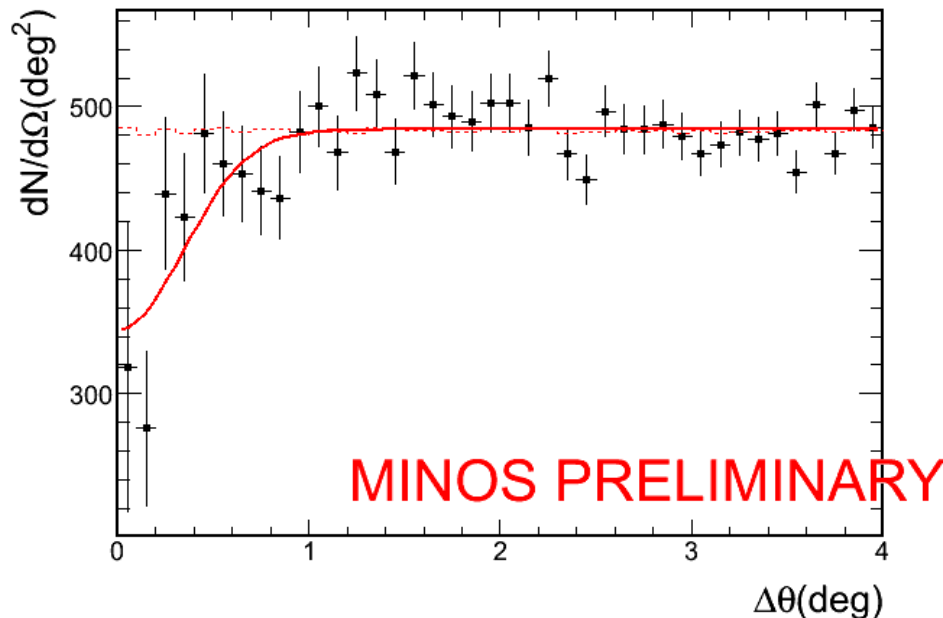
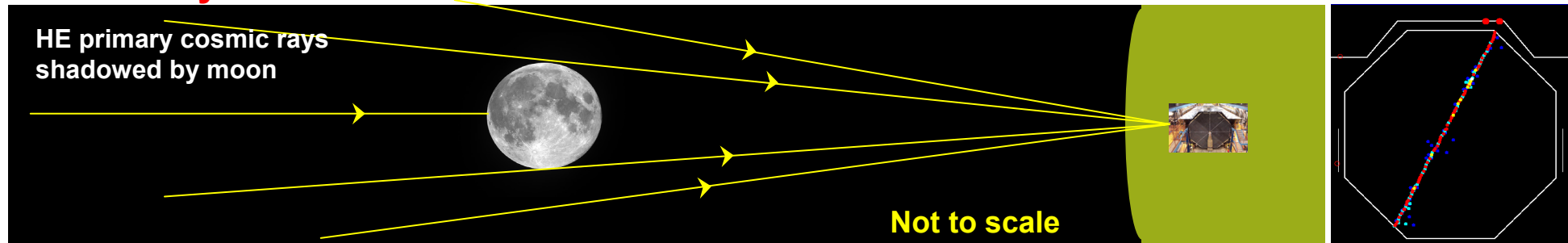


Other Results: Cosmic-Ray Physics



- ★ MINOS is a large deep underground detector and can make a number of interesting cosmic-ray measurements, e.g.

Cosmic-ray Moon shadow



Demonstrates:

- ◆ that the moon exists
- ◆ we understand our reconstruction

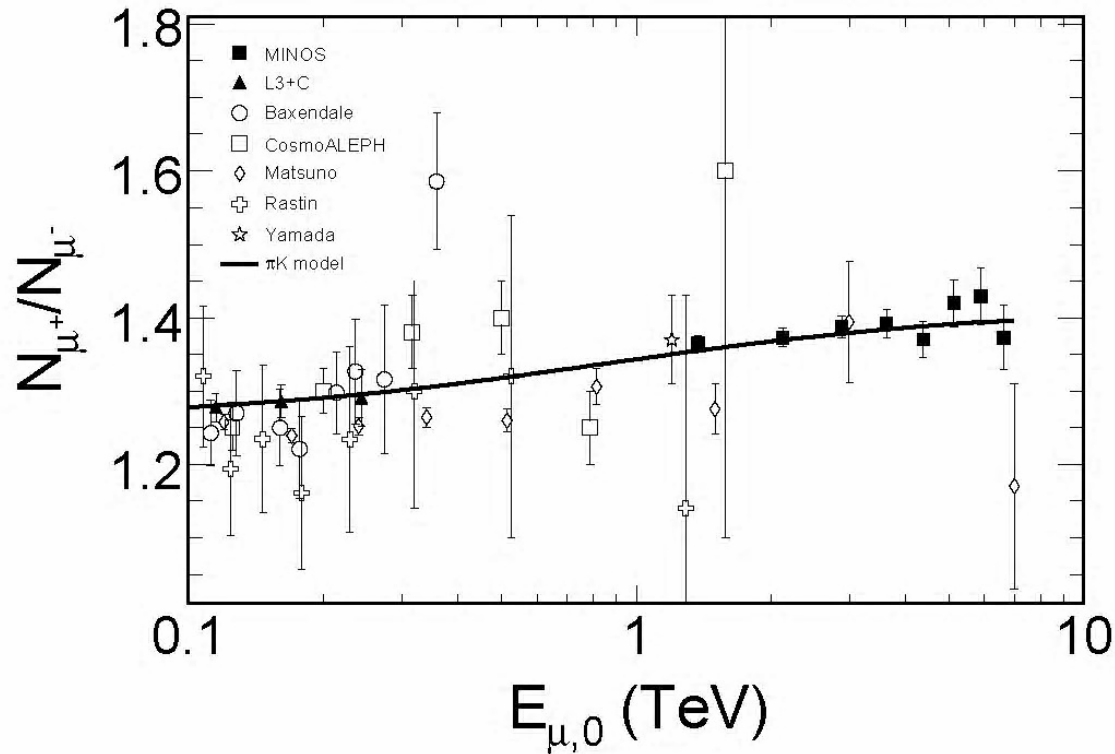
Looking for sun-shadow is more interesting as it probes solar magnetic field (ongoing work)



Cosmic-Ray Physics cont.



- ★ Can also take advantage of Magnetic field and measure the cosmic-ray muon charge ratio



- ★ MINOS data shows rise in ratio – sensitive to kaon fraction in cosmic-ray induced air showers



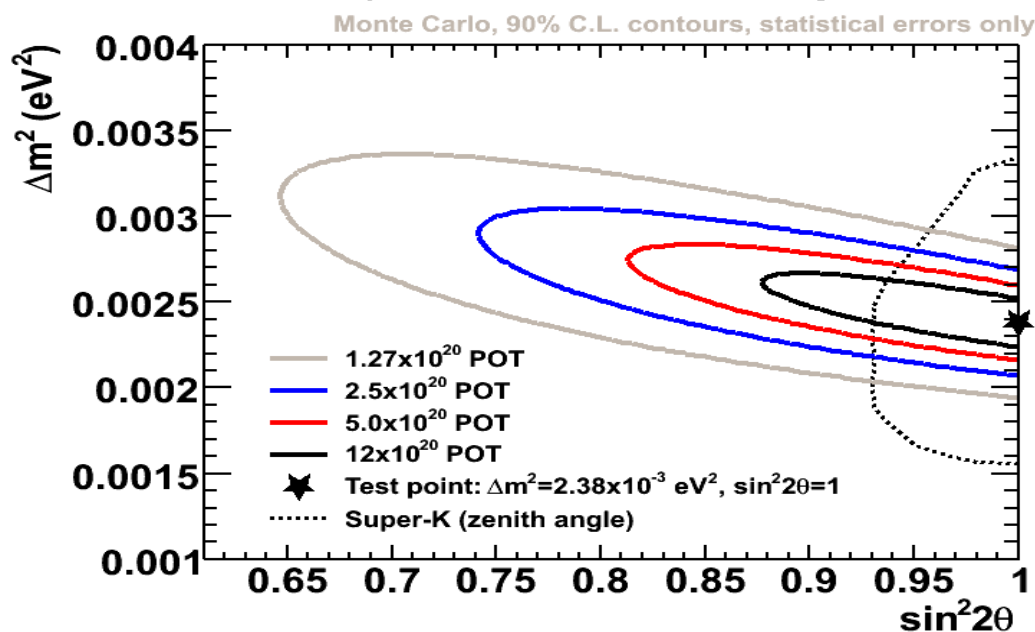
7 Future Prospects



CC Disappearance:

- ★ Currently expected to accumulate 12×10^{20} POT of data by end of 2009
- ★ Significant improvements expected

MINOS Sensitivity as a function of Integrated POT



- + significant improvements in analysis already in hand
- ★ hope to have sensitivity to $\sin^2 2\theta_{23}$ which is competitive with Super-K
will depend on success of analysis improvements



Alternative Scenarios

- ★ MINOS is the first high statistics long-baseline experiment
- ★ Can study shape of oscillation curve in detail
- ★ Compare standard oscillation hypothesis to other scenarios, e.g.

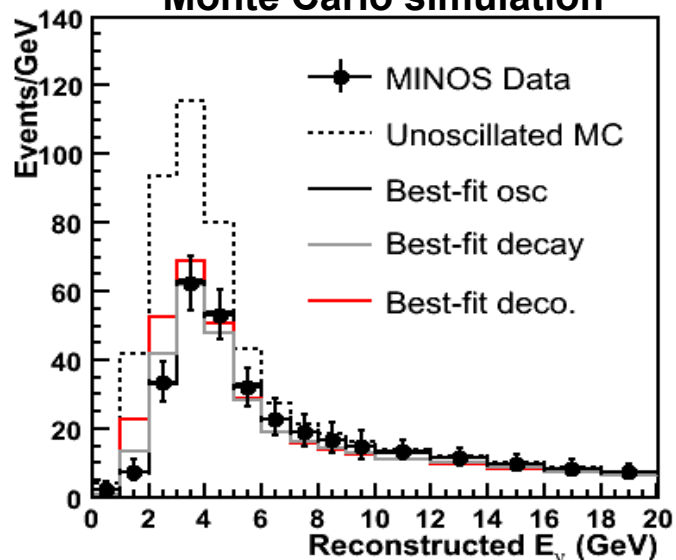
Neutrino Decay

$$P(\nu_\mu \rightarrow \nu_\mu) = (\sin^2 \theta + \cos^2 \theta e^{-\frac{\alpha L}{2E}})^2$$

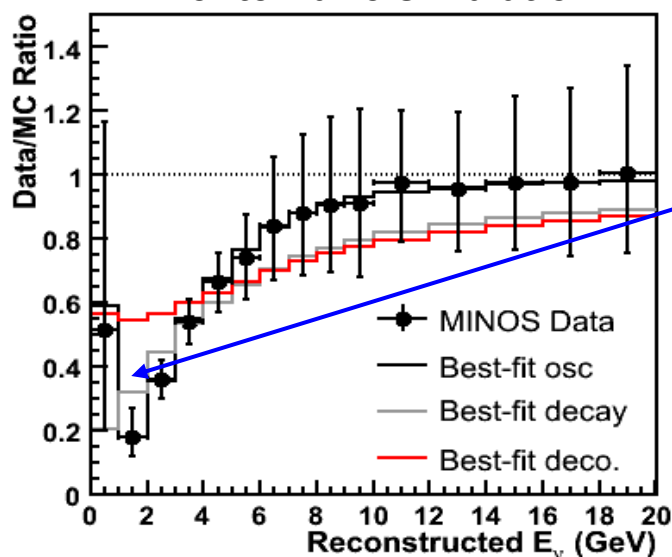
Neutrino Decoherence

$$P(\nu_\mu \rightarrow \nu_\mu) = 1 - \frac{\sin^2 2\theta}{2} \left(1 - e^{-\frac{\mu^2 L}{2E}} \right)$$

Monte Carlo simulation



Monte Carlo simulation



★ First results due this Summer...



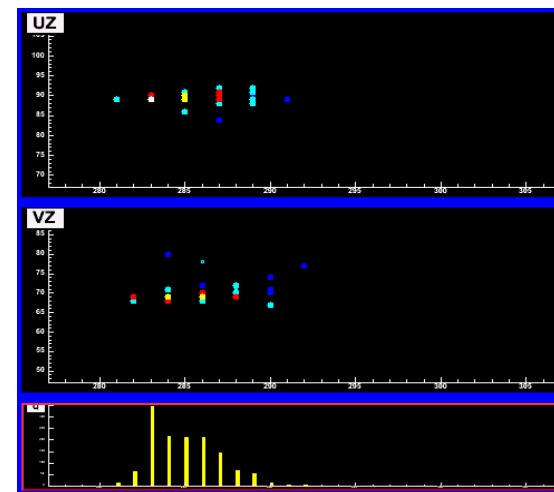
Electron Neutrino Appearance

- ★ Search for $\nu_\mu \rightarrow \nu_e$ oscillations is a hot-topic in neutrino physics
- ★ The next generation of neutrino experiments (T2K, Double-Chooz, Nona) all designed to search for $\nu_\mu \rightarrow \nu_e$ oscillations and measure θ_{13}
- ★ Vital for longer term projects to probe **CP** violation in the neutrino sector as **CP** violating terms in **PMNS** matrix enter multiplied by $\sin \theta_{13}$

$$\begin{pmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu 1} & U_{\mu 2} & U_{\mu 3} \\ U_{\tau 1} & U_{\tau 2} & U_{\tau 3} \end{pmatrix} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix} \times \begin{pmatrix} c_{13} & 0 & s_{13}e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{-i\delta} & 0 & c_{13} \end{pmatrix} \times \begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

- ★ This is a very challenging analysis in **MINOS**

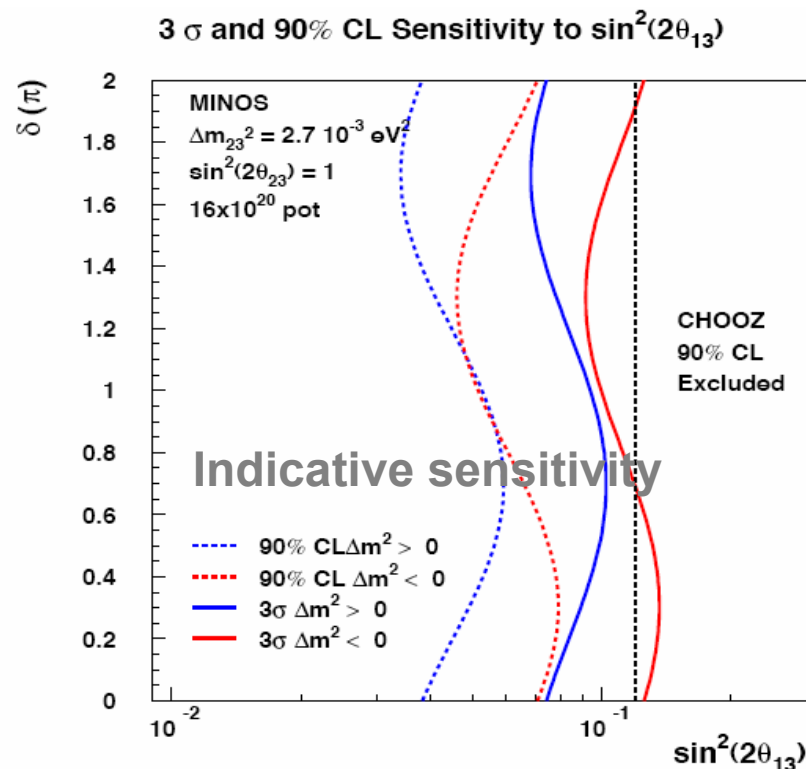
- course sampling
- events have relatively few “hits”
- event rate low <20 events in current data
- large background from NC interactions:
 π^0 in hadronic shower \Rightarrow EM shower





- ★ Sophisticated analyses being designed to efficiently separate signal from background
- ★ For example, **Monte Carlo Nearest Neighbour (MCNN)** method:
 - rather than perform multivariate analysis on reconstructed quantities
 - directly compare patterns of hits in event to large MC libraries of NC and ν_e events (~50 million)
 - identify best matches and for discriminant variable from fraction of N best MC matches that were ν_e

- Expect first results before Summer
- MINOS has the possibility to discover $\nu_\mu \rightarrow \nu_e$ oscillations
- Will need entire MINOS data set
- However, MINOS is ahead of the game; will publish before Double-Chooz/T2K





Future Prospects, cont.



In addition,...

- ★ Updated atmospheric neutrino results on anti-neutrinos
- ★ Cosmic-ray measurements
- ★ Structure function/cross section measurements - in progress.

★ Possibility of anti-neutrino running to provide first precise measurement of $\bar{\nu}_\mu \rightarrow \bar{\nu}_\tau$: 10 % precision achievable with ~6 months running

Is this sufficiently interesting ?



8 Summary



Summary

- ★ MINOS/Numi running since mid-2005
- ★ Already accumulated a large data sample, 3.2×10^{20} POT
- ★ Data analysis in advanced stage
 - Good understanding of beam – 2 detectors vital !
 - First results on beam data published (PRL)
 - MINOS already has most precise measurement of $|\Delta m_{32}^2|$
- ★ Many other analyses reaching maturity
 - Search for sub-dominant $\nu_{\mu} \rightarrow \nu_e$ oscillations is perhaps the most exciting

Outlook

- ★ MINOS remains a very high priority for Fermilab
- ★ Expect to run through US FY2010, i.e. until October 2010
- ★ Final data sample of $> 1.2 \times 10^{21}$ POT
- ★ With these data:
 - 5 % measurement of $|\Delta m_{32}^2|$
 - And maybe the first observation of $\nu_{\mu} \rightarrow \nu_e$ oscillations
 - + much more



The End



Thank you