

The new MC event generator Herwig++

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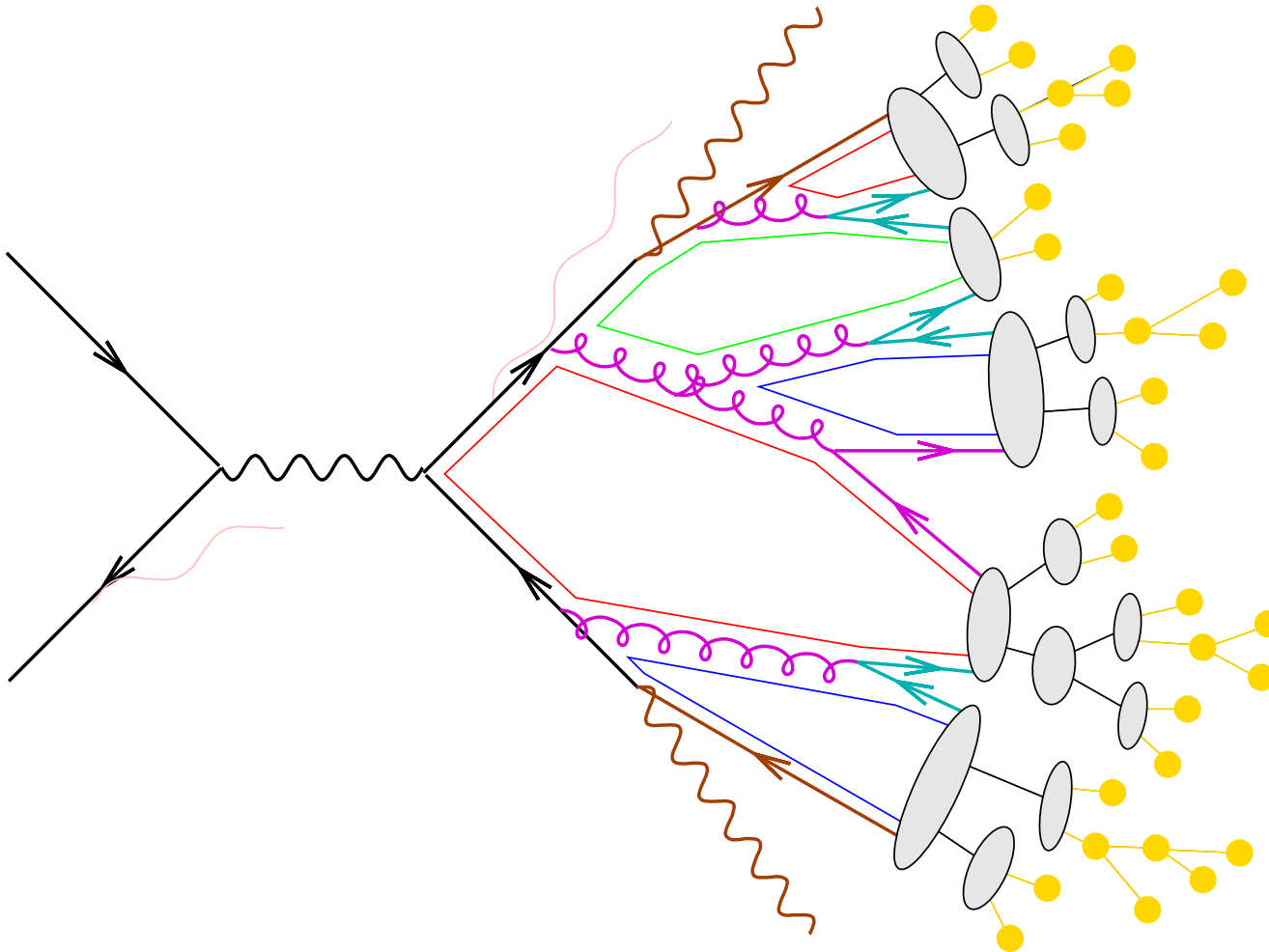
work with A Ribon, MH Seymour, P Stephens, BR Webber (Cambridge, CERN)

- Introduction
- Tour of Herwig++
- Results for e^+e^- Annihilation
- Outlook

SG, P. Stephens and B. Webber, JHEP **0312** (2003) 045 [hep-ph/0310083]

SG, A. Ribon, M. H. Seymour, P. Stephens and B. Webber, JHEP **0402** (2003) 005 [hep-ph/0311208]

e^+e^- Event Generator

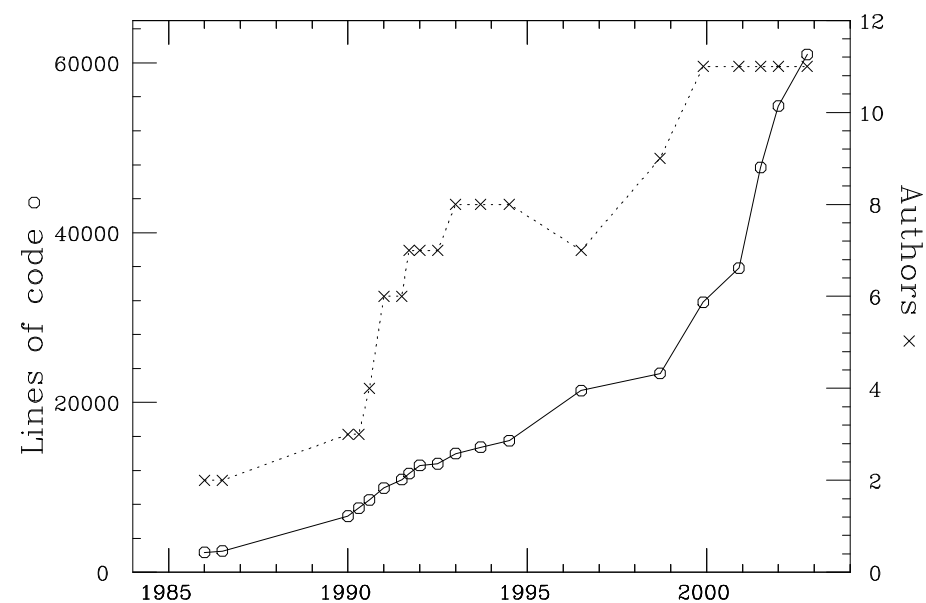


- hard scattering
- (QED) initial/final state radiation
- partonic decays, e.g. $t \rightarrow bW$
- parton shower evolution
- nonperturbative gluon splitting
- colour singlets
- colourless clusters
- cluster fission
- cluster $\rightarrow$ hadrons
- hadronic decays

The new generator Herwig++

Complete rewrite of HERWIG in C++

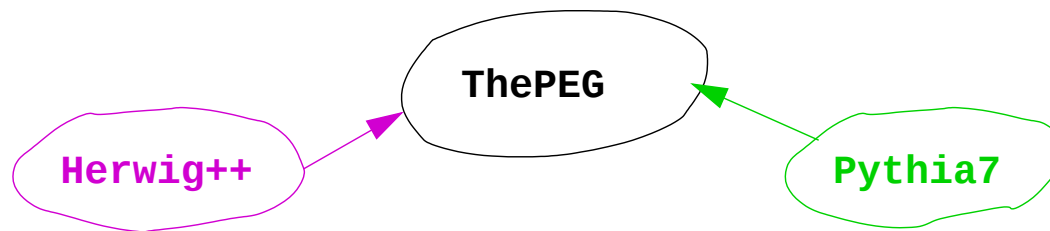
- aiming at full multi-purpose generator for LHC and future colliders.
- Preserve main features of HERWIG such as
 - angular ordered parton shower
 - Cluster Hadronization
- New features and improvements
 - improved parton shower evolution for heavy quarks
 - consistent radiation from unstable particles



HERWIG's growth. . .

Use of ThePEG in Herwig++

ThePEG = Toolkit for high energy Physics Event Generation



Won't re-invent the wheel

Share administrative overhead, common to event generators with Pythia7

Independent *physics* implementation

Large but very flexible implementation

Common basis for Pythia7/Herwig++:

- ✗ Lack of independence.
- ✗ Miss the possibility to test codes against each other.
- ✓ Physics, however, is still independent.
- ✓ Beneficial for the user to have the same framework.
- ✓ Running Herwig++ with the Lund String Fragmentation from Pythia7 is very simple!

Hard interactions

- Basic ME's included in **ThePEG**, such as:

$$e^+e^- \rightarrow q\bar{q}, \text{ partonic } 2 \rightarrow 2,$$

we use them.

- Soft and hard **matrix element corrections** implemented for $e^+e^- \rightarrow q\bar{q}g$.
- **AMEGIC++** will provide arbitrary ME's for multiparton final states via **AMEGICInterface**.
- CKKW ME+PS foreseen.
- Other authors can easily include their own matrix elements ($\rightarrow$ *safety* of OO code)

Quasi-Collinear Limit (Heavy Quarks)

Sudakov-basis p, n with $p^2 = M^2$ ('forward'), $n^2 = 0$ ('backward'),

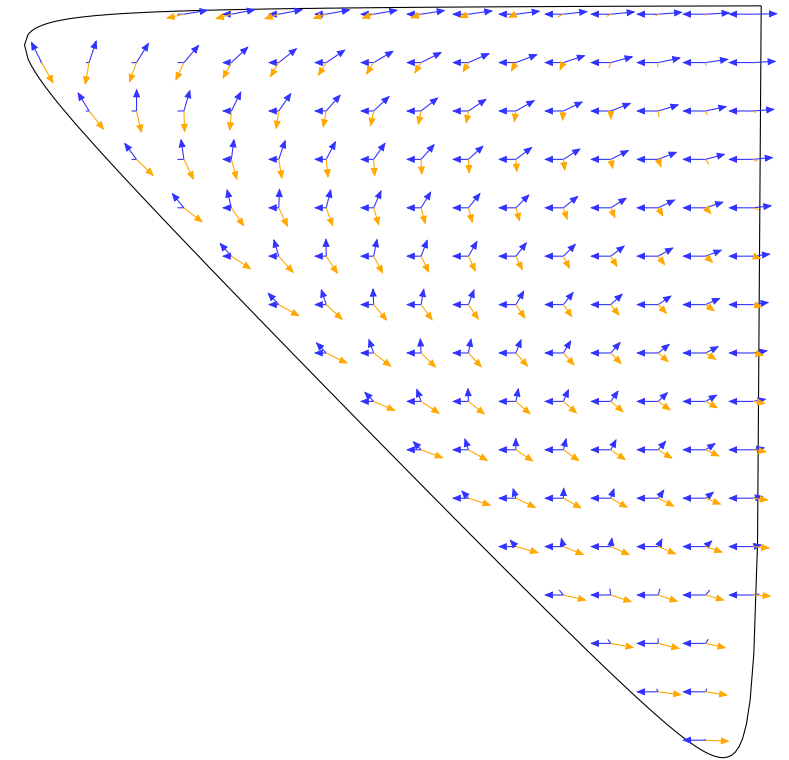
$$\begin{aligned} p_q &= zp + \beta_q n - q_\perp \\ p_g &= (1 - z)p + \beta_g n + q_\perp \end{aligned}$$

Collinear limit for radiation off heavy quark,

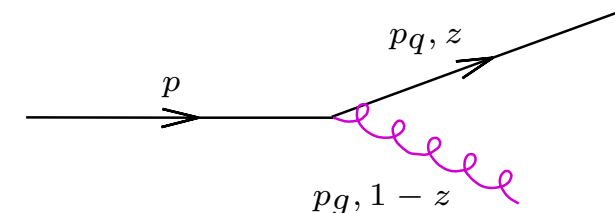
$$\begin{aligned} P_{gq}(z, \mathbf{q}^2, m^2) &= C_F \left[\frac{1 + z^2}{1 - z} - \frac{2z(1 - z)m^2}{\mathbf{q}^2 + (1 - z)^2 m^2} \right] \\ &= \frac{C_F}{1 - z} \left[1 + z^2 - \frac{2m^2}{z\tilde{q}^2} \right] \end{aligned}$$

→ $\tilde{q}^2 \sim \mathbf{q}^2$ may be used as evolution variable.

$q\bar{q}g$ -Phase space $(x, \bar{x})$

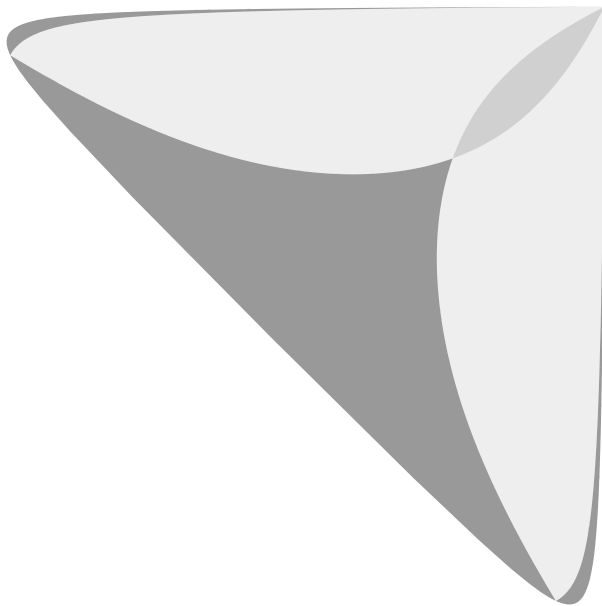


Single emission:

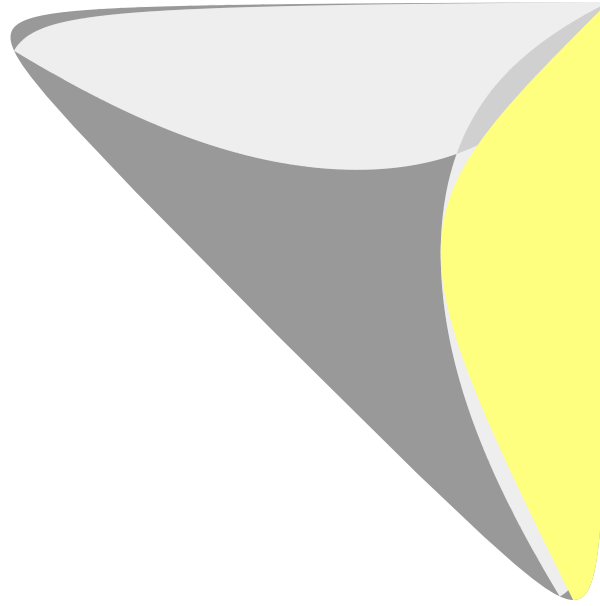


$q\bar{q}g$ Phase Space old vs new variables

Consider $(x, \bar{x})$ phase space for $e^+e^- \rightarrow q\bar{q}g$



HERWIG



Comparison



Herwig++

- ✗ Larger dead region with new variables.
- ✓ Smooth coverage of soft gluon region.
- ✓ No overlapping regions in phase space.

Universal cutoff parameter δ

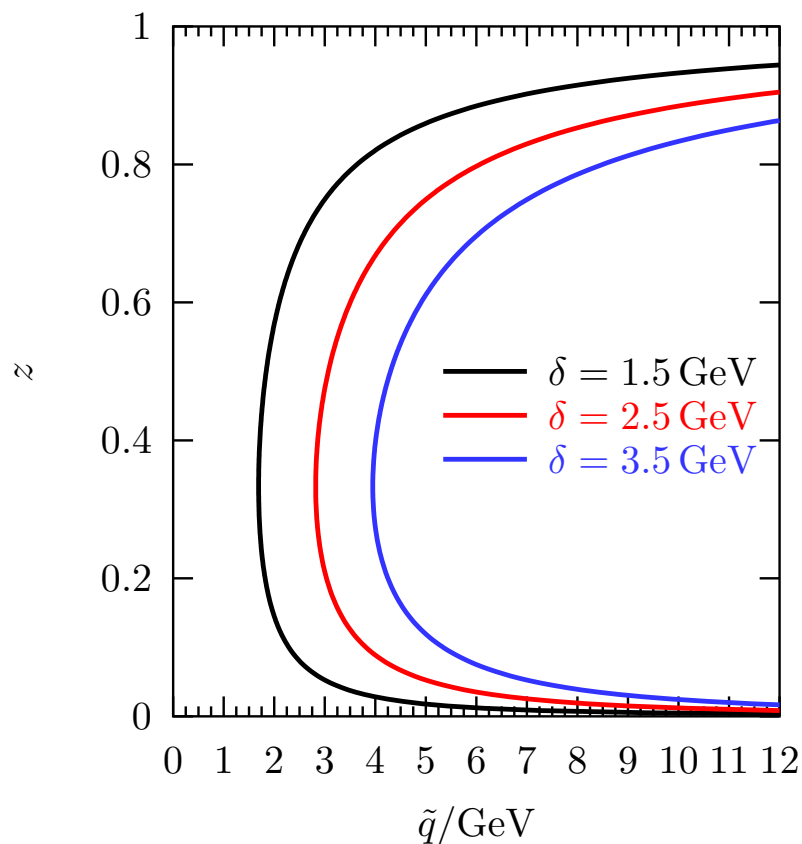
Require threshold in parton shower phase space

$$Q_{\text{thr}} = \beta m_q + \delta \quad (\beta = 0.85)$$

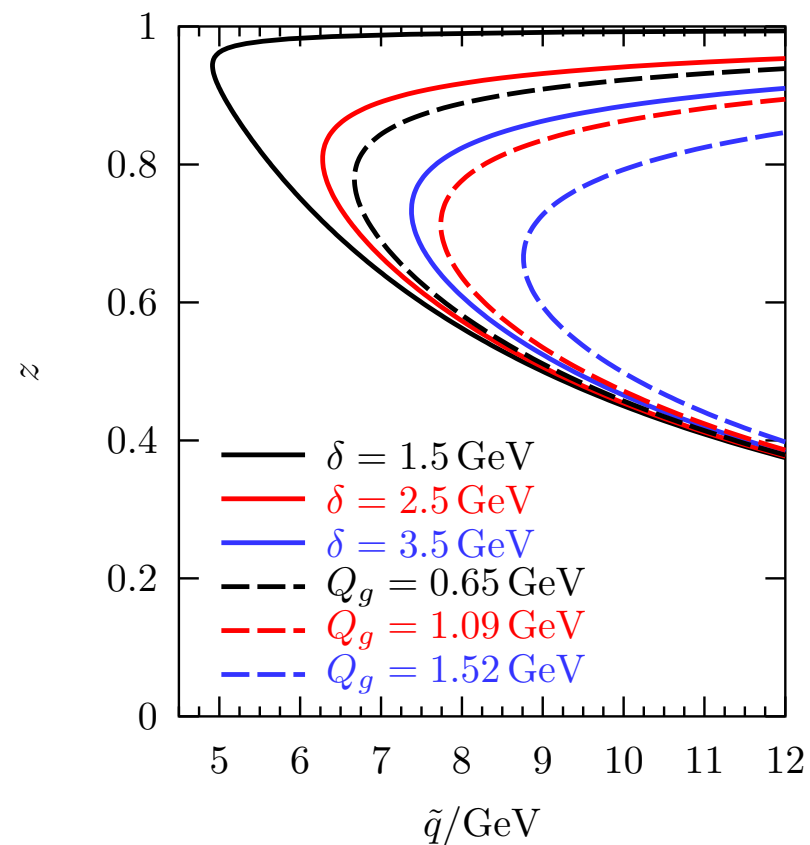
parametrization of Q_g in terms of δ, m_q

$$Q_g = \frac{\delta - 0.3m_q}{2.3}.$$

light quarks:

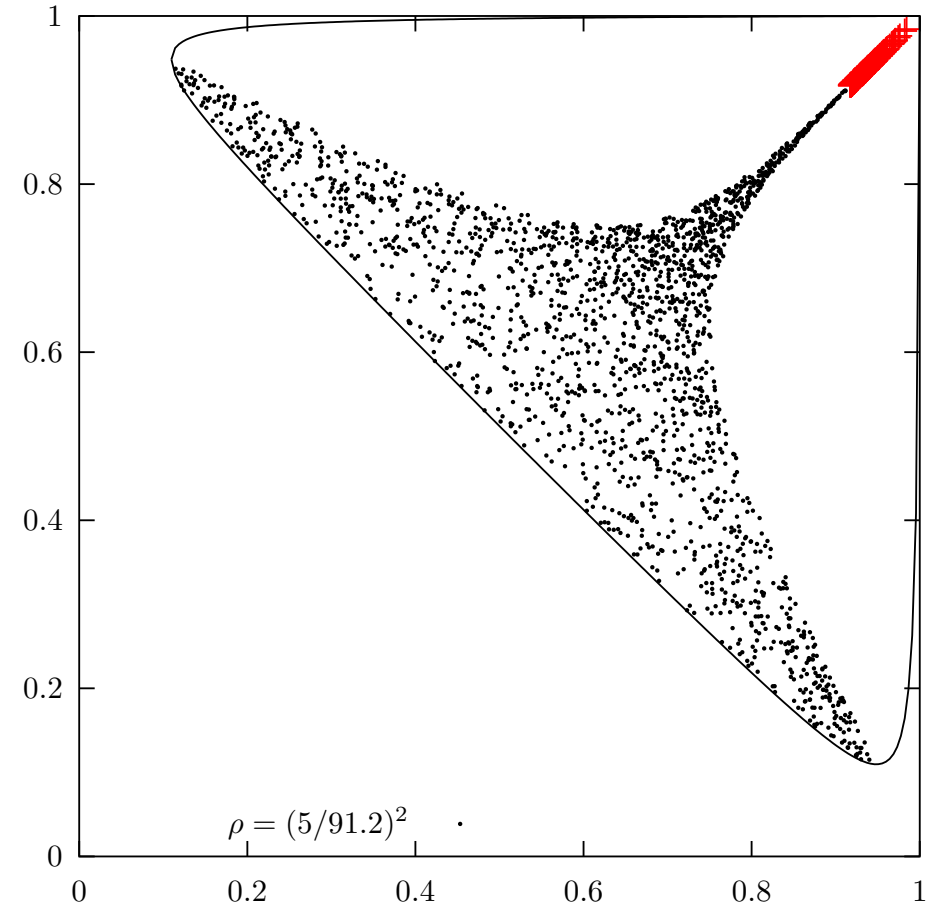


b quarks:



Hard Matrix Element Corrections

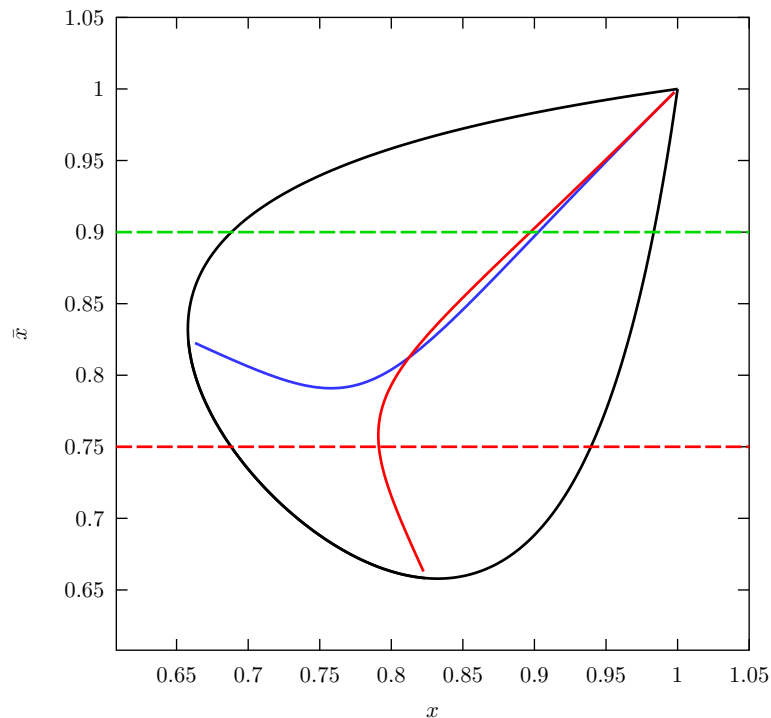
- Points $(x, \bar{x})$ in **dead region** chosen acc to LO $e^+e^- \rightarrow q\bar{q}g$ matrix element and accepted acc to ME weight.
- About **3%** of all events are actually hard $q\bar{q}g$ events.
- Red points have **weight** > 1 , practically no error by setting weight to one.
- Event **oriented** according to given $q\bar{q}$ geometry. Quark direction is kept with weight $x^2/(x^2 + \bar{x}^2)$.



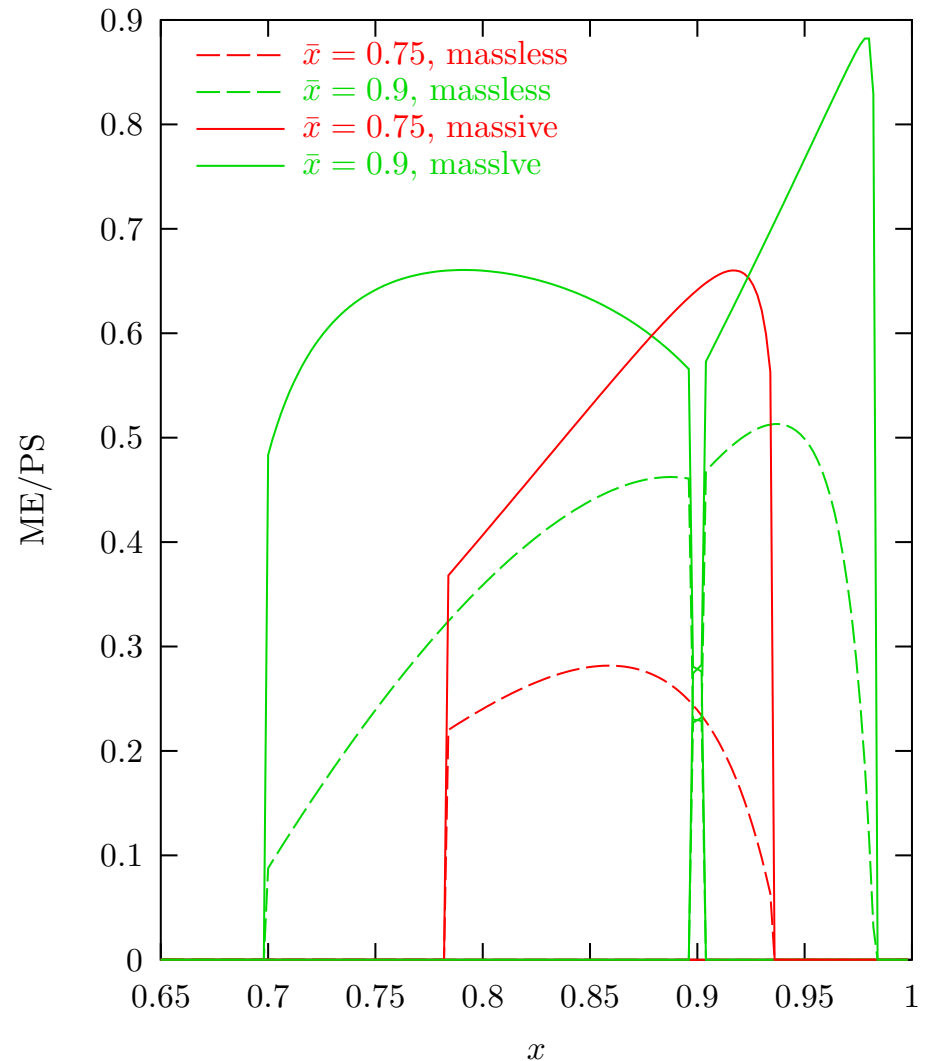
Soft Matrix Element Corrections

- Ratio ME/PS compares emission with result from true ME if slightly away from soft/collinear region.
- **Veto** on 'hardest emission so far' in $p_{\perp}$.
- **Massive splitting function** *very important*!

Example with heavy quark, $m^2/Q^2 = 0.1$:



Comparison with massless splitting function



Cluster hadronization in a nutshell

- **Nonperturbative $g \rightarrow q\bar{q}$ splitting** ($q = uds$) isotropically.
Here, $m_g \approx 750 \text{ MeV} > 2m_q$.
- **Cluster formation**, universal spectrum (see below)
- **Cluster fission**, until

$$M^p < M_{\text{max}}^p + (m_1 + m_2)^p$$

where masses are chosen from

$$M_i = \left[\left(M^p - (m_i + m_3)^p \right) r_i + (m_i + m_3)^p \right]^{1/p},$$

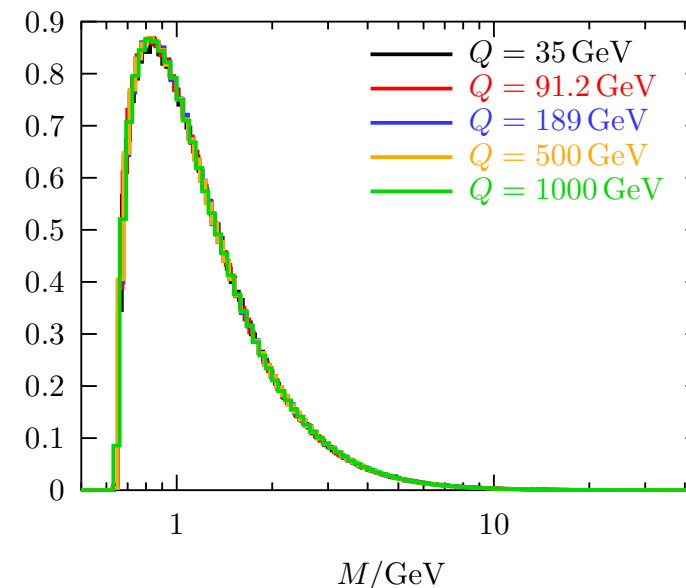
with additional phase space constraints. Constituents keep moving in their original direction.

- **Cluster Decay**

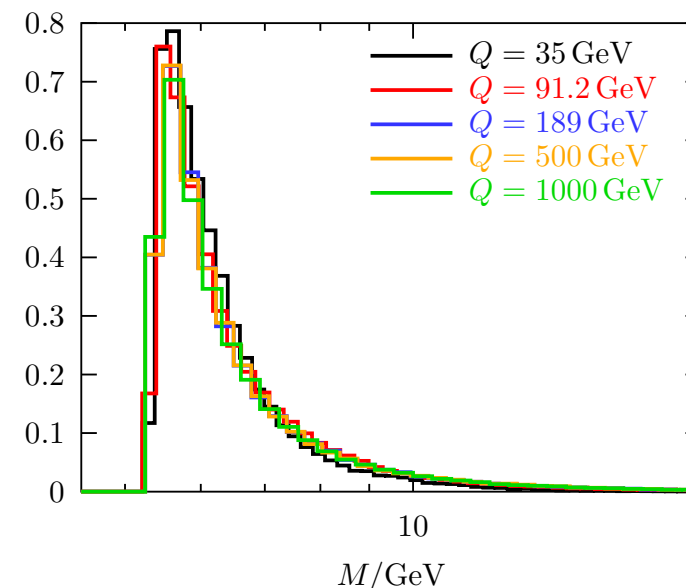
$$P(a_{i,q}, b_{q,j} | i, j) = \frac{W(a_{i,q}, b_{q,j} | i, j)}{\sum_{M/B} W(c_{i,q'}, d_{q',j} | i, j)}.$$

New! Meson/Baryon ratio is parametrized in terms of diquark weight. In HERWIG the sum ran over all possible hadrons.

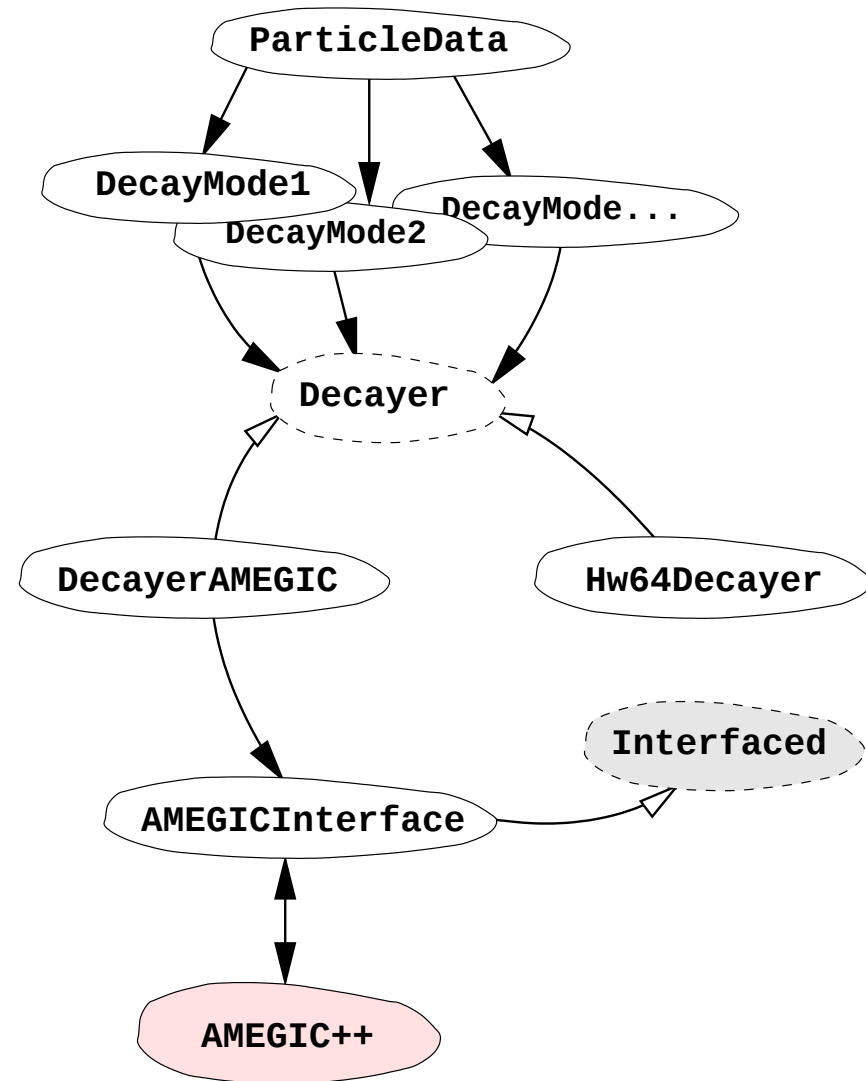
Primary Light Clusters



Primary b -Clusters



Decays



- FORTRAN HERWIG is reproduced with **Hw64Decayer** using the same Matrix element codes as before (will be used for hadronic decays right now)
- **DecayerAMEGIC** gets final states for a decay mode directly from **AMEGIC++**
- Room for improvement. . .

Results

Aim: test new parton shower and its interplay with other parts of the program in great detail.

- $e^+e^- \rightarrow$ jets, mainly at Z^0 .
- hard and soft ME corrections.
- Cutoff parameter δ varied.

Analysis:

- Each bin i : data $D_i \pm \delta D_i$ and MC result $M_i \pm \delta D_i$.
- Normalization N MC/data not necessarily given.
- $\delta D_i/D_i > 5\%$ to emphasize *global* strategy.
- Then

$$\chi^2 = \sum_i \chi_i^2 = \sum_i \frac{(D_i - N M_i)^2}{\delta D_i^2 + \delta M_i^2}$$
$$R_i = \frac{M_i - D_i}{D_i} \pm \left(\frac{\delta M_i}{D_i} \oplus \frac{M_i \delta D_i}{D_i^2} \right)$$

Most plots in the following include

- Histograms with $\delta = 1.7, 2.3, 3.2$ GeV and data.
- R_i compared to $\delta D_i/D_i$ (yellow band).
- χ_i^2/χ^2 for each bin.

Hadron Multiplicities

Particle	Experiment	Measured	Old Model	Herwig++	Fortran
All Charged	M,A,D,L,O	20.924 ± 0.117	20.22*	20.814	20.532*
γ	A,O	21.27 ± 0.6	23.032	22.67	20.74
π^0	A,D,L,O	9.59 ± 0.33	10.27	10.08	9.88
$\rho(770)^0$	A,D	1.295 ± 0.125	1.235	1.316	1.07
$\pi^\pm$	A,O	17.04 ± 0.25	16.30	16.95	16.74
$\rho(770)^\pm$	O	2.4 ± 0.43	1.99	2.14	2.06
η	A,L,O	0.956 ± 0.049	0.886	0.893	0.669*
$\omega(782)$	A,L,O	1.083 ± 0.088	0.859	0.916	1.044
$\eta'(958)$	A,L,O	0.152 ± 0.03	0.13	0.136	0.106
K^0	S,A,D,L,O	2.027 ± 0.025	2.121*	2.062	2.026
$K^*(892)^0$	A,D,O	0.761 ± 0.032	0.667	0.681	0.583*
$K^*(1430)^0$	D,O	0.106 ± 0.06	0.065	0.079	0.072
$K^\pm$	A,D,O	2.319 ± 0.079	2.335	2.286	2.250
$K^*(892)^\pm$	A,D,O	0.731 ± 0.058	0.637	0.657	0.578
$\phi(1020)$	A,D,O	0.097 ± 0.007	0.107	0.114	0.134*
p	A,D,O	0.991 ± 0.054	0.981	0.947	1.027
Δ^{++}	D,O	0.088 ± 0.034	0.185	0.092	0.209*
Σ^-	O	0.083 ± 0.011	0.063	0.071	0.071
Λ	A,D,L,O	0.373 ± 0.008	0.325*	0.384	0.347*
Σ^0	A,D,O	0.074 ± 0.009	0.078	0.091	0.063
Σ^+	O	0.099 ± 0.015	0.067	0.077	0.088
$\Sigma(1385)^\pm$	A,D,O	0.0471 ± 0.0046	0.057	0.0312*	0.061*
Ξ^-	A,D,O	0.0262 ± 0.001	0.024	0.0286	0.029
$\Xi(1530)^0$	A,D,O	0.0058 ± 0.001	0.026*	0.0288*	0.009*
Ω^-	A,D,O	0.00125 ± 0.00024	0.001	0.00144	0.0009

Hadron Multiplicities (ctd')

Particle	Experiment	Measured	Old Model	Herwig++	Fortran
$f_2(1270)$	D,L,O	0.168 ± 0.021	0.113	0.150	0.173
$f_2'(1525)$	D	0.02 ± 0.008	0.003	0.012	0.012
$D^\pm$	A,D,O	0.184 ± 0.018	0.322*	0.319*	0.283*
$D^*(2010)^\pm$	A,D,O	0.182 ± 0.009	0.168	0.180	0.151*
D^0	A,D,O	0.473 ± 0.026	0.625*	0.570*	0.501
$D_s^\pm$	A,O	0.129 ± 0.013	0.218*	0.195*	0.127
$D_s^{*\pm}$	O	0.096 ± 0.046	0.082	0.066	0.043
J/Ψ	A,D,L,O	0.00544 ± 0.00029	0.006	0.00361*	0.002*
Λ_c^+	D,O	0.077 ± 0.016	0.006*	0.023*	0.001*
$\Psi'(3685)$	D,L,O	0.00229 ± 0.00041	0.001*	0.00178	0.0008*

of *'s = observables with more than 3σ deviation:

OldModel : Herwig++ : Fortran = 9 : 7 : 13

$k_{\perp}$ Algorithm (“Durham”–Algorithm)

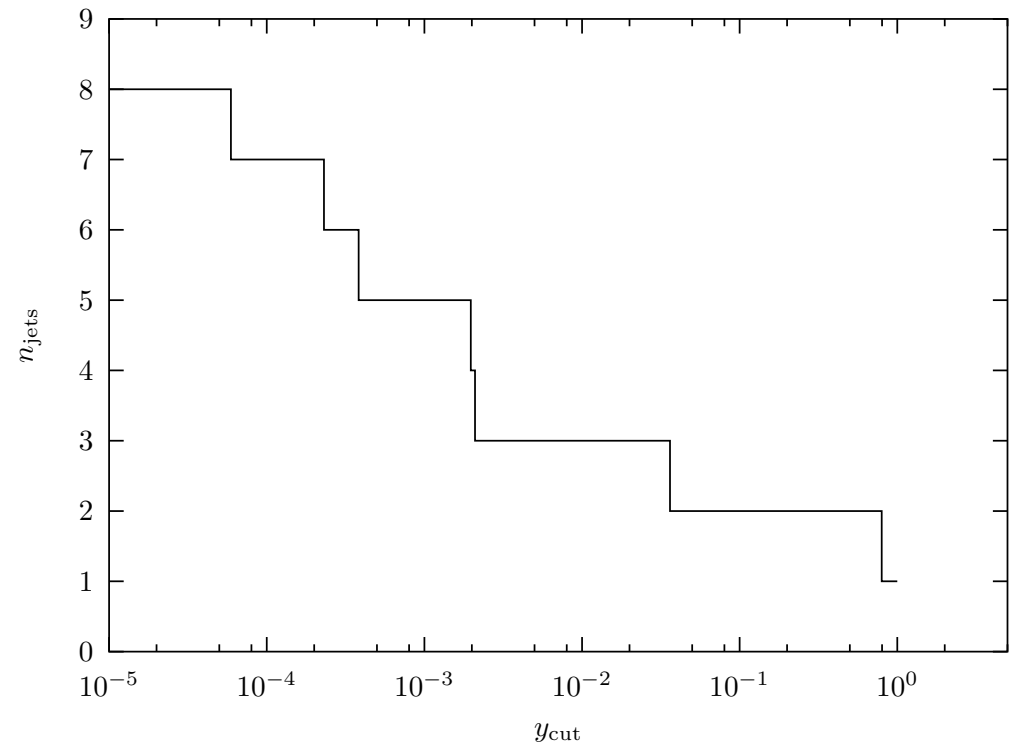
- For each pair i, j of particles in an event calculate the ‘distance’

$$y_{ij} = \frac{2 \min(E_i^2, E_j^2)}{Q^2} (1 - \cos \theta_{ij}).$$

[resolution scale $Q_{ij} = Q \sqrt{y_{ij}}$]

- The pair with minimum y_{ij} is **clustered** into a pseudoparticle with momentum $p = p_i + p_j$.
- Stop** when all particles are clustered or all $y_{ij} > y_{\text{cut}}$.

We have an **n -jet** event if n particles or pseudoparticles are left at a given y_{cut} .

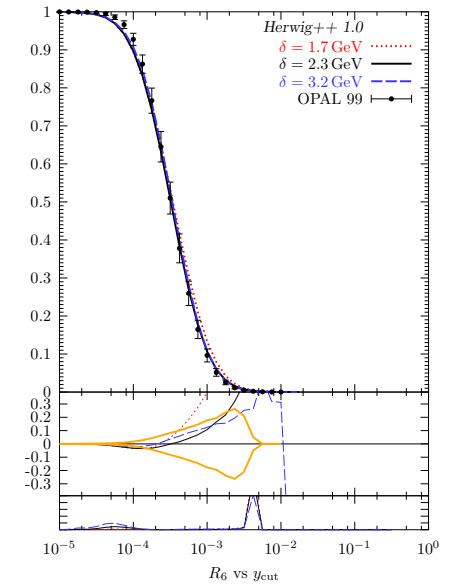
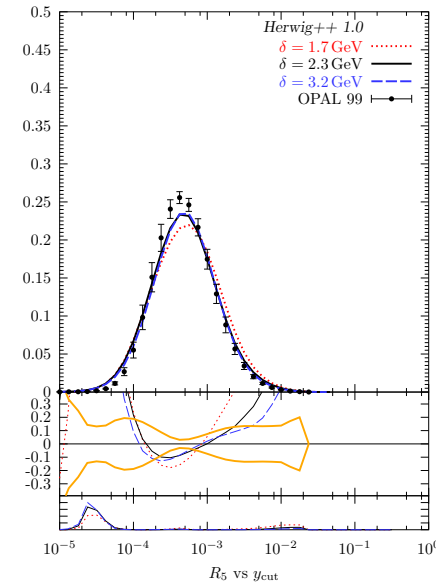
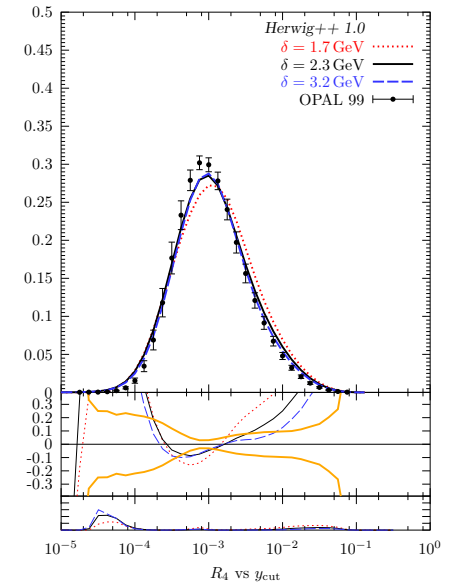
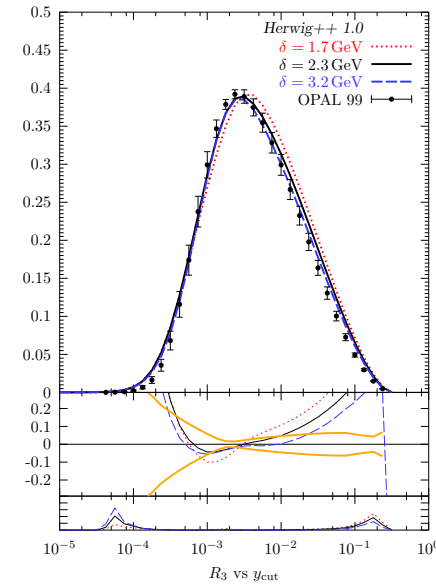
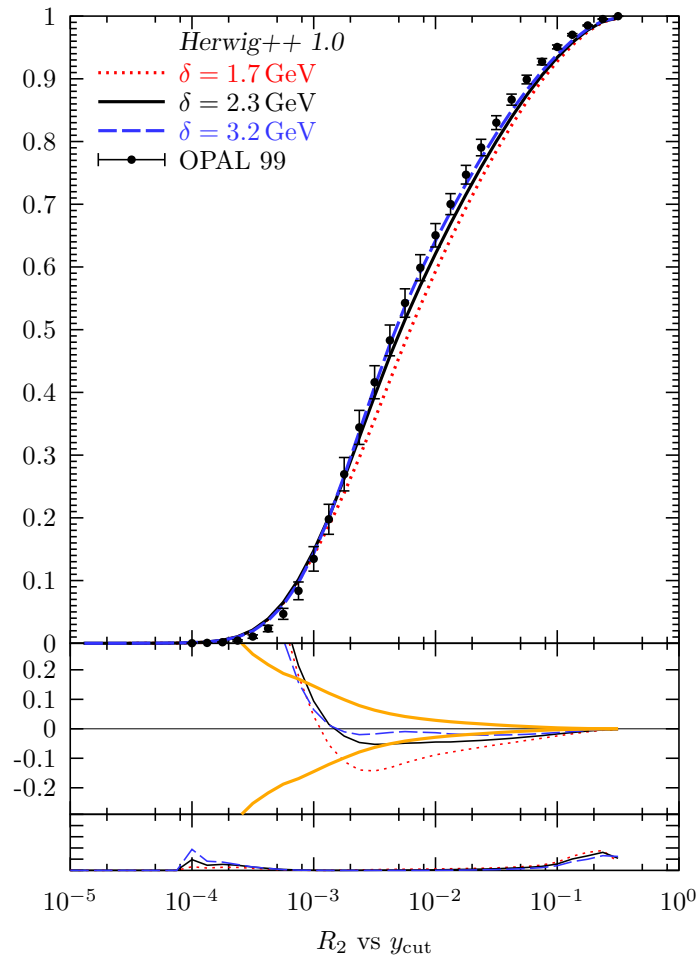


There are other jet algorithms with a different distance measure (JADE) or recombination scheme. We have interfaced the KtJet package for jet clustering.

Jet Rates

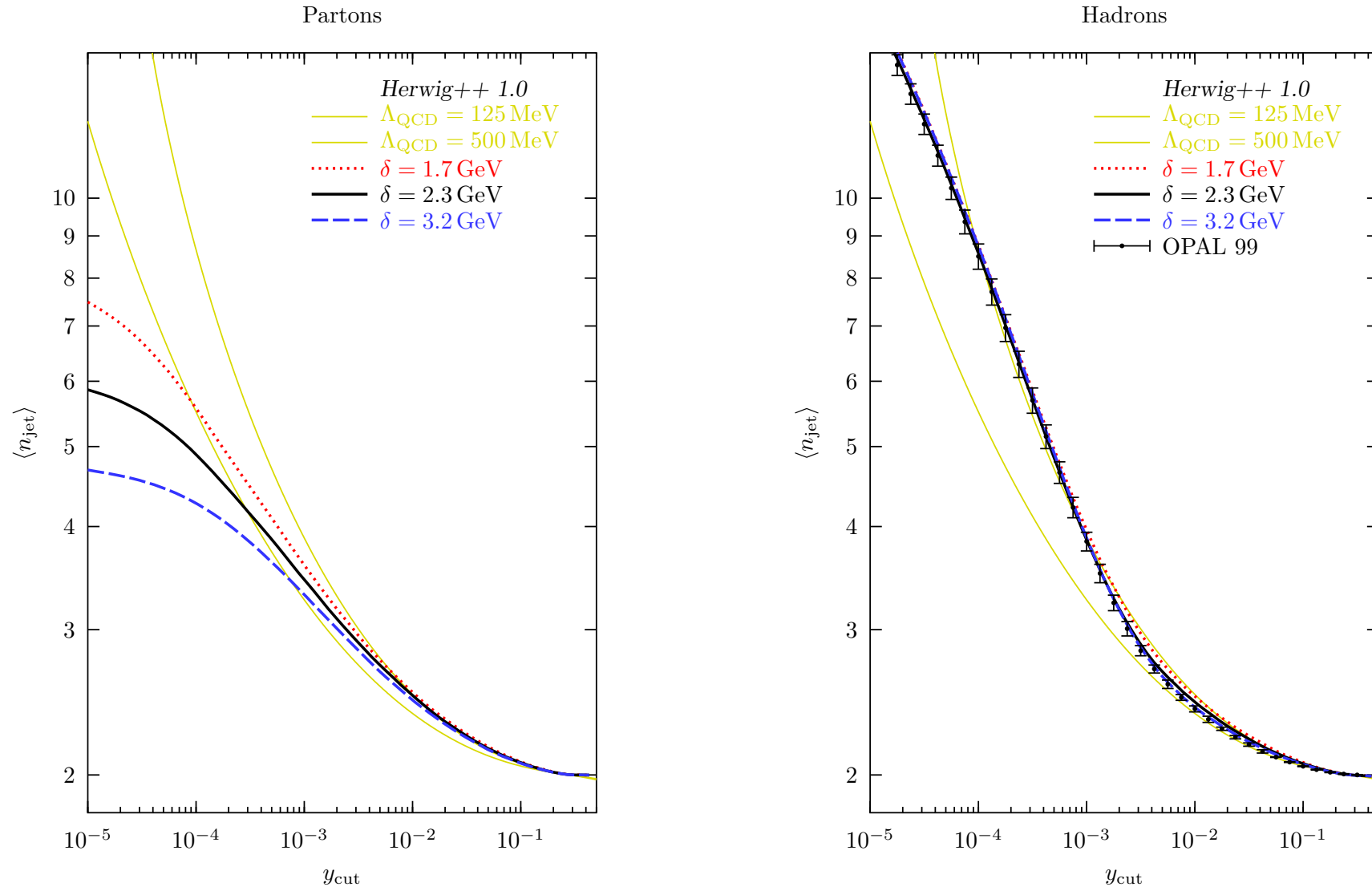
$$R_n = \sigma(n\text{-jets})/\sigma(\text{jets}) \quad (n = 2..5)$$

$$R_6 = \sigma(> 5\text{-jets})/\sigma(\text{jets})$$



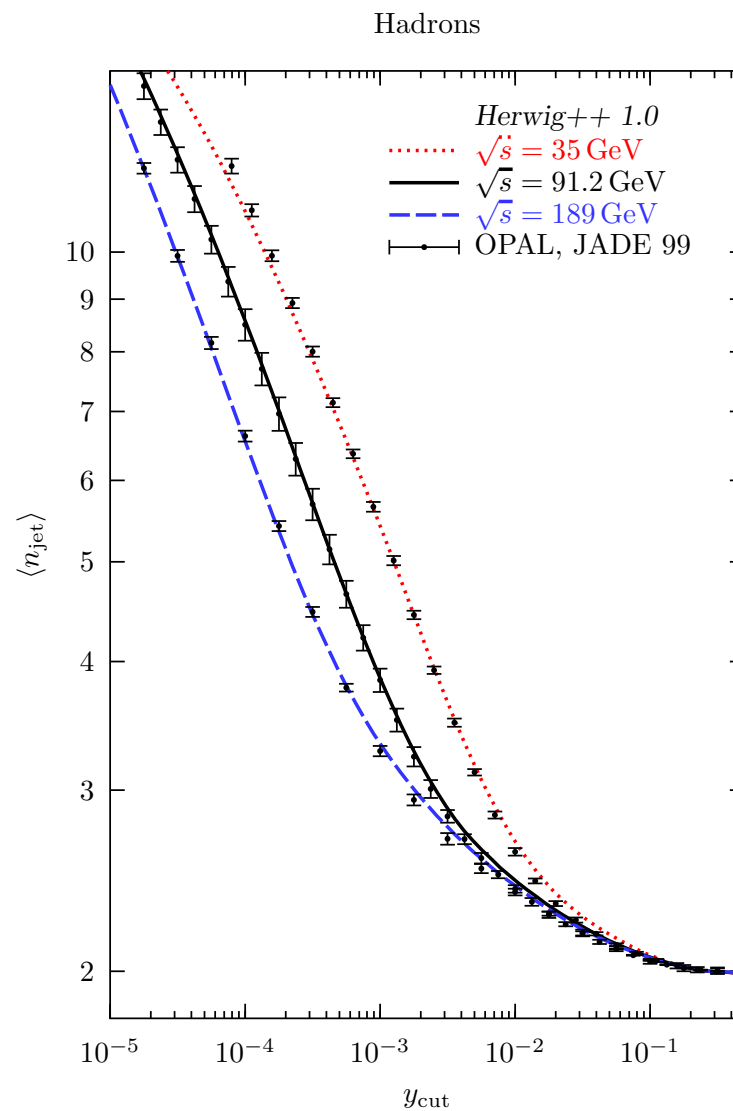
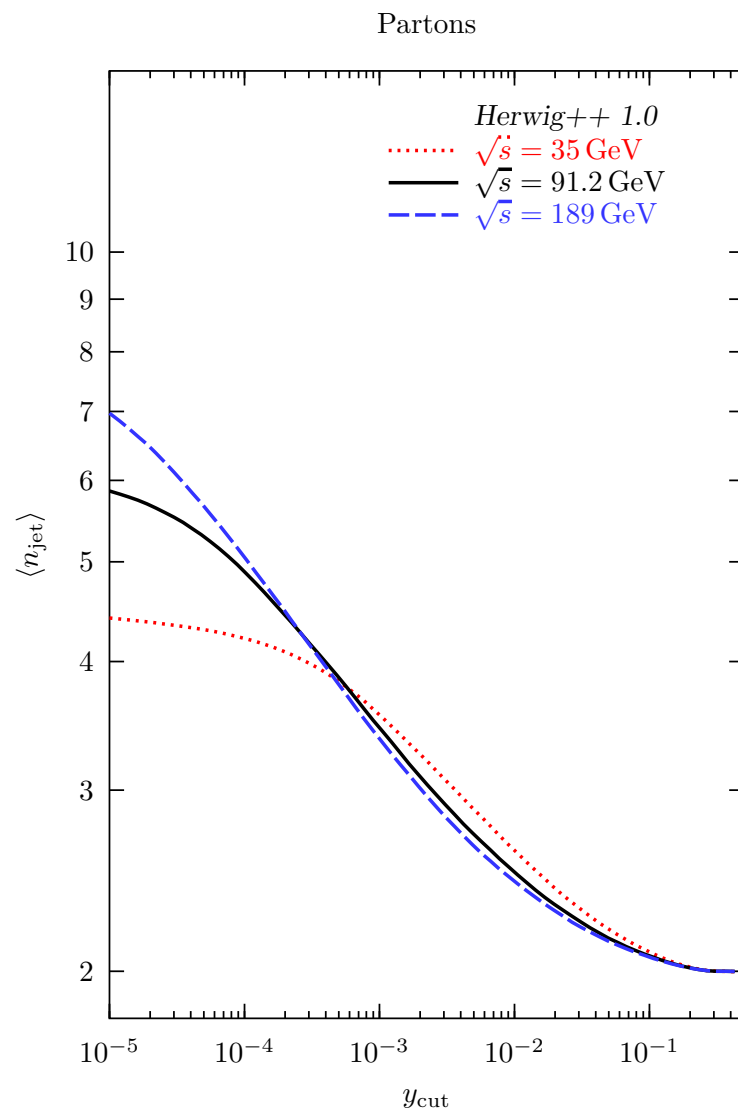
Jet Multiplicity

Durham algorithm. Smooth interplay between shower and hadronization.



Jet Multiplicity (PETRA, LEP, LEP II)

$$\sqrt{s} = \{35, 91.2, 189\} \text{ GeV}$$



Event Shape Variables, Definition

Thrust

$$F(\mathbf{n}) = \frac{\sum_{\alpha} |\mathbf{p}_{\alpha} \cdot \mathbf{n}|}{\sum_{\alpha} |\mathbf{p}_{\alpha}|}$$

Find $\mathbf{n}$, such that thrust

$$\begin{aligned} T &= \max_{\mathbf{n}} F(\mathbf{n}) \\ &= F(\mathbf{n}_T) , \end{aligned}$$

thrust major

$$\begin{aligned} M &= \max_{\mathbf{n} \perp \mathbf{n}_T} F(\mathbf{n}) \\ &= F(\mathbf{n}_M) , \end{aligned}$$

thrust minor

$$\begin{aligned} \mathbf{n}_m &= \mathbf{n}_T \times \mathbf{n}_M \\ m &= F(\mathbf{n}_m) \end{aligned}$$

Sphericity

$$Q_{ij} = \frac{\sum_{\alpha} (\mathbf{p}_{\alpha})_i (\mathbf{p}_{\alpha})_j}{\sum_{\alpha} \mathbf{p}_{\alpha}^2}$$

Diagonalize, eigenvalues

$$\begin{aligned} \lambda_1 &> \lambda_2 > \lambda_3 \\ \lambda_1 + \lambda_2 + \lambda_3 &= 1 \end{aligned}$$

Then

$$\begin{aligned} S &= \frac{3}{2}(\lambda_2 + \lambda_3) \\ P &= \lambda_2 - \lambda_3 \\ A &= \frac{3}{2}\lambda_3 \end{aligned}$$

Eigenvector $\mathbf{n}_S$ sphericity axis
etc.

C, D parameter

$$L_{ij} = \frac{\sum_{\alpha} (\mathbf{p}_{\alpha})_i (\mathbf{p}_{\alpha})_j / |\mathbf{p}_{\alpha}|}{\sum_{\alpha} |\mathbf{p}_{\alpha}|}$$

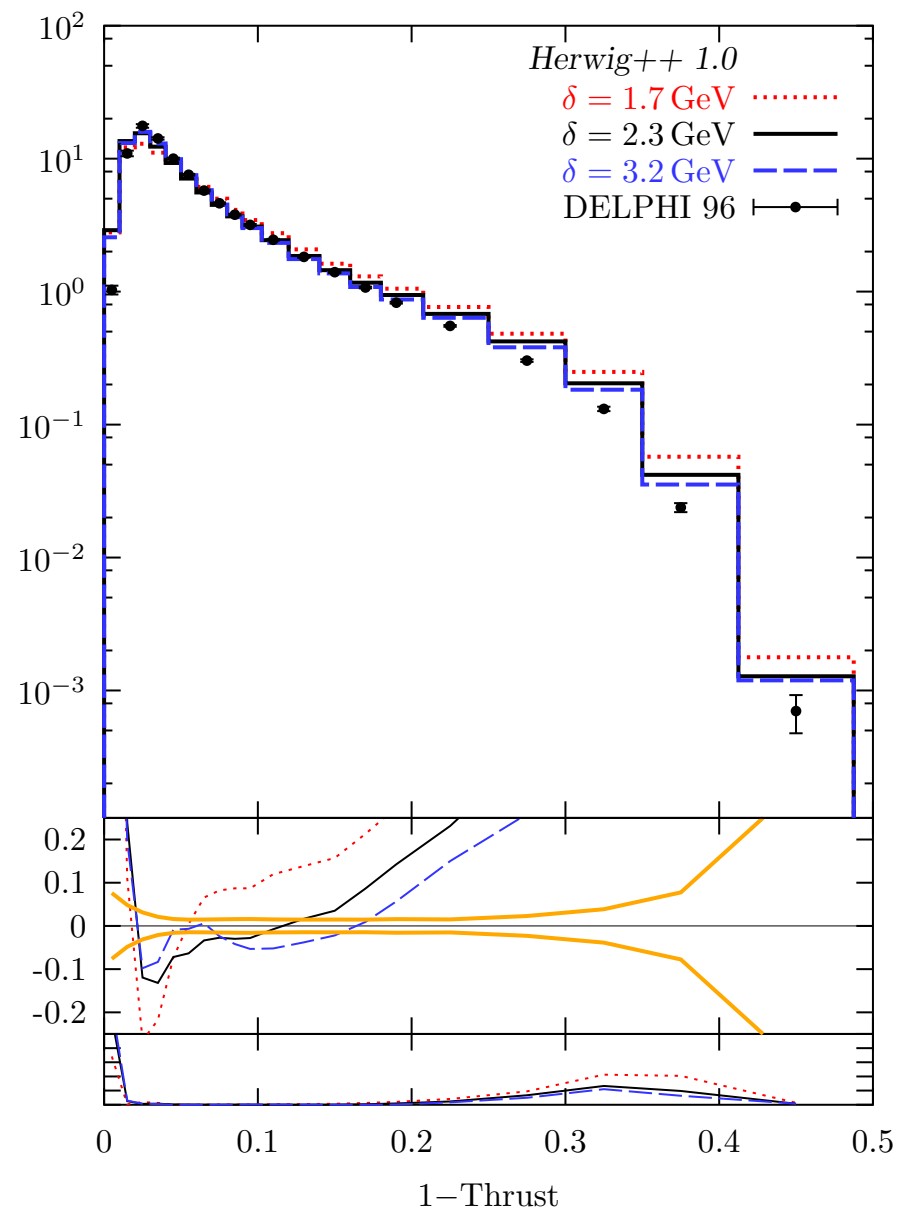
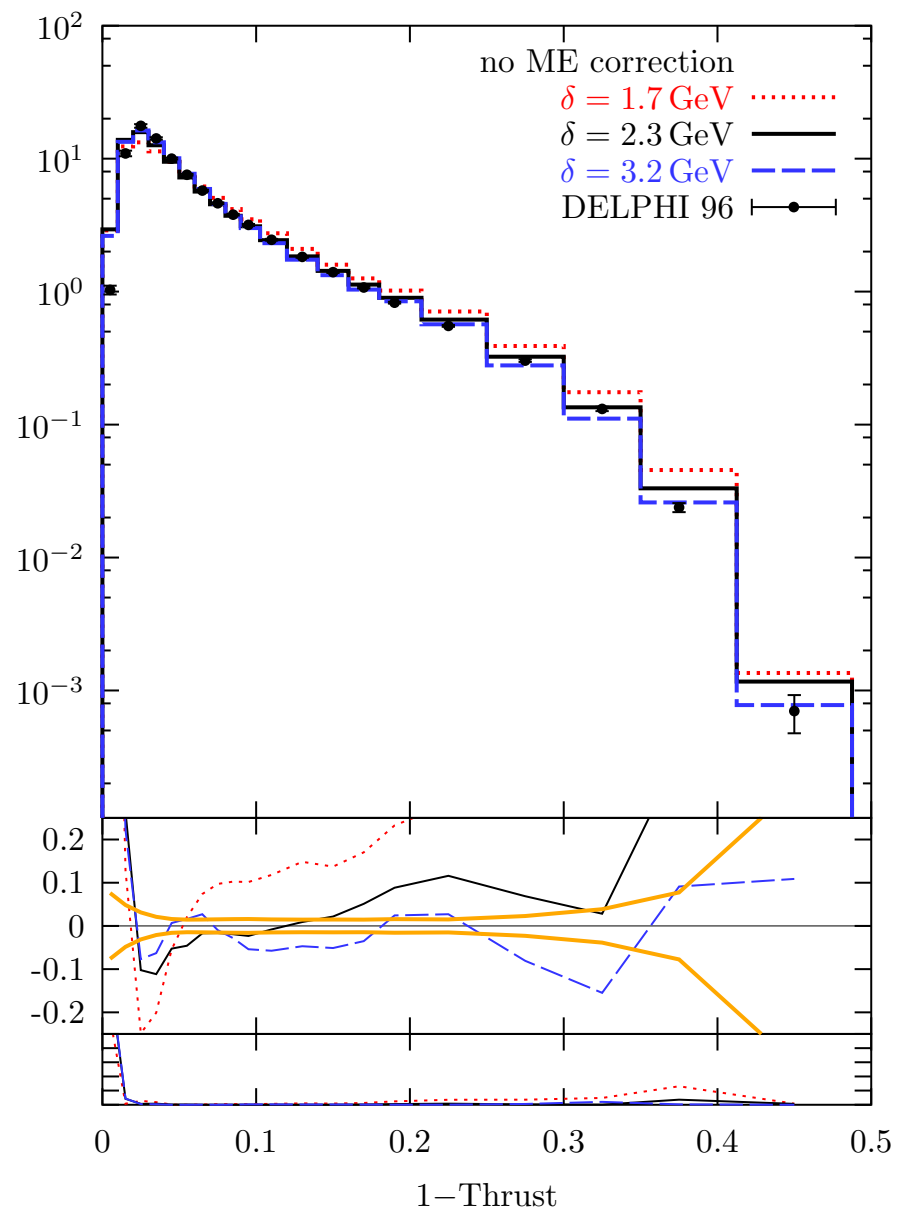
Diagonalize, eigenvalues

$$\lambda_1 + \lambda_2 + \lambda_3 = 1$$

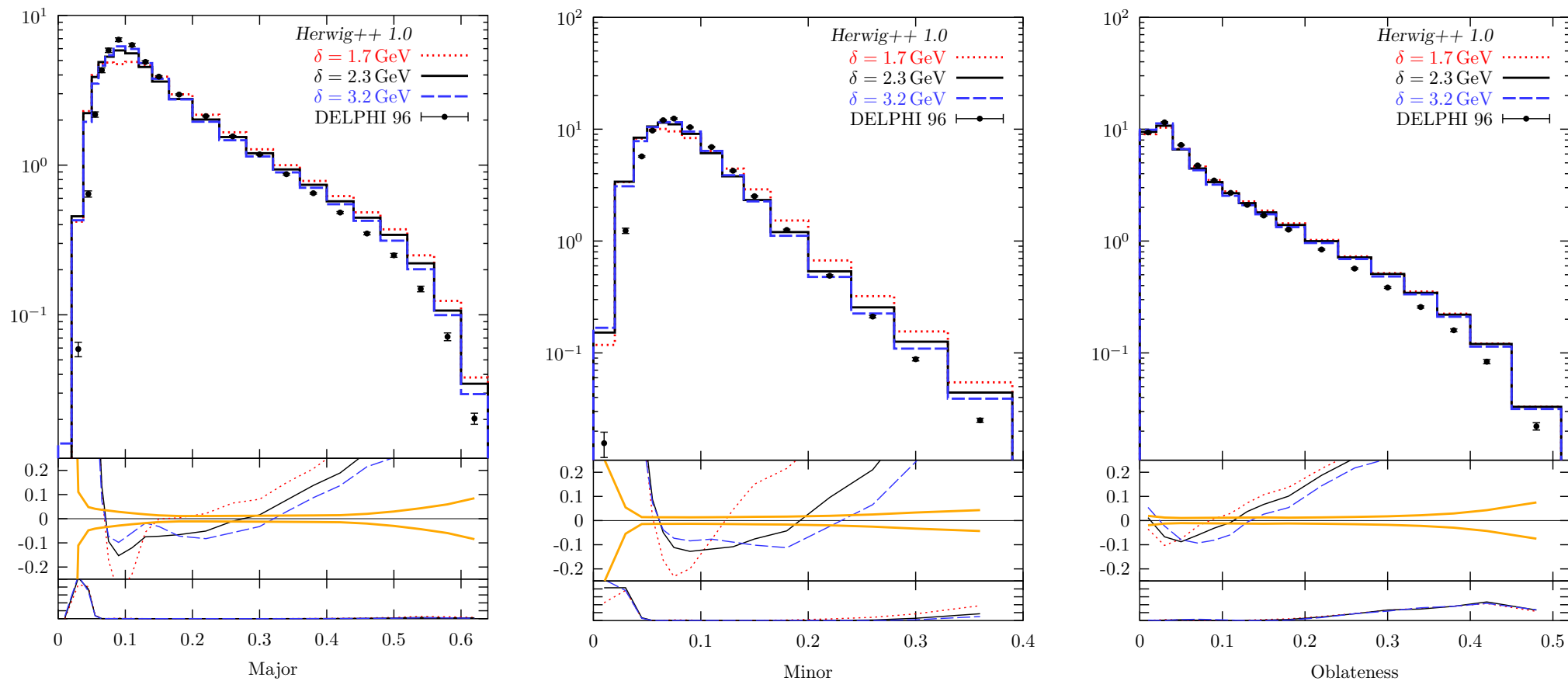
and define

$$\begin{aligned} C &= 3(\lambda_1 \lambda_2 + \lambda_2 \lambda_3 + \lambda_3 \lambda_1) \\ D &= 27 \lambda_1 \lambda_2 \lambda_3 \end{aligned}$$

Thrust — ME Corrections off/on



Major, Minor, Oblateness



All Thrust-related distributions slightly wide, ie too many 2-jet like on one side and too many spherical events on the other side.

Four Jet Angles — Definitions

Bengtsson–Zerwas angle

$$\chi_{BZ} = \angle(\mathbf{p}_1 \times \mathbf{p}_2, \mathbf{p}_3 \times \mathbf{p}_4)$$

Körner–Schierholz–Willrodt angle

$$\Phi_{KSW} = \frac{1}{2} [\angle(\mathbf{p}_1 \times \mathbf{p}_3, \mathbf{p}_2 \times \mathbf{p}_4) + \angle(\mathbf{p}_1 \times \mathbf{p}_4, \mathbf{p}_2 \times \mathbf{p}_3)]$$

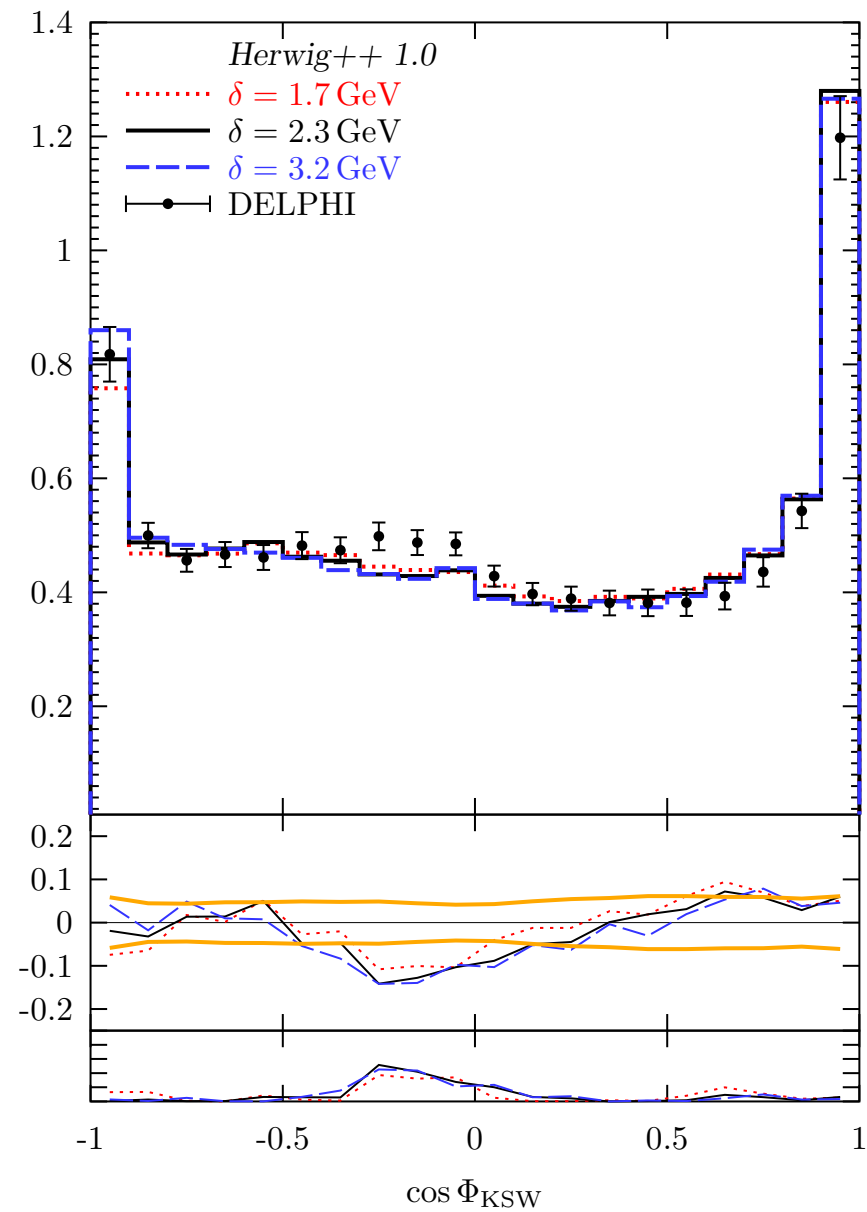
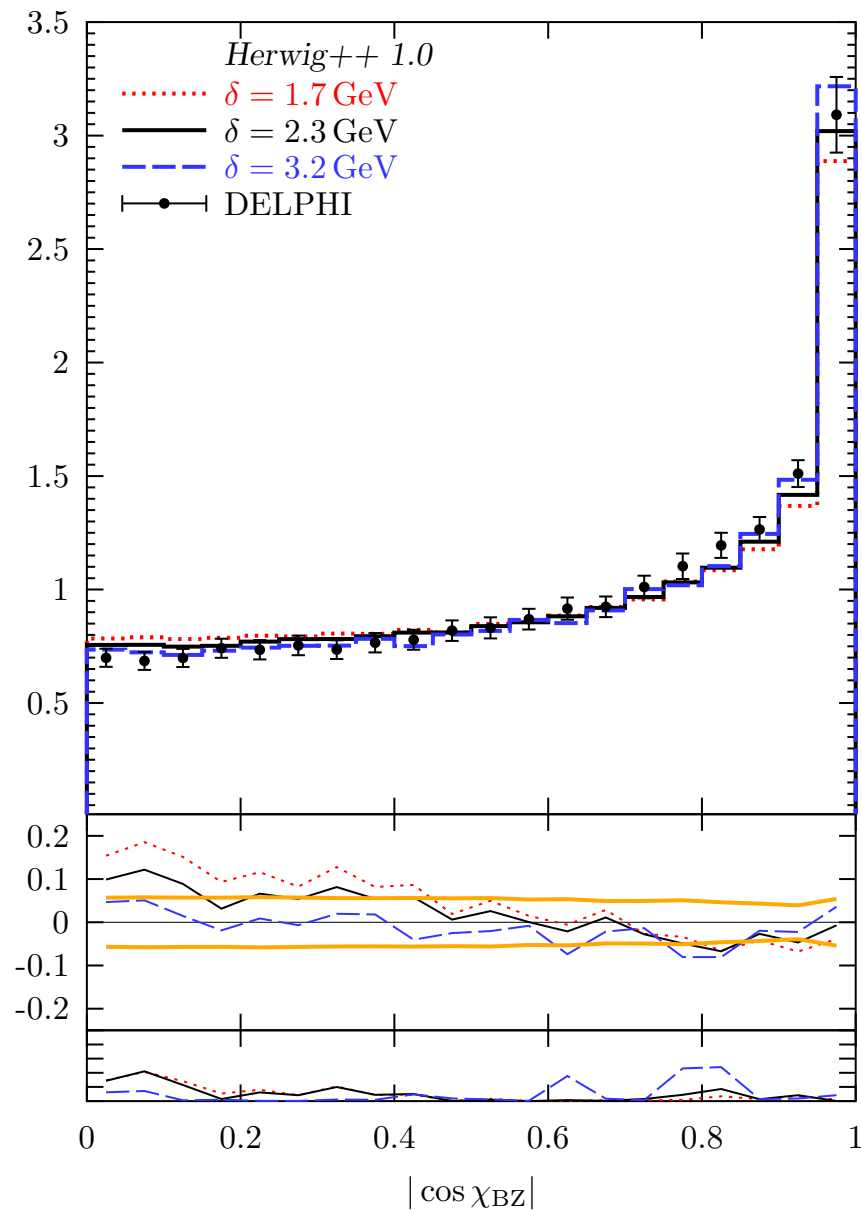
(modified) Nachtmann–Reiter angle

$$\theta_{NR}^* = \angle(\mathbf{p}_1 - \mathbf{p}_2, \mathbf{p}_3 - \mathbf{p}_4)$$

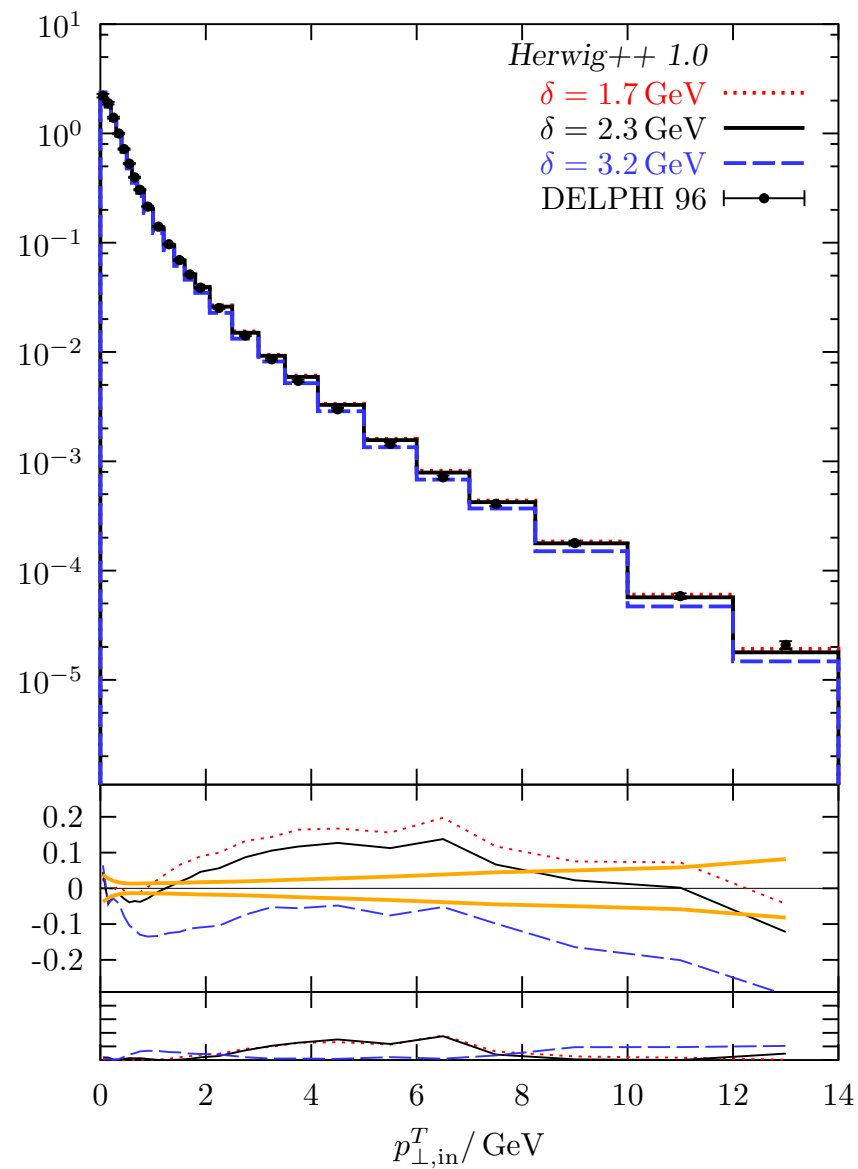
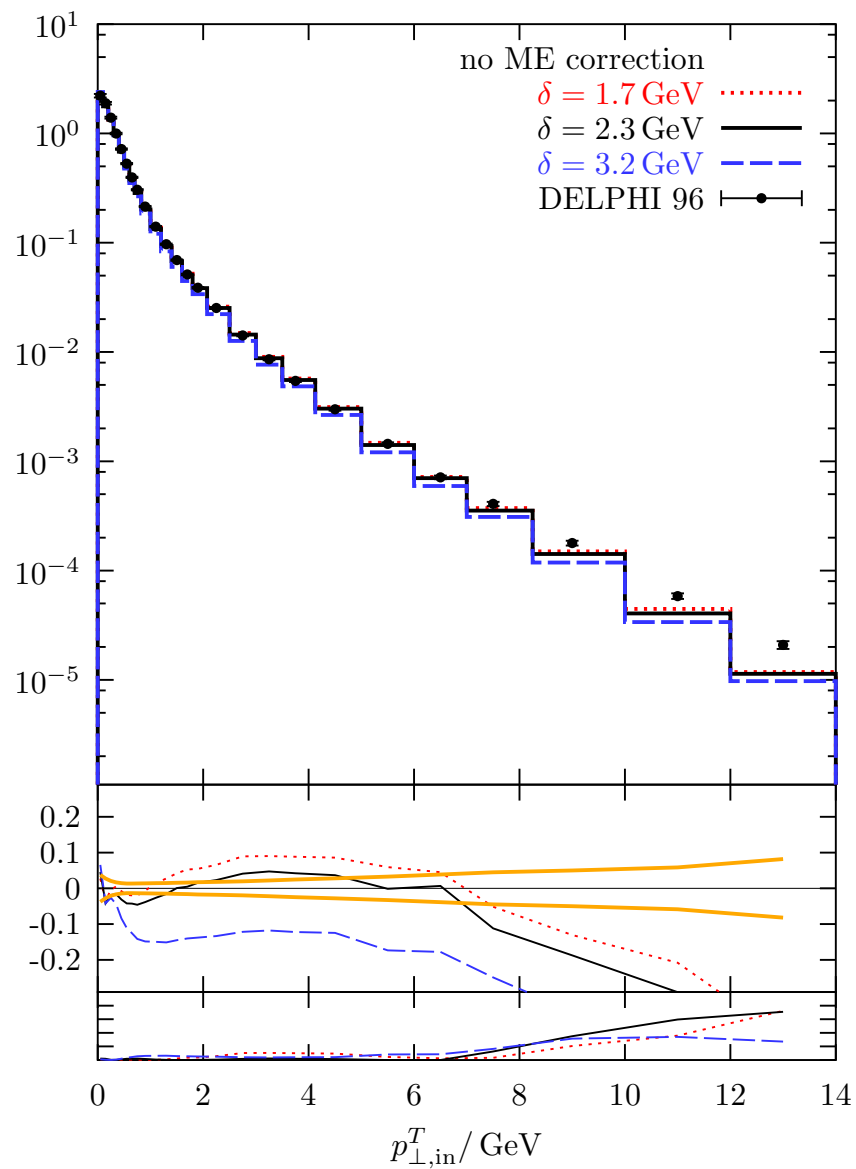
α_{34} :

$$\alpha_{34} = \angle(\mathbf{p}_3, \mathbf{p}_4)$$

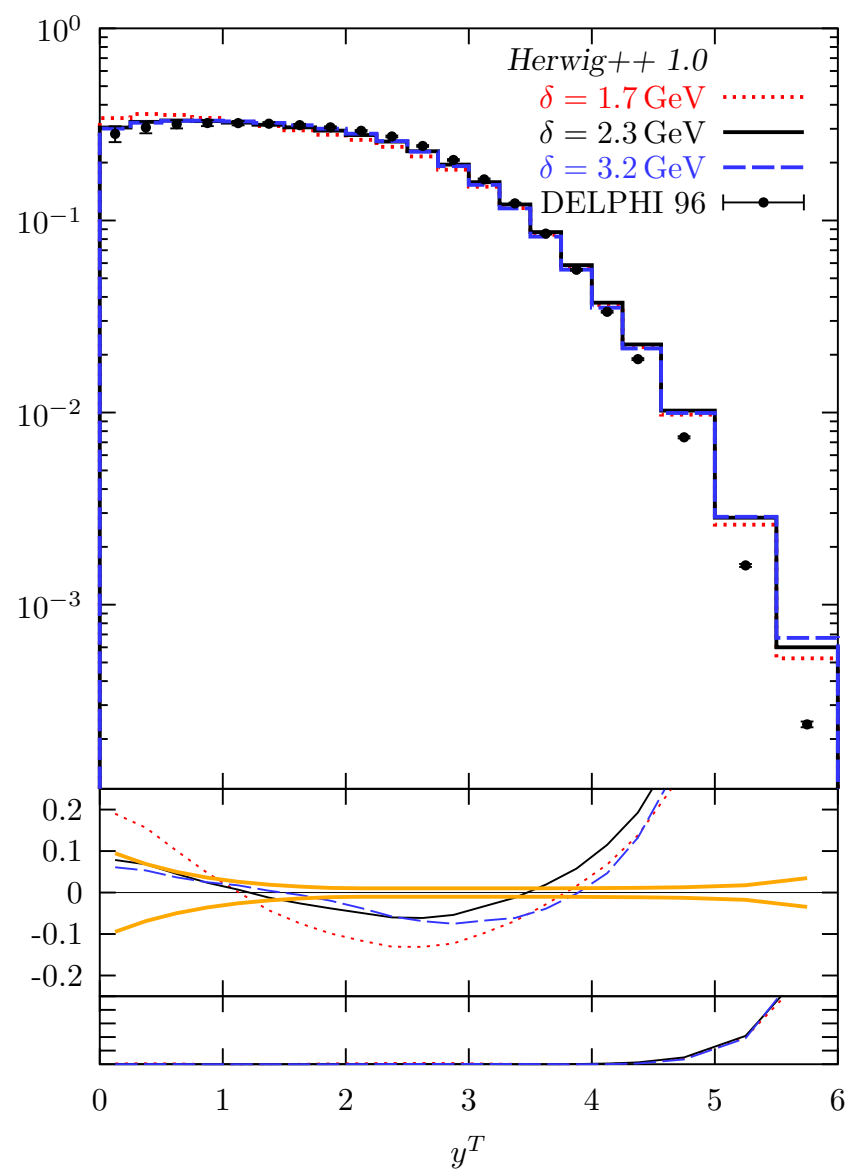
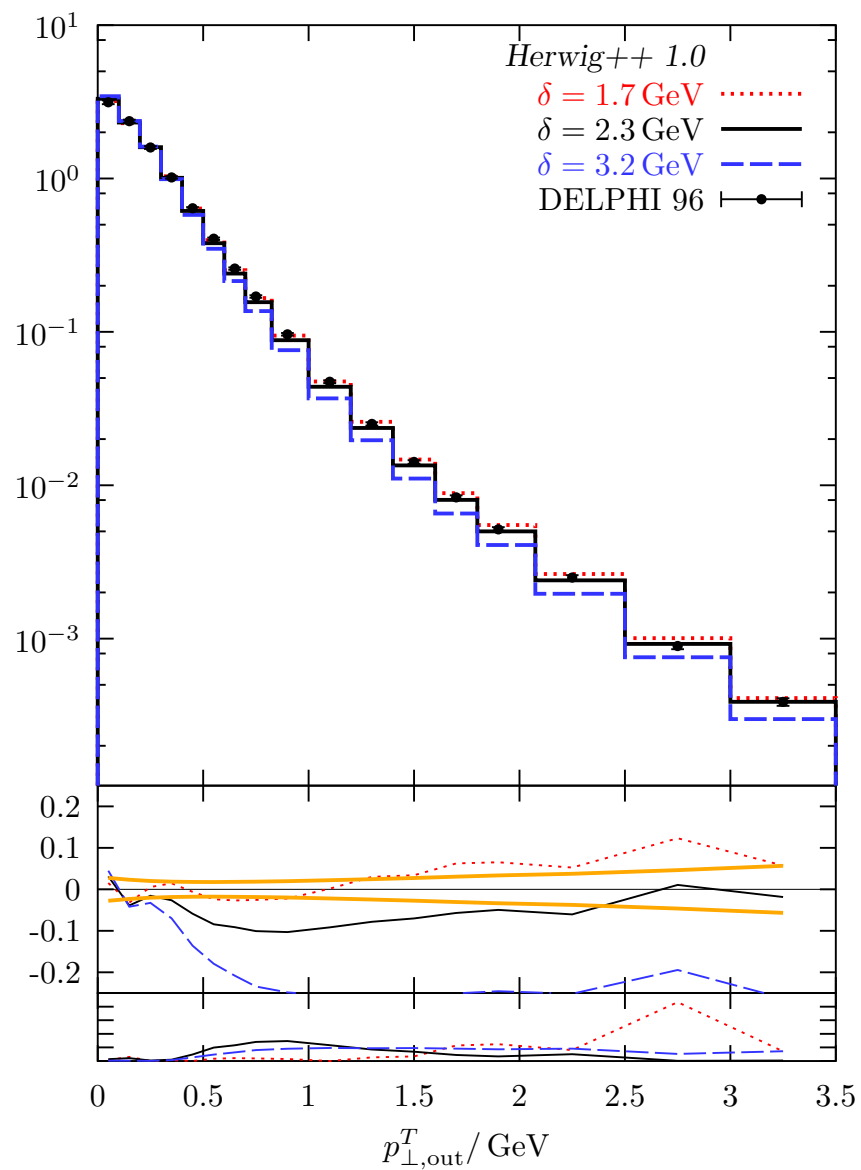
Four Jet Angles



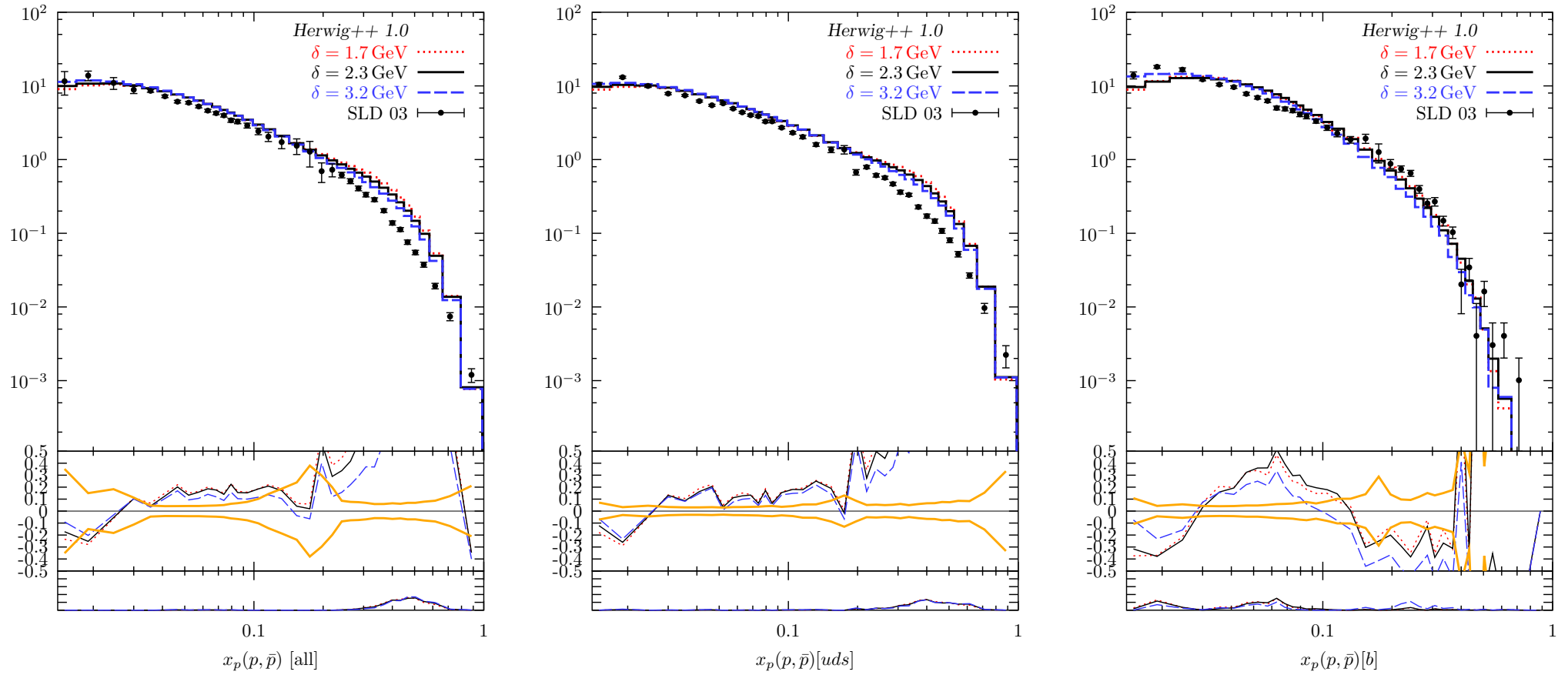
$p_{\perp,\text{in}}^T$ — ME corrections off/on



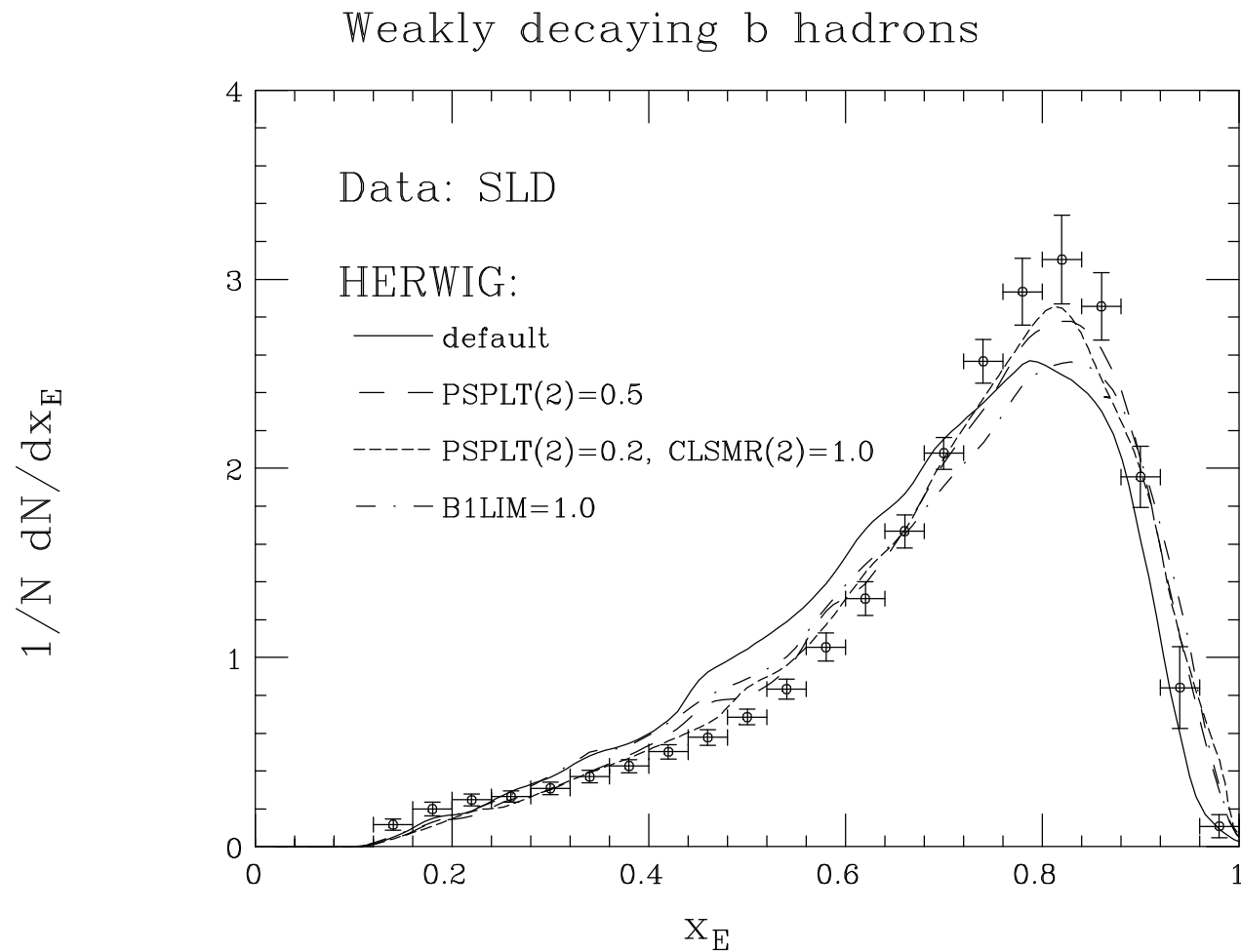
$$p_{\perp,\text{out}}^T, y^T$$



Proton Momentum (all, uds , b)

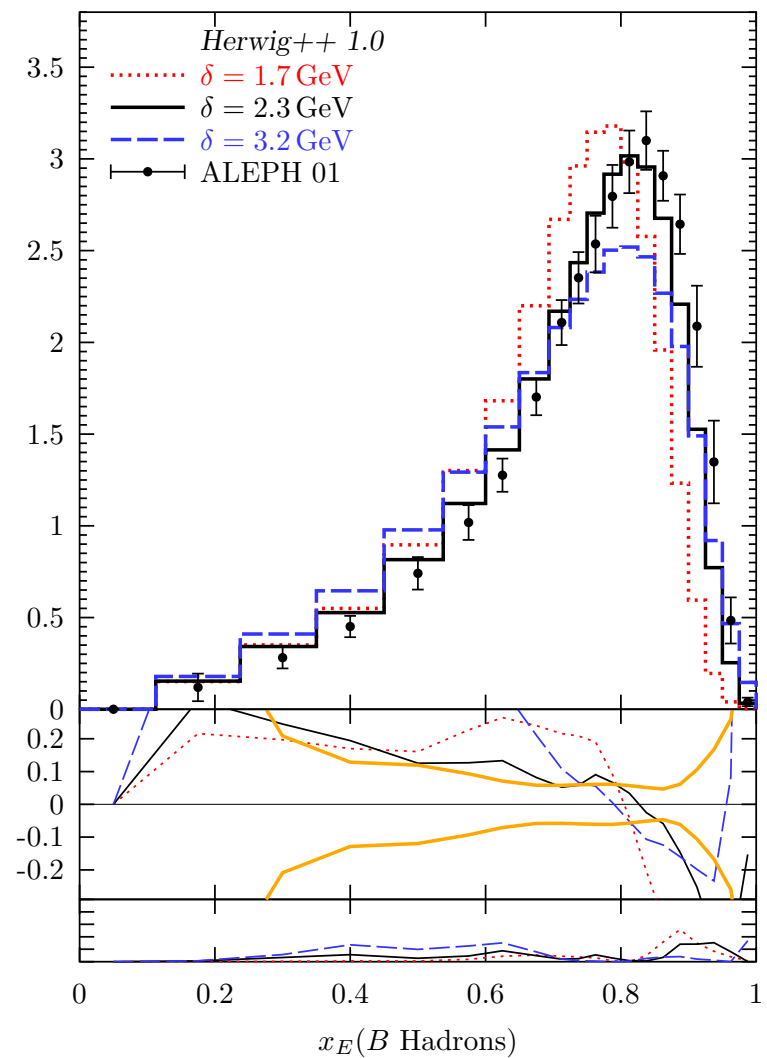
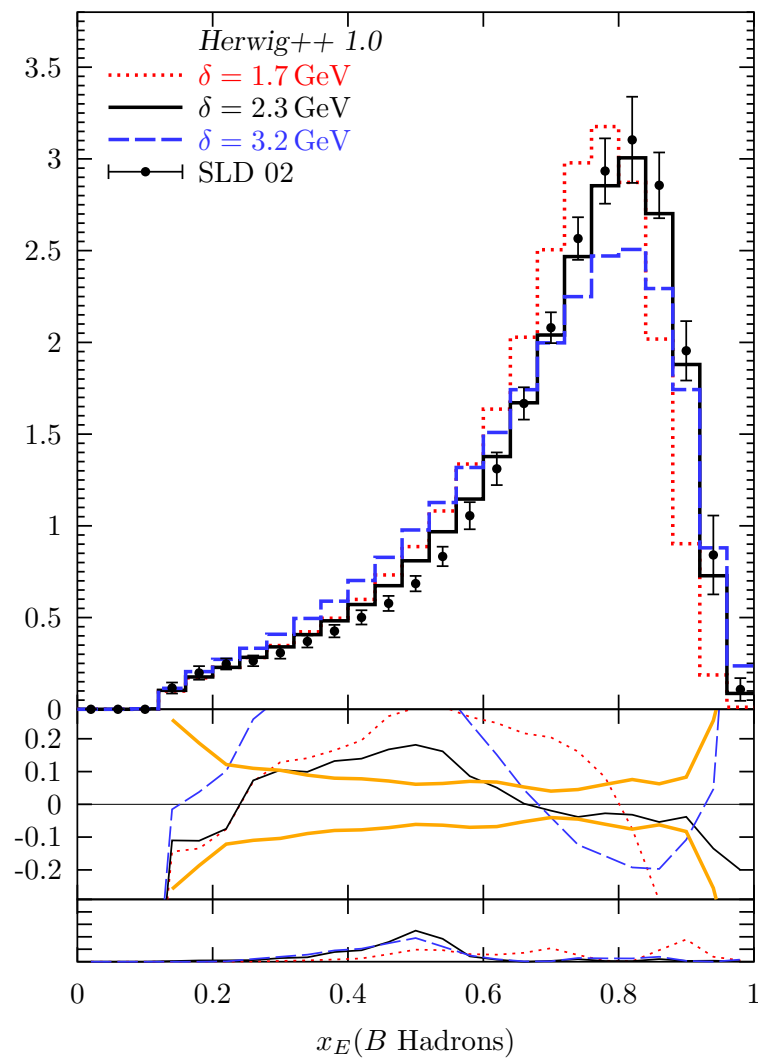


B-fragmentation function



HERWIG 6.4, very sensitive on hadronization!

B -fragmentation function



Only parton shower parameters varied!

Recommendation

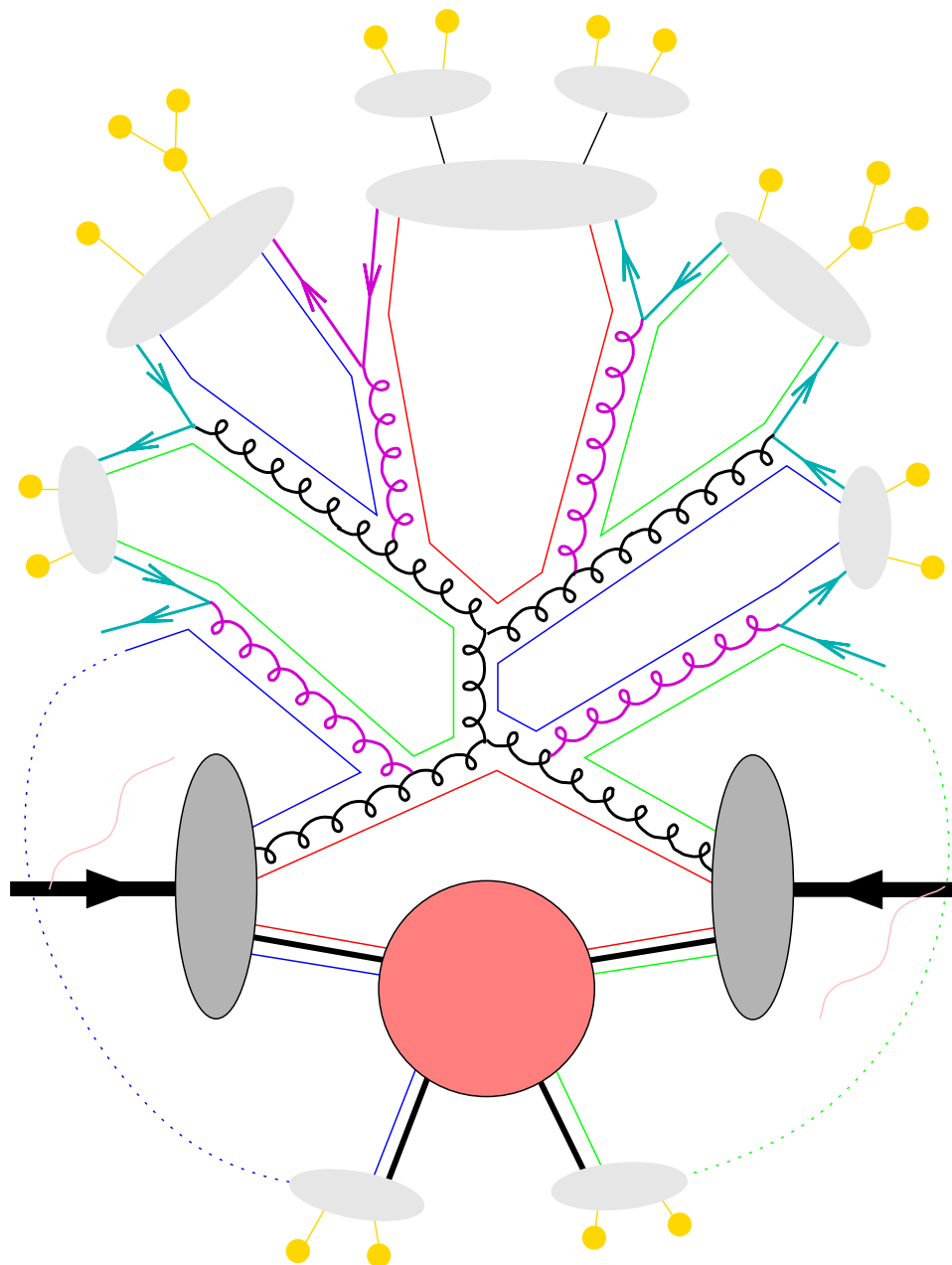
. . . However, the different observables have to be weighted sensibly.

- Low cutoff preferred by event shapes, jet rates, differential jet rates.
- High cutoff preferred by Single Particle distributions along thrust or sphericity axis.
- Either high or low cutoff for y_{nm} .
- High cutoff preferred by Identified particle spectra, particularly for heavy flavour events.
- Intermediate cutoff preferred by B fragmentation function.

We recommend the intermediate value.

Parameter	Default	Initial
$\alpha_s(M_Z)$	0.118	0.114
δ/GeV	2.3	—
m_g/GeV	0.750	—
$Q_{\min}/\text{GeV}$ in $\alpha_s(Q_{\min})$	0.631	—
CIMax/GeV	3.2	3.35
CIPow	2.0	—
PSplt1	1	—
PSplt2	0.33	—
B1Lim	0.0	—
CIDir1	1	—
CIDir2	1	—
CISmr1	0.40	—
CISmr2	0.0	—
Pwt _d	1.0	—
Pwt _u	1.0	—
Pwt _s	0.85	1.0
Pwt _c	1.0	—
Pwt _b	1.0	—
Pwt _{di}	0.55	1.0
Singlet Weight	1.0	—
Decuplet Weight	0.7	1.0

Additional Complications in pp



- + backward parton evolution
- + soft underlying event

What's next?

Near Future. . .

- ★ Initial state shower:
 - Complete implementation and tests.
- ★ Refine e^+e^- :
 - Full CKKW ME+PS matching.
 - Precision tune to LEP data should be possible.
- ★ with IS and FS showers running:
 - we can start to test Drell–Yan and jets in pp collisions.
 - cross check with Tevatron data and finally make predictions for the LHC.
 - Study of **DIS** possible.
- ★ Underlying Event.
- ★ Hadronic Decays: **NEW!** τ -decays, Spin correlations (P Richardson).
- ★ **New Ideas:** soft gluons, improved shower algorithm, NLO, . . .

Schedule?

- Ready for LHC!

Conclusion

We have completed a new event generator for e^+e^- Annihilation:

Herwig++ 1.0

<http://www.hep.phy.cam.ac.uk/theory/Herwig++>