The contribution of diagram (a) is given by

$$M = \frac{(-ig_W)^2}{8} \bar{v}(p_+) \not\in (q_+)(1-\gamma_5) \frac{i}{\not\!p_- - \not\!q_-} \not\in (q_-)(1-\gamma_5)u(p_-) .$$
(8.114)

Inserting the leading term for the polarizations from Eq. (8.113) and dropping non-leading terms, we obtain

$$M = -i \frac{(-ig_W)^2}{8M_W^2} \bar{v}(p_+)(q_+ - q_-)(1 - \gamma_5)u(p_-) . \qquad (8.115)$$

The contribution of the photon exchange diagram (b), using the Feynman rules of Figs. 8.2(updated) and 8.5, is

$$M = (-ig_W)^2 Q_e \sin^2 \theta_W \bar{v}(p_+) \gamma^{\rho} u(p_-) \frac{-ig_{\rho\alpha}}{(q_+ + q_-)^2} \times V^{\beta\delta\alpha}(-q_+, -q_-, q_+ + q_-) \varepsilon_{\beta}(q_+) \varepsilon_{\delta}(q_-) .$$
(8.116)

The contribution of the Z exchange diagram (c) is

$$M = \frac{(-ig_W)^2}{2} \bar{v}(p_+) \gamma^{\rho} (V_e - A_e \gamma_5) u(p_-) \frac{-ig_{\rho\alpha}}{(q_+ + q_-)^2 - M_Z^2} \times V^{\beta\delta\alpha} (-q_+, -q_-, q_+ + q_-) \varepsilon_{\beta}(q_+) \varepsilon_{\delta}(q_-) .$$
(8.117)

where

$$V^{\beta\delta\alpha}(p,q,r) = g^{\beta\delta}(p^{\alpha} - q^{\alpha}) + g^{\delta\alpha}(q^{\beta} - r^{\beta}) + g^{\alpha\beta}(r^{\delta} - p^{\delta}) .$$
(8.118)

Hence we find, using the approximate longitudinal polarization vector of Eq. (8.113),

$$V^{\beta\delta\alpha}(-q_{+}, -q_{-}, q_{+} + q_{-})\varepsilon_{\beta}(q_{+})\varepsilon_{\delta}(q_{-}) = \frac{(q_{+} + q_{-})^{2}}{2M_{W}^{2}} \Big[q_{+}^{\alpha} - q_{-}^{\alpha}\Big] + O(1) .$$
(8.119)

The result for the sum of the Z and γ exchange contributions, after inserting the values for V_e and A_e , is

$$M = i \frac{(-ig_W)^2}{8M_W^2} \bar{v}(p_+)(q_+ - q_-)(1 - \gamma_5)u(p_-) , \qquad (8.120)$$

which cancels exactly with Eq. (7.115). So we see that the leading high-energy behaviour is cancelled, as a result of the relationship between the three-boson couplings and the couplings to fermions imposed by the gauge structure.