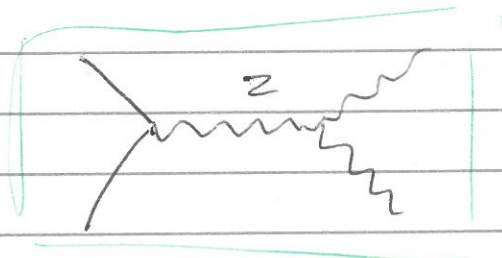
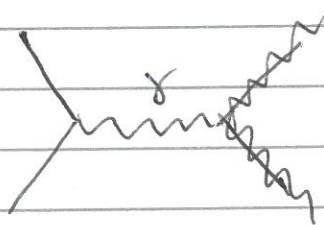
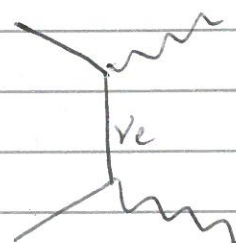


Section 10

Slide 3



Both have massless propagators

$\therefore \sigma \propto E \Rightarrow$  unitarity violation!

\* Add in Z diagram \*

Individual contributions do diverge  
BUT now they cancel each other

$\sigma_{TOTAL}$  well behaved at high  $\sqrt{s}$

Can only cancel if couplings are related

slide 4/5

$SU(2)$  acts on weak isospin  $I$

$U(1)$  acts on weak hypercharge  $Y = 2(Q - I_3)$

$\uparrow$   
third component  
of isospin

Act on LH doublets  $\begin{pmatrix} \nu_e \\ e^- \end{pmatrix}_L \quad \begin{pmatrix} u \\ d \end{pmatrix}_L$

but not on  $\nu_{eR} \quad e^-_R \quad u_R \quad d_R$   
RH singlets

	LH			RH		
	Q	$I_3$	Y	Q	$I_3$	Y
$\nu_e, \nu_\mu, \nu_\tau$	0	$+\frac{1}{2}$	-1	0	0	0
$e^-, \mu^-, \tau^-$	-1	$-\frac{1}{2}$	-1	-1	0	-2
u c t	$+\frac{2}{3}$	$+\frac{1}{2}$	$+\frac{1}{3}$	$+\frac{2}{3}$	0	$+\frac{4}{3}$
d s b	$-\frac{1}{3}$	$-\frac{1}{2}$	$+\frac{1}{3}$	$-\frac{1}{3}$	0	$-\frac{2}{3}$

$\uparrow$   
no weak interaction

Imposing  $SU(2) \times U(1)$  symmetry

$\Rightarrow$  4 massless gauge bosons

$W^+ \quad W^- \quad W_3 \quad B$

$\underbrace{\hspace{2cm}}_{\text{mix}} \Rightarrow \gamma, Z$

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only if true

$$\langle \psi | \hat{M}^2 | \psi \rangle = m^2$$

$$m_Z^2 = \langle Z | \hat{M}^2 | Z \rangle = \langle W_3 \cos \theta - B \sin \theta | \hat{M}^2 | W_3 \cos \theta - B \sin \theta \rangle \\ = m_W^2 \cos^2 \theta + m_B^2 \sin^2 \theta - 2m_W m_B \cos \theta \sin \theta$$

$$m_A^2 = \langle A | \hat{M}^2 | A \rangle = \langle W_3 \sin \theta + B \cos \theta | \hat{M}^2 | W_3 \sin \theta + B \cos \theta \rangle \\ = m_W^2 \sin^2 \theta + m_B^2 \cos^2 \theta + 2m_W m_B \cos \theta \sin \theta$$

$$m_Z^2 + m_A^2 = m_W^2 + m_B^2 \quad \text{if } A \equiv \gamma, \text{ then } m_A = 0$$

$$\Rightarrow m_Z^2 = m_W^2 + m_B^2$$

find  $m_B^2$  from  $m_A^2 = 0$

find  $m_W m_B$  from unphysical state  $\langle Z | \hat{M}^2 | A \rangle \equiv 0$

$$m_{ZA}^2 = (m_W^2 - m_B^2) \cos \theta \sin \theta + m_W m_B (\cos^2 \theta - \sin^2 \theta) \equiv 0 \\ \rightarrow m_W m_B = \frac{(m_B^2 - m_W^2) \cos \theta \sin \theta}{\cos^2 \theta - \sin^2 \theta}$$

$$m_A^2 = m_W^2 \sin^2 \theta + m_B^2 \cos^2 \theta + \frac{2(m_B^2 - m_W^2) \cos^2 \theta \sin^2 \theta}{\cos^2 \theta - \sin^2 \theta} \equiv 0 \\ \rightarrow m_B^2 = m_W^2 \frac{\sin^2 \theta}{\cos^2 \theta}$$

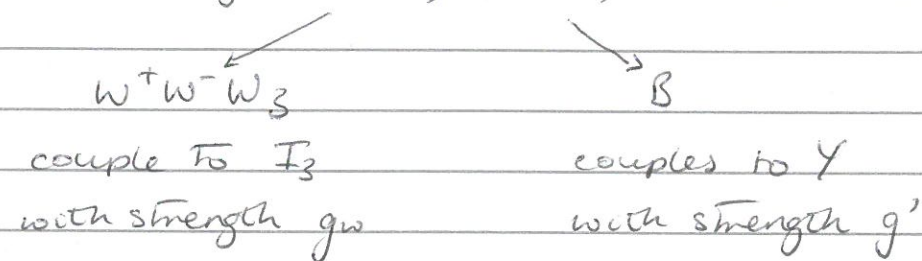
$$\Rightarrow m_Z^2 = m_W^2 \frac{\sin^2 \theta}{\cos^2 \theta} + m_W^2 = m_W^2 \left( 1 + \frac{\sin^2 \theta}{\cos^2 \theta} \right)$$

$$= \frac{m_W^2}{\cos^2 \theta}$$

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only if true

EWK unification using  $su(3) \times u(1)$



Photon coupling  $A = W_3 \sin \theta + B \cos \theta$

$$Q_e = g_W I_3 \sin \theta + g' Y \cos \theta = g_W I_3 \sin \theta + g' (Q - I_3) \cos \theta$$

$$\text{for } Q=0 \quad 0 = g_W I_3 \sin \theta - 2g' I_3 \cos \theta \quad g' = \frac{g_W}{2} \tan \theta$$

$$\text{for } I_3=0 \quad Q_e = 2g' Q \cos \theta \quad g' = \frac{e}{2 \cos \theta}$$

$$\Rightarrow g_W = \frac{e}{\sin \theta}$$

Z coupling  $Z = W_3 \cos \theta - B \sin \theta$

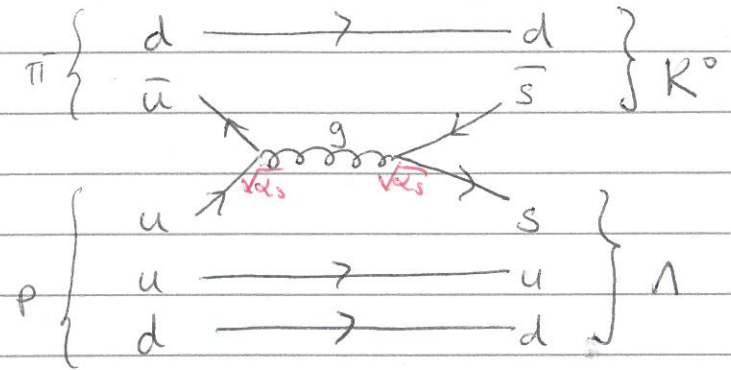
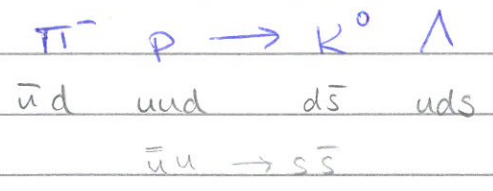
$$g_W I_3 \cos \theta - g' 2(Q - I_3) \sin \theta \quad \text{Subs } g_W = \frac{e}{\sin \theta} \quad g' = \frac{e}{2 \cos \theta}$$

$$= \frac{e I_3 \cos \theta}{\sin \theta} - \frac{e}{\cos \theta} (Q - I_3) \sin \theta$$

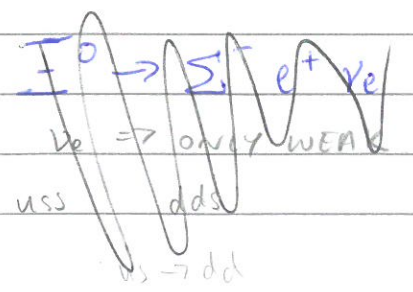
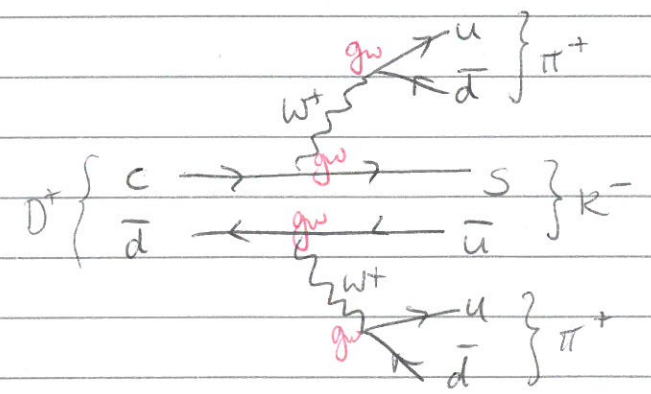
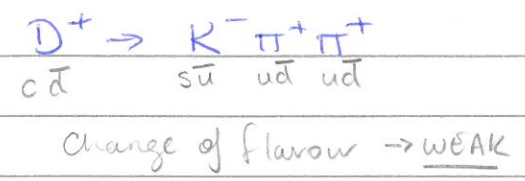
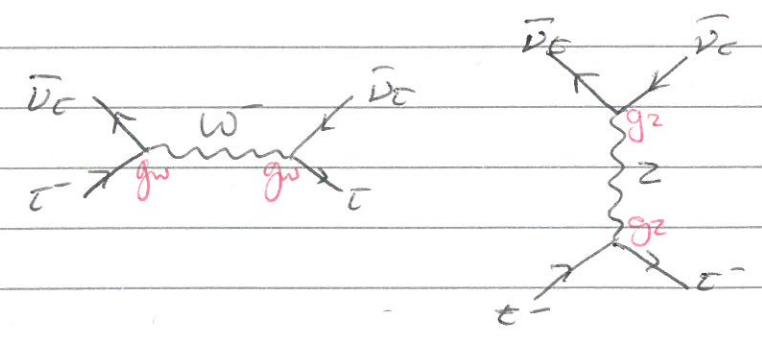
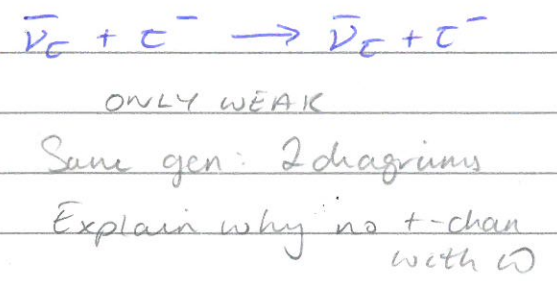
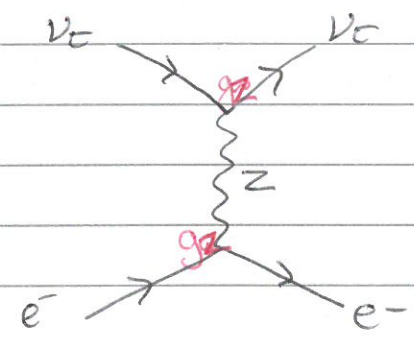
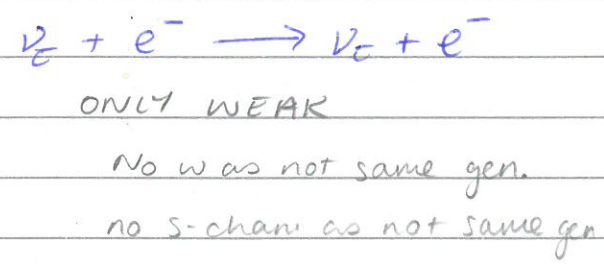
$$= \frac{e}{\sin \theta \cos \theta} \left[ I_3 - Q \sin^2 \theta \right]$$

$g_Z$  couples to  $I_3$  &  $Y$  with strength  $g_Z$

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Strong, EM or Weak



slide 20/21/22

$$\sigma(e^+e^- \rightarrow Z \rightarrow f\bar{f}) = \frac{g^2 \pi}{E_e^2} \frac{\Gamma_{ee} \Gamma_{ff}}{(\bar{E}_{cm} - m_Z)^2 + \Gamma_Z^2/4}$$

$g$ : spin factor: make spin 2 from 2x spin 1/2 electrons  
 $g = \frac{2J_z + 1}{(2J_e + 1)(2J_e + 1)} = \frac{3}{4}$

$E_e^2$  collide  $\bar{E}_{cm}^2 = (\bar{E}_1 + \bar{E}_2)^2 - (\vec{p}_1 + \vec{p}_2)^2$  head on collisions  $\bar{E}_1 = \bar{E}_2$   
 $= 4E_e^2$   
 $\Rightarrow \bar{E}_e^2 = \bar{E}_{cm}^2 / 4$

$$\sigma = \frac{3\pi}{\bar{E}_{cm}^2} \frac{\Gamma_{ee} \Gamma_{ff}}{(\bar{E}_{cm} - M_Z)^2 + \Gamma_Z^2/4}$$

at  $\bar{E}_{cm} = M_Z$   $\sigma = \frac{3\pi}{M_Z^2} \frac{\Gamma_{ee} \Gamma_{ff}}{\Gamma_Z^2/4}$

$\Gamma_Z$  is the total width =  $\Gamma_{ee} + \Gamma_{\mu\mu} + \Gamma_{\tau\tau} + \Gamma_{qq} + \Gamma_{\nu\nu}$

$\therefore \sigma(e^+e^- \rightarrow Z \rightarrow \text{anything}) = \frac{12\pi}{M_Z^2} \frac{\Gamma_{ee}}{\Gamma_Z}$   $\Gamma_{ff} = \Gamma_Z$

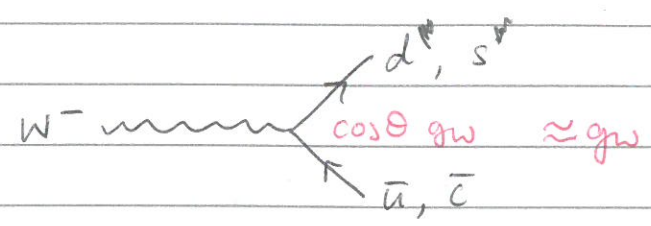
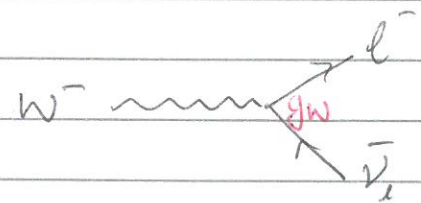
compare to  $\sigma(e^+e^- \rightarrow \gamma \rightarrow \text{anything}) = \frac{4\pi\alpha^2}{3\bar{E}_{cm}^2} = \frac{4\pi\alpha^2}{3M_Z^2}$   $\bar{E}_{cm} = M_Z$

$\frac{\Gamma_Z}{\sigma_{\text{QED}}} \sim 5700$

$\Gamma_{ee} = 85 \text{ keV}$   
 $\Gamma_Z = 2.5 \text{ GeV}$   
 $\alpha = 1/137$

Z diagram dominates at  $\sim m_Z = \sqrt{s}$

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Assume  $\cos\theta = 1$  : only Cabibbo allowed decays.

Expect  $W^- \rightarrow e^- \bar{\nu}_e$   
 $\mu^- \bar{\nu}_\mu$   
 $\tau^- \bar{\nu}_\tau$   
 $\bar{u} d \times 3$  for color  
 $\bar{c} s \times 3$  "

} Should have ~ same coupling

Cabibbo suppressed final states also contribute in a small way

$$W^- \rightarrow l \nu = \frac{3}{9} = \frac{1}{3} \qquad W \rightarrow q \bar{q} = \frac{6}{9} = \frac{2}{3}$$

$$\begin{aligned}
 W^+ W^- &\rightarrow l \nu \quad \frac{1}{3} \times \frac{1}{3} = \frac{1}{9} \\
 &\rightarrow l \nu q \bar{q} \quad \frac{1}{3} \times \frac{2}{3} \times 2 = \frac{4}{9} \\
 &\rightarrow q \bar{q} q \bar{q} \quad \frac{2}{3} \times \frac{2}{3} = \frac{4}{9}
 \end{aligned}$$