# Lattice investigation of charmed and bottom hadrons

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#### Motivation

### LHCb discovery of excited $\Omega_c^0$ baryons: 2017



Motivation

### A year of heavy baryons (2017) : $\Xi_{cc}$



Aaij et al. (LHCb) 1707.01621

Mattson et al. (SELEX) hep-ex/0208014

Other papers from LHCb on properties of  $\Xi_{cc}$ : 1905.02421; 1807.01919; 1806.02744

Charmed-bottom hadrons from lattice QCD

Universität Regensburg (3 of 27)

Motivation

### Doubly heavy tetraquarks : possibly stable system



Discovery of the Doubly Charmed  $\Xi_{cc}$  Baryon Implies a Stable  $bb\bar{u}\bar{d}$  Tetraquark

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# Heavy Quark Symmetry implies for a sufficiently heavy "heavy" quark, the doubly heavy tetraquarks should be stable to strong decays.



Carlson et al. PRD37 744 (1988); Eichten and Quigg 1707.09575

For a heavy-light meson-like system of doubly charm baryons

$$\frac{M(\Xi_{QQ}^*) - M(\Xi_{QQ})}{M(V_Q) - M(PS_Q)} \rightarrow \frac{3}{4}$$

Brambilla et al hep-ph/0506065

Data : Brown et al 1409.0497

Figure : Lewis CHARM-2018

### Heavy hadrons

- Clean signature for many observables, even though rarity in events.
- A large number of discoveries in heavy hadron sector over the past decade.
  - usual quarkonium mesons like  $h_c(1P)$ ,  $h_b(1P)$ ,  $h_b(2P)$ ,  $\psi(1^3D_2)$ .
  - heavy baryons like  $\Xi_{cc}$  and highly excited states of  $\Omega_c$  (LHCb).
  - new beauty baryons like  $\Sigma_b^{\pm}(6097)$ ,  $\Lambda_b(6146)$  and  $\Lambda_b(6152)$  by LHCb.
  - tetraquarks (LHCb, Belle) and pentaquarks (LHCb).
  - many other states with less theoretical understanding (XYZ).
- Charmed-bottom hadrons : largely unexplored unlike others. Discovered :  $B_c(1S, 0^-)$  at 6275(1) MeV  $B_c(2S, 0^-)$  at 6842(6) MeV (ATLAS)  $B_c(2S, 0^-)$  at 6872(6) MeV and  $B_c(2S, 1^-)$  (LHCb) at 6841(6) MeV\*

### Lattice QCD : theoretical prospects

LQCD : A non-perturbative, gauge invariant regulator for the QCD path integrals.

- Quarks lives on sites
- Gauge fields lives on links
- Lattice spacing : UV cut off
- Lattice size : IR cut off



Discretization  $\Rightarrow$  Finite number of degrees of freedom

 $\Rightarrow$  Infinite dimensional path integrals  $\rightarrow$  finite dimensional integrals.

Employ Monte Carlo importance sampling methods on Euclidean metric for numerical studies.

### QCD spectrum from Lattice QCD

- Aim : to extract the physical states of QCD.
- Euclidean two point current-current correlation functions

$$\mathcal{C}_{ji}(t_f-t_i)=\langle 0|\Phi_j(t_f)ar{\Phi}_i(t_i)|0
angle =\sum_nrac{Z_i^{n*}Z_j^n}{2m_n}e^{-m_n(t_f-t_i)}$$

where  $\Phi_j(t_f)$  and  $\overline{\Phi}_i(t_i)$  are the desired interpolating operators and  $Z_i^n = \langle 0 | \Phi_j | n \rangle$ .

• Effective mass defined as  $\log[\frac{C(t)}{C(t+1)}]$ 



• The ground states : from the exponential fall off at large times. Non-linear fitting techniques.

### LHCb discovery of excited $\Omega_c^0$ baryons: 2017



### Excited baryon spectroscopy

- Aim : Extraction of highly excited states. Local operators  $\rightarrow$  low lying states. Extended operators  $\rightarrow$  Radial and orbital excitations.
- Continuum operators with well defined quantum nos. Reduce/subduce into the irreps of the reduced symmetry.



- Variational analysis of correlation matrices, C<sub>ji</sub>.
- Rigorous spin identification procedure using operator state overlaps, Z.

### Results : $\Omega_c$ spectrum (L1)



MP & Mathur 1704.00259, Charm 2013 (1311.4806), Charm 2015 (1508.07168).

Magenta ellipses : States with strong non-relativistic content. The low lying spectrum same as non-relativistic expectations.

### Experiment vs. lattice predictions (L1)



Expt.

422(1)

395(1)

Lattice

464(20)

409(19)

 $\begin{array}{l} \mbox{Spin 1/2, 3/2, 5/2} \\ \Omega_{ccc}: \ \mbox{HSC 1307.7022} \\ \Xi_{cc} \ \mbox{and} \ \Omega_{cc}: \ \ \mbox{HSC 1502.01845} \end{array}$ 

Here  $\Delta E = E - E_{\Omega_c^0}$ .

The new states correspond to the excited *p*-wave excitations.

MP & Mathur 1704.00259, Charm 2013, Charm 2015.

On anisotropic  $N_f=2+1$  lattices  $L\sim 1.9~fm,~a_tm_c=0.114$  and  $m_\pi=391~MeV$ 

Edwards et al. 0803.3960

Energy

 $\overline{\Delta E}_{\Omega^0_2(3119)}$ 

 $\Delta E_{\Omega^0_c(3090)}$ 

### Heavy quarks on the lattice

Fermion action:

$$ar{\psi}(\gamma.\mathsf{D}+\mathsf{m})\psi$$

Only dimensionless quantities defined on the lattice Lattice fermion action:  $\bar{\psi}(\mathbf{x})(\gamma.\mathbf{D_L} + \mathbf{am})\psi(\mathbf{x})$ 

$$\label{eq:Discretization} \begin{split} \text{Discretization} &\to \textbf{D}_{\textbf{L}} \text{; clover, staggered, overlap, domain-wall, etc.} \\ \text{Discretization errors on observables } \mathcal{O}(\textbf{am}). \end{split}$$

Charm quarks:  $m_c \sim 1.275~{
m GeV}$ 

 $am_c=0.5 \ \Rightarrow \ a\sim 0.075 \ {
m fm}$ 

 $am_c = 0.3 \ \Rightarrow \ a \sim 0.046 \ {
m fm}$ 

Bottom quarks:  $m_b \sim 4.66$  GeV

$$am_b = 0.5 \Rightarrow a \sim 0.021 \text{ fm}$$

$$am_b = 0.3 \Rightarrow a \sim 0.013 \text{ fm}$$

### Another lattice calculation: L2

- State-of-the-art ensembles :  $N_f = 2+1+1$  HISQ (MILC) a = 0.1207(11), 0.0888(8) and 0.0582(5) fm. MILC Collaboration 1212.4768
- Chiral fermion action for quark masses from light to charm. Exact chiral symmetry at finite lattice spacing. No O(am) errors.
- Heavy quark mass tuning Dowdall *et al.* 1110.6887  $\Delta E_{hfs}^{1S,\bar{c}c} = 115(2)(3) \text{ MeV } (Lattice) \& \Delta E_{hfs}^{1S,\bar{c}c} = 114 \text{ MeV } (Expt).$
- We work with energy splittings

$$\Delta M_H = [M_H^L - n_c \overline{1S}_c/2] a^{-1}.$$

- Continuum extrapolation fit forms  $Q^f = A + a^2 B$ ,  $L^f = A + a^2 log(a)B$ .
- We extract the mass of hadrons from

$$M_H^c = \Delta M_H^c + n_c (\overline{1S}_c)_{phys}/2.$$

### Comparison between two lattice determinations



MP & Mathur (HSC) 1704.00259

Charmed-bottom hadrons from lattice QCD

Doubly charmed baryons

### A year of heavy baryons (2017) : $\Xi_{cc}$



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Charmed-bottom hadrons from lattice QCD

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### $\Xi_{cc}$ from lattice QCD



### $\Xi_{cc}$ isospin splitting (LQCD), 2.16(11)(17) MeV : BMW 1406.4088 SELEX measurement (3519 MeV) : Mattson *et al.* hep-ex/0208014

Charmed-bottom hadrons from lattice QCD

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Doubly charmed baryons

### Doubly charmed baryons from lattice QCD



Figure from arXiv:1905.09651 [hep-lat]

 $\Omega_{cc}$ : probably the next doubly charm baryon to be observed in LHCb.  $\Xi_{cc}^*$ : radiative decay might be predominant over strong decays. So might be difficult for LHCb to discover.

### Singly charmed baryons from lattice QCD



Figure from arXiv:1905.09651 [hep-lat]

### Charmed-bottom hadrons from lattice QCD: L2

- State-of-the-art ensembles :  $N_f = 2+1+1$  HISQ (MILC) 24<sup>3</sup> × 64, 32<sup>3</sup> × 96 and 48<sup>3</sup> × 144 a = 0.1207(11), 0.0888(8) and 0.0582(5) fm. MILC Collaboration 1212.4768
- Chiral fermion action for quark masses from light to charm. Exact chiral symmetry at finite lattice spacing. No O(am) errors.
- Bottom quarks realized using NRQCD formulation. Includes terms in the Hamiltonian through  $\mathcal{O}(\alpha \nu_b^4)$ . HPQCD Collaboration 1110.6887, Hammant *et al* 1303.3234.
- Heavy quark mass tuning Dowdall *et al.* 1110.6887  $\Delta E_{hfs}^{1S,\bar{c}c} = 115(2)(3) \text{ MeV and } \Delta E_{hfs}^{1S,\bar{b}b} = 63(3)(5) \text{ MeV} \quad \text{Lattice}$   $\Delta E_{hfs}^{1S,\bar{c}c} = 114 \text{ MeV and } \Delta E_{hfs}^{1S,\bar{b}b} = 62(1) \text{ MeV} \quad \text{Expt.}$

### Charmed-bottom hadrons from lattice QCD: L2

• We work with energy splittings and dimensionless ratios

$$\Delta M_H = [M_H^L - n_b \overline{1S}_b/2 - n_c \overline{1S}_c/2] a^{-1} \quad \text{and} \\ R_H = \frac{M_H^L - n_b \overline{1S}_b/2}{M_{B_c(0^-)}^L - n_b \overline{1S}_b/2}.$$

- Chiral extrapolation  $(\Xi_{bc})$  using  $A + m_{\pi}^2 B$  and other fancy chiral extrapolation forms as in Brown *et al* 1409.0497.
- Continuum extrapolation fit forms  $Q^f = A + a^2 B$ ,  $C^f = A + a^3 B$ ,  $L^f = A + a^2 log(a)B$ .
- We extract the mass of hadrons from

$$\begin{aligned} M_H^c &= \Delta M_H^c + n_b (\overline{1S}_b)_{phys}/2 + n_c (\overline{1S}_c)_{phys}/2 & \text{and} \\ M_H^c &= R_H^c \times (M_{B(0^-)} - n_b \overline{1S}_b/2)_{phys} + n_b (\overline{1S}_b)_{phys}/2. \end{aligned}$$

### $B_c$ meson results (L2)

Hadrons	Lattice	HPQCD	Wurtz <i>et al</i>	Experiment
$B_{c}(0^{-})$	6276(3)(6)	6278(9)	-	6274.9(8)
$B_{c}^{*}(1^{-})$	6331(4)(6)	6332(9)	6332.5(3)	?
$B_{c}(0^{+})$	6712(18)(7)	6707(17)	6711(2)	?
$B_{c}(1^{+})$	6736(17)(7)	6742(16)	6752(2)	?

Mathur, MP, Mondal 1806.04151; Dowdall et al. (HPQCD) 1207.5149; Wurtz et al. 1505.04410



### bc baryons: L2 Vs other lattice results

Hadrons	Lattice	Brown <i>et al</i>	Experiment
$\Xi_{cb}(cbu)(1/2^+)$	6945(22)(14)	6943(33)(28)	?
$\Xi_{cb}^{\prime}(cbu)(1/2^+)$	6966(23)(14)	6959(36)(28)	?
$\Xi_{cb}^{*}(cbu)(3/2^{+})$	6989(24)(14)	6985(36)(28)	?
$\Omega_{cb}(cbs)(1/2^+)$	6994(15)(13)	6998(27)(20)	?
$\Omega_{cb}^{\prime}(cbs)(1/2^+)$	7045(16)(13)	7032(28)(20)	?
$\Omega^*_{cb}(cbs)(3/2^+)$	7056(17)(13)	7059(28)(21)	?
$\Omega_{ccb}(1/2^+)$	8005(6)(11)	8007(9)(20)	?
$\Omega^*_{ccb}(3/2^+)$	8026(7)(11)	8037(9)(20)	?
$\Omega_{cbb}(1/2^+)$	11194(5)(12)	11195(8)(20)	?
$\Omega^*_{cbb}(3/2^+)$	11211(6)(12)	11229(8)(20)	?

Brown *et al* 1409.0497

Mathur, MP, Mondal 1806.04151

### bc baryons: L2 Vs other lattice results



Brown et al 1409.0497

Mathur, MP, Mondal 1806.04151

Charmed-bottom hadrons from lattice QCD

### Summary in predictions for *bc* hadrons (L2)



Mathur, MP, Mondal 1806.04151

### Doubly heavy tetraquarks : possibly stable system



Discovery of the Doubly Charmed  $\Xi_{cc}$  Baryon Implies a Stable  $bb\bar{u}\bar{d}$  Tetraquark

Marek Karliner1,\* and Jonathan L. Rosner2,†

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For a heavy-light meson-like system of doubly charm baryons

$$\frac{M(\Xi_{QQ}^*) - M(\Xi_{QQ})}{M(V_Q) - M(PS_Q)} \rightarrow \frac{3}{4}$$

Brambilla et al hep-ph/0506065

Data : Brown et al 1409.0497

Figure : Lewis CHARM-2018

## $\bar{b}\bar{b}ud$ and $\bar{b}\bar{b}ls$ tetraquarks on lattice

Francis et al. 1607.05214



Study of pseudovector doubly bottom tetraquarks on  $N_f = 2 + 1$  PACS-CS lattices Correlation matrices out of tetraquark as well as meson-meson interpolators Binding energies :  $\Delta E_{\bar{b}\bar{b}ud} = 189(10)$  MeV and  $\Delta E_{\bar{b}\bar{b}us} = 98(7)$  MeV

Other existing lattice calculations :

Bicudo et al. 1510.03441, Leskovec et al. 1904.04197

### Doubly heavy tetraquarks from L2



Junnarkar, Mathur & MP (ILGTI) 1810.12285

- Extensive study of quark mass dependence on binding energies of doubly heavy tetraquarks.
- Correlation matrices out of tetraquark as well as meson-meson interpolators.
- Binding energies from difference in the lowest non-interacting levels.

### Summary doubly heavy tetraquarks $(J^P = 1^+)$



1810.12285 (ILGTI, This work), 1607.05214 (Francis),

1707.09575 (Eichten and Quigg) 1707.07666 (Karliner and Rosner),

1510.03441 (Bicudo), 1709.01417 (HSC, Cheung), 1904.04197 (Leskovec).

### Summary

- Discussed exploratory investigation of excited  $\Omega_c$  baryon spectrum. Highlighted predictions for the quantum numbers of LHCb discovered excited  $\Omega_c$  baryons.
- Summarized lattice QCD predictions for singly and doubly charm baryons. Compared lattice results with experimental values.
- Presented precision lattice QCD predictions/postdictions for masses of charmed-bottom hadron ground states. Predictions compared with other existing lattice estimates and postdictions compared with the respective experimental values and other lattice results.
- Beyond conventional hadrons : Promising platforms to study. Doubly heavy tetra-quarks: possibly deeply bound states. Discussed our study on  $m_q$  dependence of these binding energies. Ongoing investigations for  $bc\bar{u}\bar{d}$  tetraquarks.

Thank you...

## Spin identification : $J > \frac{3}{2}$ (L1)

- For example, a continuum operator  $O = [ccc \otimes (\frac{3}{2}^+)_S^1 \otimes D_{L=2,S}^{[2]}]^{J=\frac{5}{2}}$ . Projects on to  $\frac{5}{2}^+$ .
- In the continuum,  $\langle 0|O|\frac{5}{2}^+\rangle = Z$ .
- On lattice, O gets subduced over two lattice irreps  $H_g$  and  $G_{2g}$ .
- Then

$$\langle 0|O_{H_g}|^{5^+}_2 \rangle = Z_1 \alpha$$
 &  $\langle 0|O_{G_{2g}}|^{5^+}_2 \rangle = Z_2 \beta$ 

where  $\alpha$  and  $\beta$  are the Clebsch-Gordan coefficients.

• If "close" to the continuum, then  $Z ~\sim~ Z_1 ~\sim~ Z_2.$ 

#### Summarv

### Overlap factors (Z) across multiple irreps : $5/2^+$ (L1)



### Errors in predictions for *bc* hadrons (L2)

- Wall sources to reduce statistical errors. All well below percent level.
- Discretization errors : Largest found to be in  $\Xi_{bc}$ ,  $\sim$  6 MeV.
- Scale setting : Independently calibrate the lattice using  $\Omega$  baryon. This work : a = 0.1192(14), 0.0877(10) and 0.0582(5) fm. MILC( $r_1$ ) : a = 0.1207(11), 0.0888(8) and 0.0582(5) fm. Largest errors to be  $\sim 6$  MeV.
- Chiral extrapolation : Robust with different extrapolation forms.
- Uncertainties in NRQCD : missing higher order terms,  $O(\alpha^2 \nu^4)$  and  $O(\alpha \nu^6)$ . ~ 4 MeV for mesons
  - $\sim$  5, 5 and 6 MeV for *bcq*, *bcc* and *bbc* baryons respectively.
- Other sources : Quark mass mis-tuning, unphysical sea quark mass effects, electromagnetism, isospin breaking effects and absence of dynamical bottom quarks collectively to be within few MeV.

Mathur, MP, Mondal 1806.04151

Summary

### $\Xi_{bc}$ baryon chiral extrapolation (L2)



On the coarse lattice. Chiral extrapolation form  $A + m_{\pi}^2 B$ .

Mathur, MP, Mondal 1806.04151

Charmed-bottom hadrons from lattice QCD