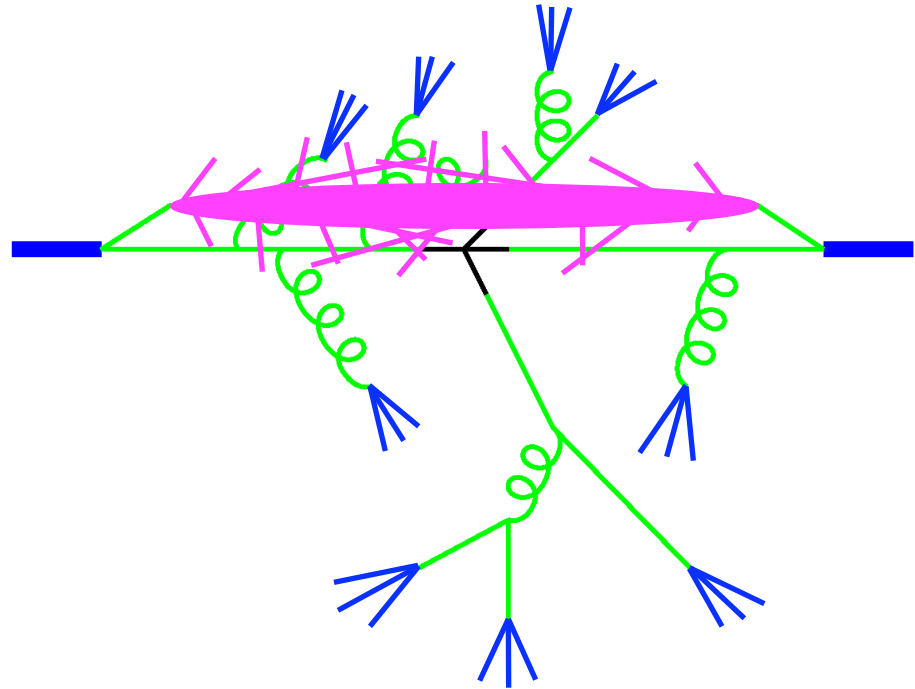


# Monte Carlo Methods in Particle Physics

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IMPRS, Munich  
19-23 November 2007

# Structure of LHC Events

1. Hard process
2. Parton shower
3. Hadronization
4. Underlying event



# Lecture 3: Hadronization

Partons are not physical particles: they cannot freely propagate.

Hadrons are.

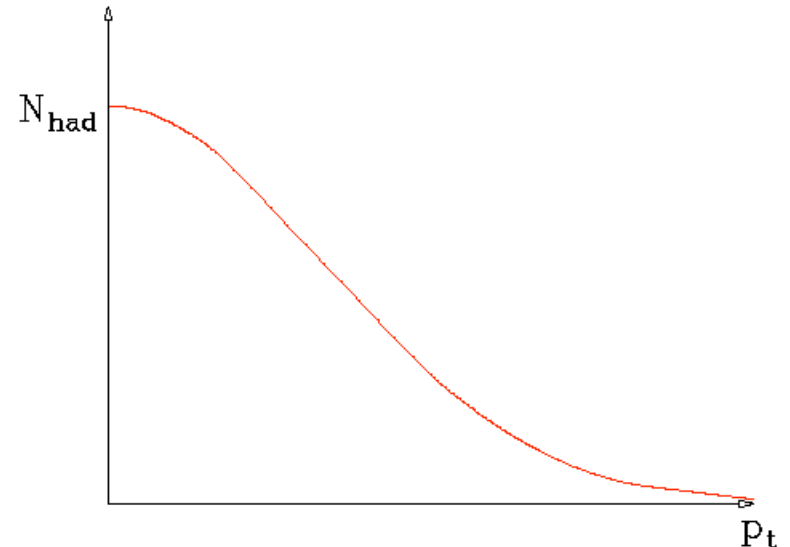
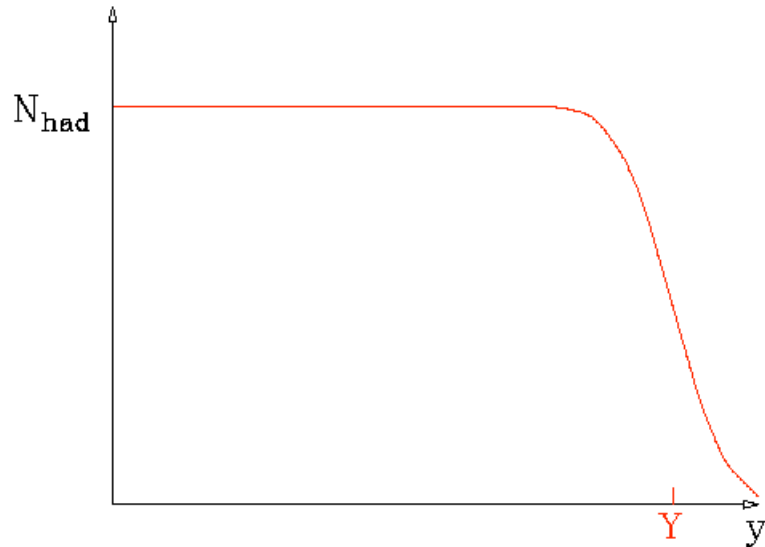
Need a model of partons' confinement into hadrons: hadronization.

1. Phenomenological models.
2. Confinement.
3. The string model.
4. Preconfinement.
5. The cluster model.
6. Underlying event models.

# Phenomenological Models

Experimentally,  $e^+e^- \rightarrow$  two jets:

Flat rapidity plateau and limited  $p_t$ ,  $\rho(p_t^2) \sim e^{-p_t^2/2p_0^2}$



# Estimate of Hadronization Effects

Using this model, can estimate hadronization correction to perturbative quantities.

Jet energy and momentum:

$$E = \int_0^Y dy d^2 p_t \rho(p_t^2) p_t \cosh y = \lambda \sinh Y$$

$$P = \int_0^Y dy d^2 p_t \rho(p_t^2) p_t \sinh y = \lambda(\cosh Y - 1) \sim E - \lambda,$$

with  $\lambda = \int d^2 p_t \rho(p_t^2) p_t$ , mean transverse momentum.

Estimate from Fermi motion  $\lambda \sim 1/R_{had} \sim m_{had}$ .

Jet acquires non-perturbative mass:  $M^2 = E^2 - P^2 \sim 2\lambda E$

Large:  $\sim 10$  GeV for 100 GeV jets.

# Independent Fragmentation Model (“Feynman—Field”)

Direct implementation of the above.

Longitudinal momentum distribution = arbitrary fragmentation function: parameterization of data.

Transverse momentum distribution = Gaussian.

Recursively apply  $q \rightarrow q' + \text{had.}$

Hook up remaining soft  $q$  and  $\bar{q}$ .

Strongly frame dependent.

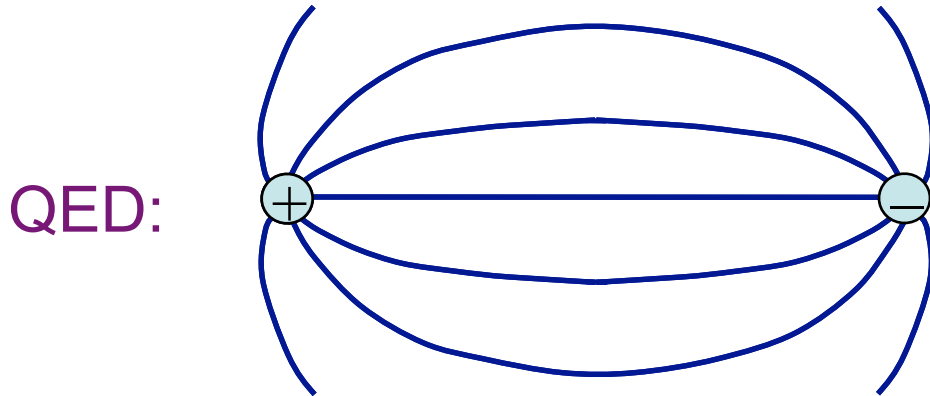
No obvious relation with perturbative emission.

Not infrared safe.

Not a model of confinement.

# Confinement

Asymptotic freedom:  $q\bar{q}$  becomes increasingly QED-like at short distances.



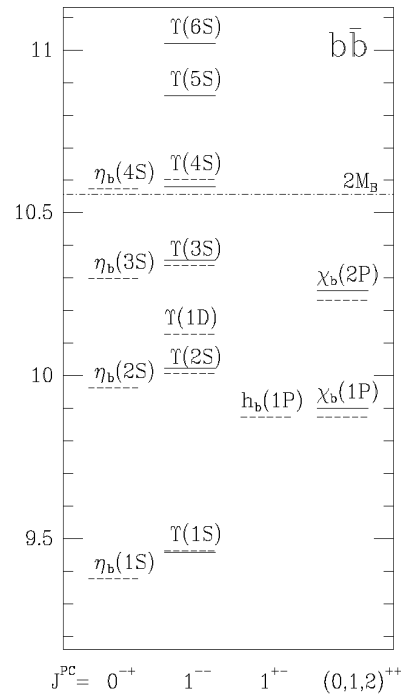
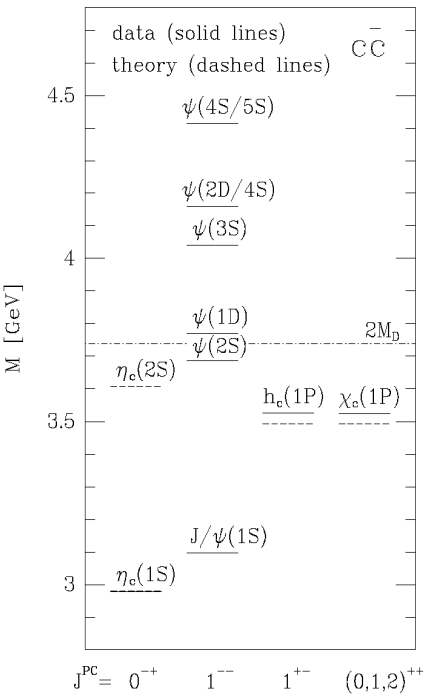
but at long distances, gluon self-interaction makes field lines attract each other:



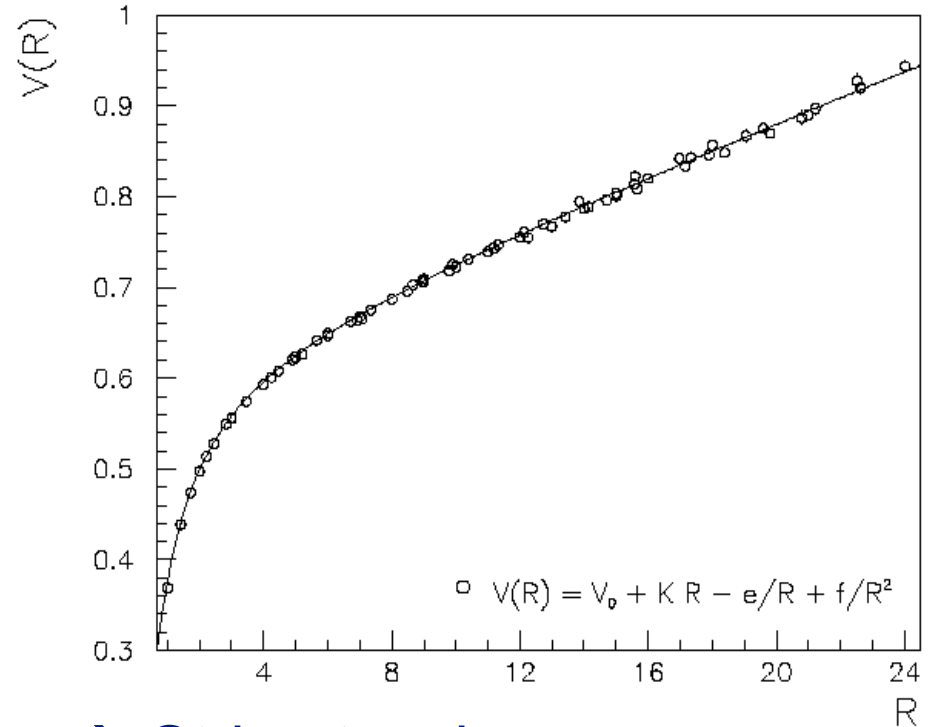
→ linear potential → confinement

# Interquark potential

Can measure from  
quarkonia spectra:



or from lattice QCD:



→ String tension

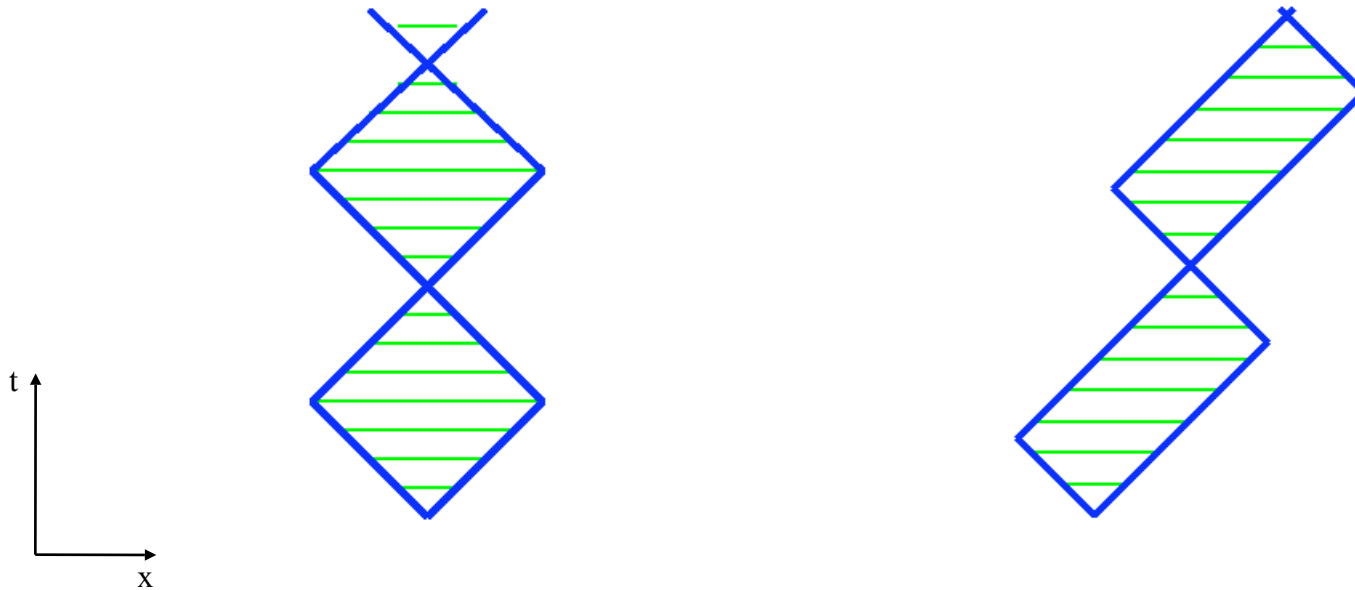
$$\kappa \approx 1 \text{ GeV/fm.}$$



# String Model of Mesons

Light quarks connected by string.

$L=0$  mesons only have 'yo-yo' modes:



Obeys area law:  $m^2 = 2\kappa^2$  area

# The Lund String Model

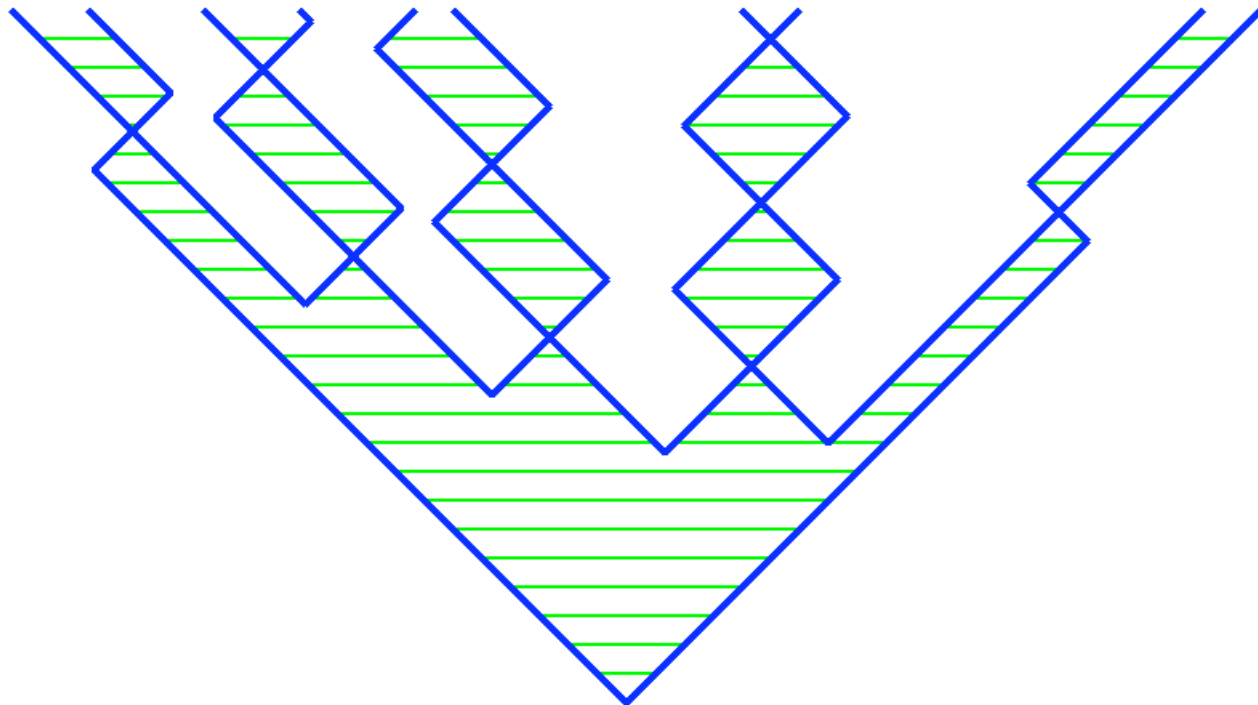
Start by ignoring gluon radiation:

$e^+e^-$  annihilation = pointlike source of  $q\bar{q}$  pairs

Intense chromomagnetic field within string  $\rightarrow q\bar{q}$  pairs created by tunnelling. Analogy with QED:

$$\frac{d(\text{Probability})}{dx dt} \propto \exp(-\pi m_q^2/\kappa)$$

Expanding string breaks into mesons long before yo-yo point.



# Lund Symmetric Fragmentation Function

String picture  $\rightarrow$  constraints on fragmentation function:

- Lorentz invariance
- Acausality
- Left—right symmetry

$$f(z) \propto z^{a_\alpha - a_\beta - 1} (1 - z)^{a_\beta}$$

$a_{\alpha,\beta}$  adjustable parameters for quarks  $\alpha$  and  $\beta$ .

Fermi motion  $\rightarrow$  Gaussian transverse momentum.

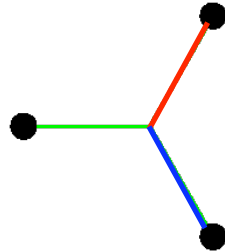
Tunnelling probability becomes

$$\exp \left[ -b(m_q^2 + p_t^2) \right]$$

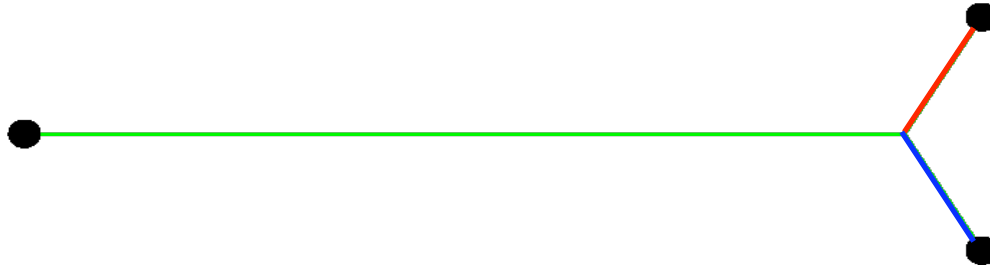
$a, b$  and  $m_q^2$  = main tuneable parameters of model

# Baryon Production

Baryon pictured as three quarks attached to a common centre:



At large separation, can consider two quarks tightly bound: diquark



→ diquark treated like antiquark.

Two quarks can tunnel nearby in phase space: baryon—antibaryon pair  
Extra adjustable parameter for each diquark!

Alternative “popcorn” model:

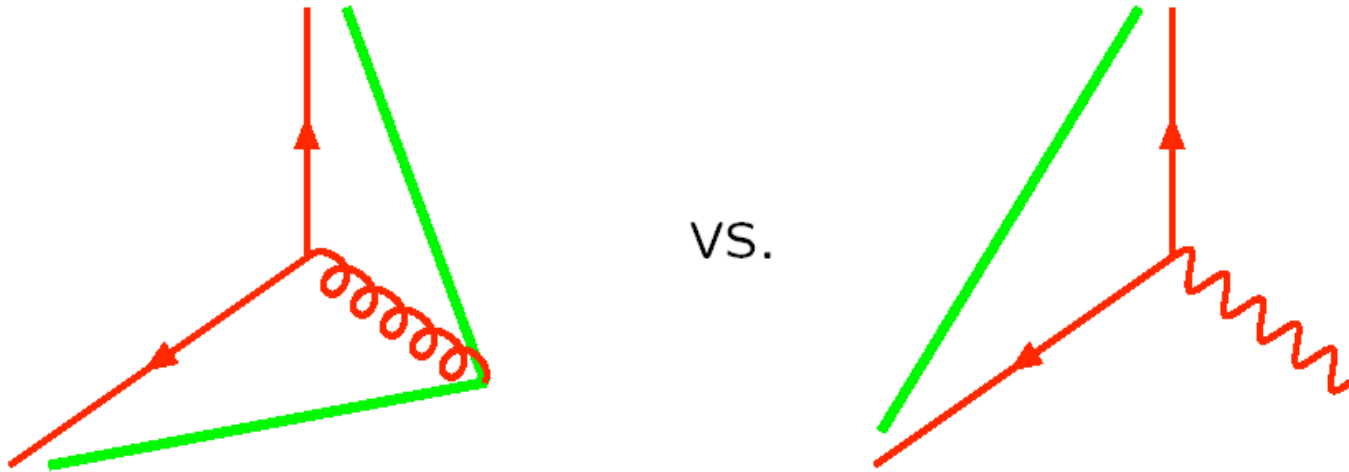


# Three-Jet Events

So far: string model = motivated, constrained independent fragmentation!

New feature: universal

Gluon = kink on string  $\rightarrow$  the string effect



Infrared safe matching with parton shower: gluons with

$k_{\perp} < \text{inverse string width}$  irrelevant.

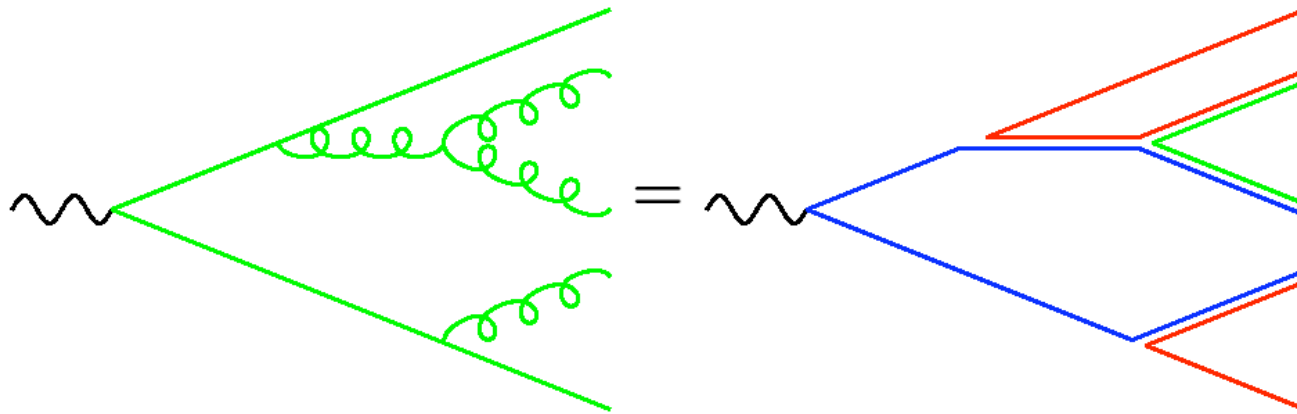
# String Model Summary

- String model strongly physically motivated.
- Very successful fit to data.
- Universal: fitted to  $e^+e^-$ , little freedom elsewhere.
- How does motivation translate to prediction?  
~ one free parameter per hadron/effect!
- Blankets too much perturbative information?
- Can we get by with a simpler model?

# Preconfinement

Planar approximation: gluon = colour—anticolour pair.

Follow colour structure of parton shower: colour-singlet pairs end up close in phase space



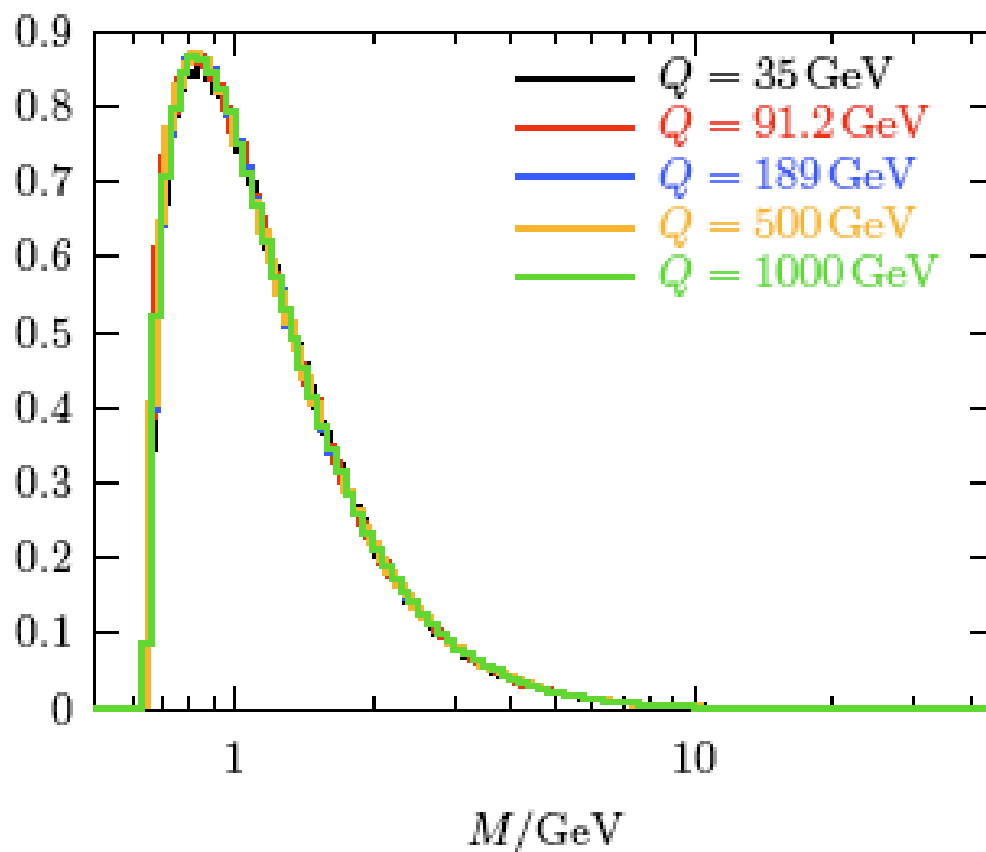
Mass spectrum of colour-singlet pairs asymptotically independent of energy, production mechanism, ...

Peaked at low mass  $\sim Q_0$ .

# Cluster mass distribution

- Independent of shower scale  $Q$ 
  - depends on  $Q_0$  and  $\Lambda$

Primary Light Clusters





# The Naïve Cluster Model

Project colour singlets onto continuum of high-mass mesonic resonances (=clusters). Decay to lighter well-known resonances and stable hadrons.

Assume spin information washed out:  
decay = pure phase space.

→ heavier hadrons suppressed

→ baryon & strangeness suppression 'for free' (i.e. untuneable).

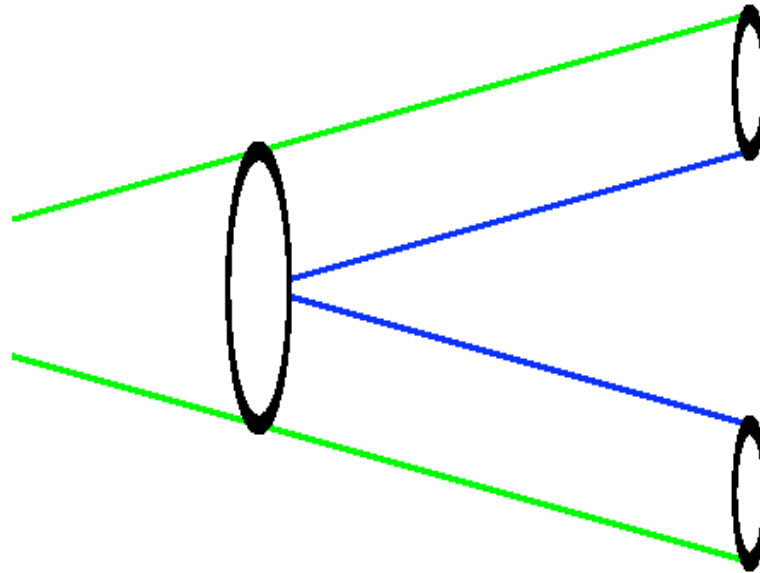
Hadron-level properties fully determined by cluster mass spectrum, i.e. by perturbative parameters.

Shower cutoff  $Q_0$  becomes parameter of model.

# The Cluster Model

Although cluster mass spectrum peaked at small  $m$ , broad tail at high  $m$ .

“Small fraction of clusters too heavy for isotropic two-body decay to be a good approximation” → Longitudinal cluster fission:



Rather string-like.

Fission threshold becomes crucial parameter.

~15% of primary clusters get split but ~50% of hadrons come from them.

# The Cluster Model

“Leading hadrons are too soft”

→ ‘perturbative’ quarks remember their direction somewhat

$$P(\theta^2) \sim \exp(-\theta^2/2\theta_0^2)$$

Rather string-like.

Extra adjustable parameter.

# Strings

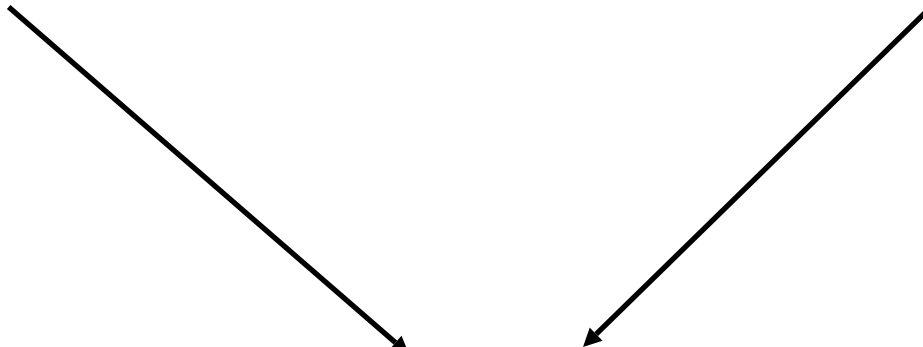
“Hadrons are produced by hadronization: you must get the non-perturbative dynamics right”

Improving data has meant successively refining perturbative phase of evolution...

# Clusters

“Get the perturbative phase right and any old hadronization model will be good enough”

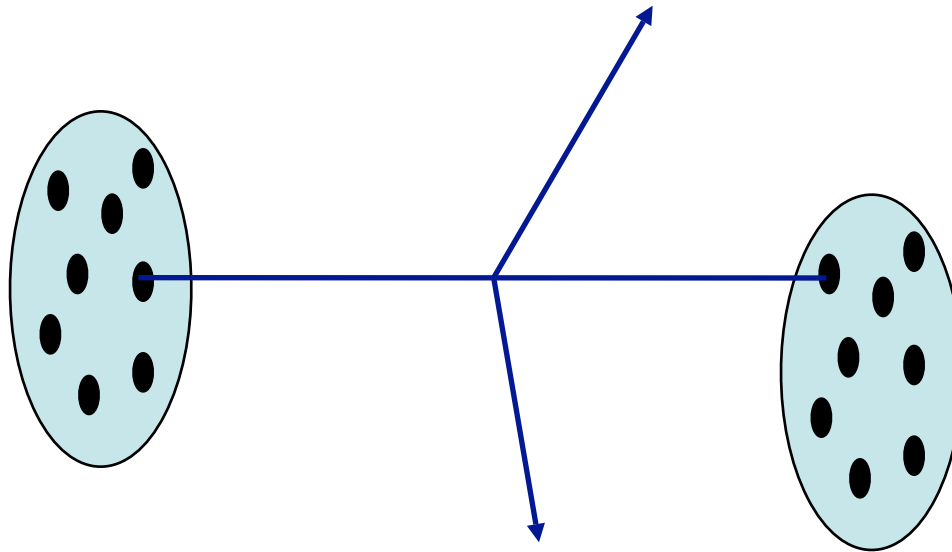
Improving data has meant successively making non-perturbative phase more string-like...



???

# The Underlying Event

- Protons are extended objects
- After a parton has been scattered out of each, what happens to the remnants?

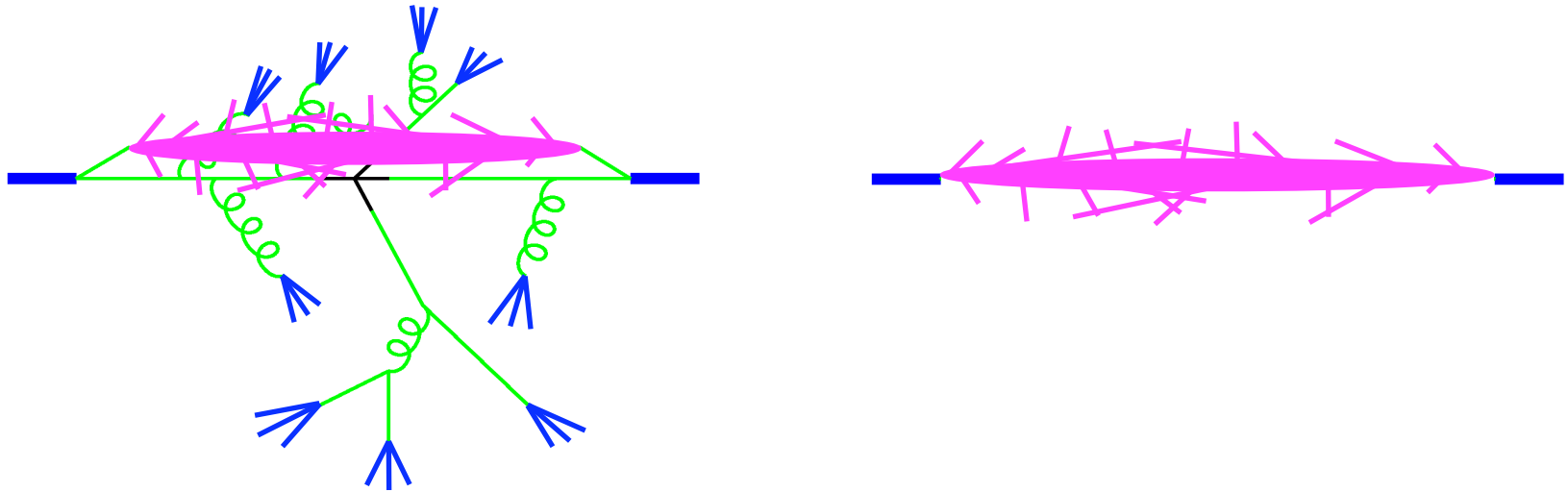


## Two models:

- **Non-perturbative:** Soft parton—parton cross section is so large that the remnants always undergo a soft collision.
- **Perturbative:** ‘Hard’ parton—parton cross section huge at low  $p_t$ , high energy, dominates inelastic cross section and is calculable.

# Soft Underlying Event Model (HERWIG)

Compare underlying event with ‘minimum bias’ collision  
(‘typical’ inelastic proton—proton collision)

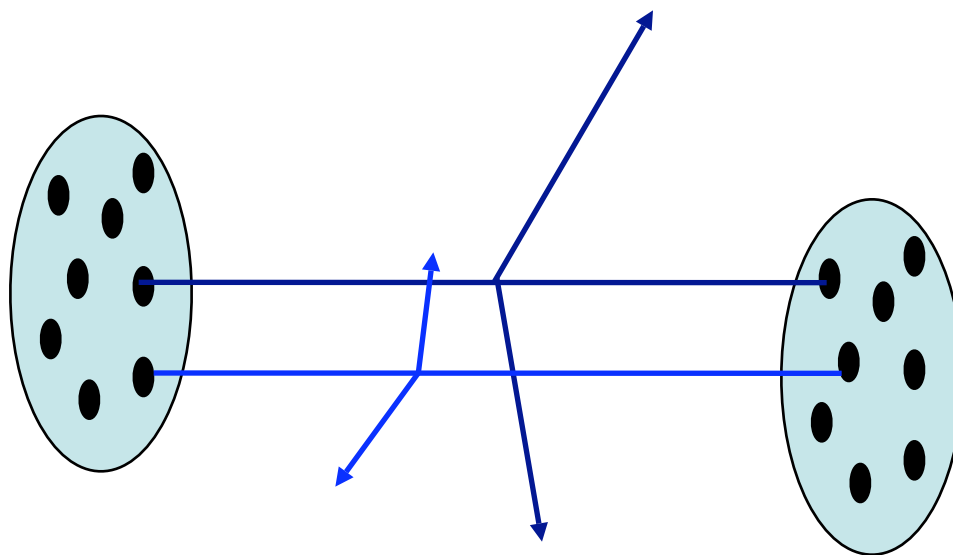


Parametrization of (UA5) data + model of energy dependence

# Multiparton Interaction Model (PYTHIA/JIMMY)

For small  $p_{t\min}$  and high energy inclusive parton—parton cross section is larger than total proton—proton cross section.

→ More than one parton—parton scatter per proton—proton



Need a model of spatial distribution within proton

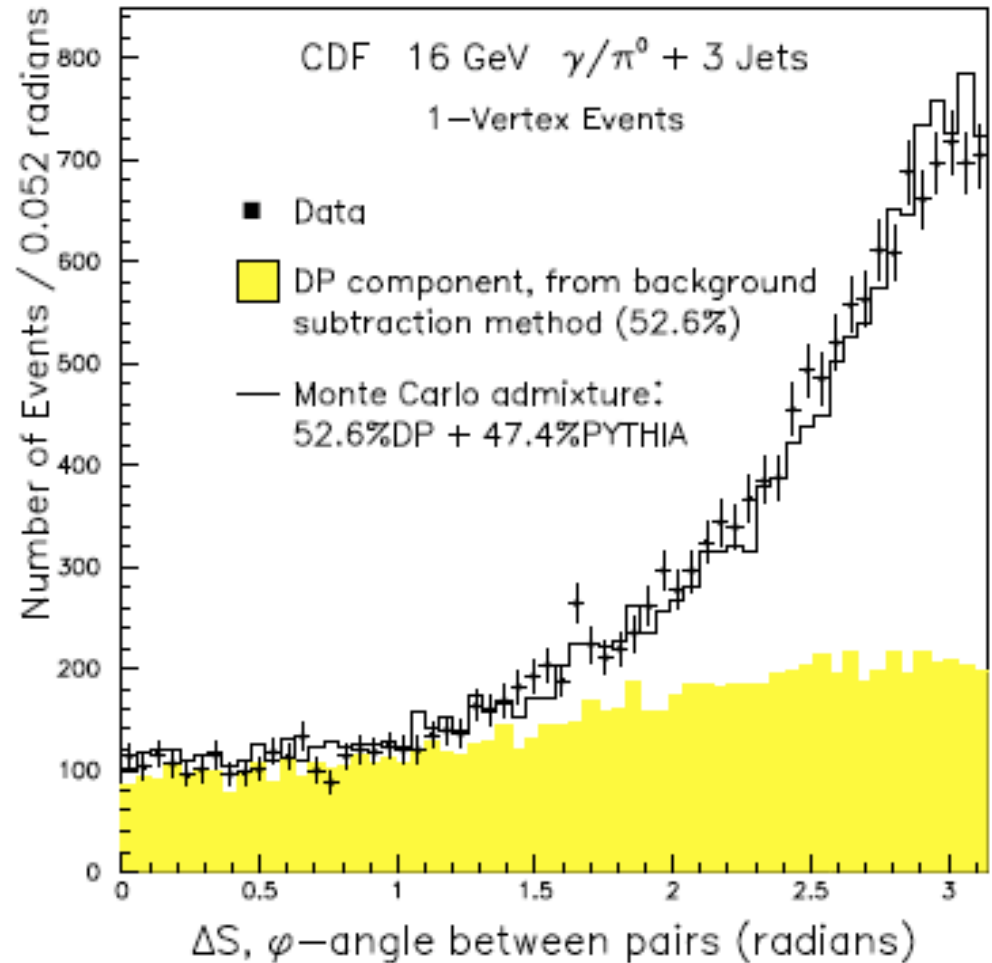
→ Perturbation theory gives you n-scatter distributions

# Double Parton Scattering

- CDF Collaboration,  
PR D56 (1997) 3811

$$\sigma_{\text{DPS}} = \frac{\sigma_{\gamma j} \sigma_{jj}}{\sigma_{\text{eff}}}$$

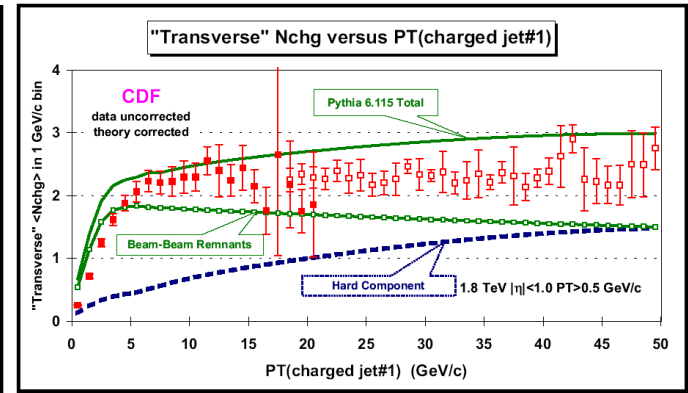
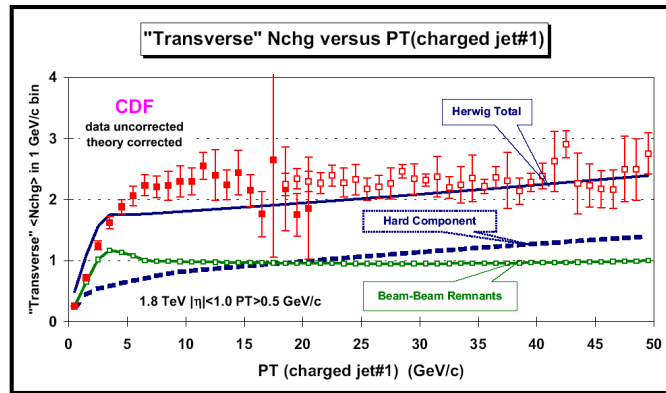
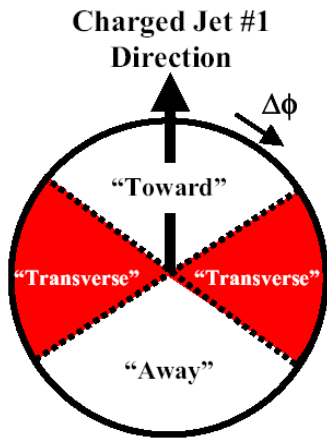
$$\sigma_{\text{eff}} = 14 \pm 1.7^{+1.7}_{-2.3} \text{ mb}$$





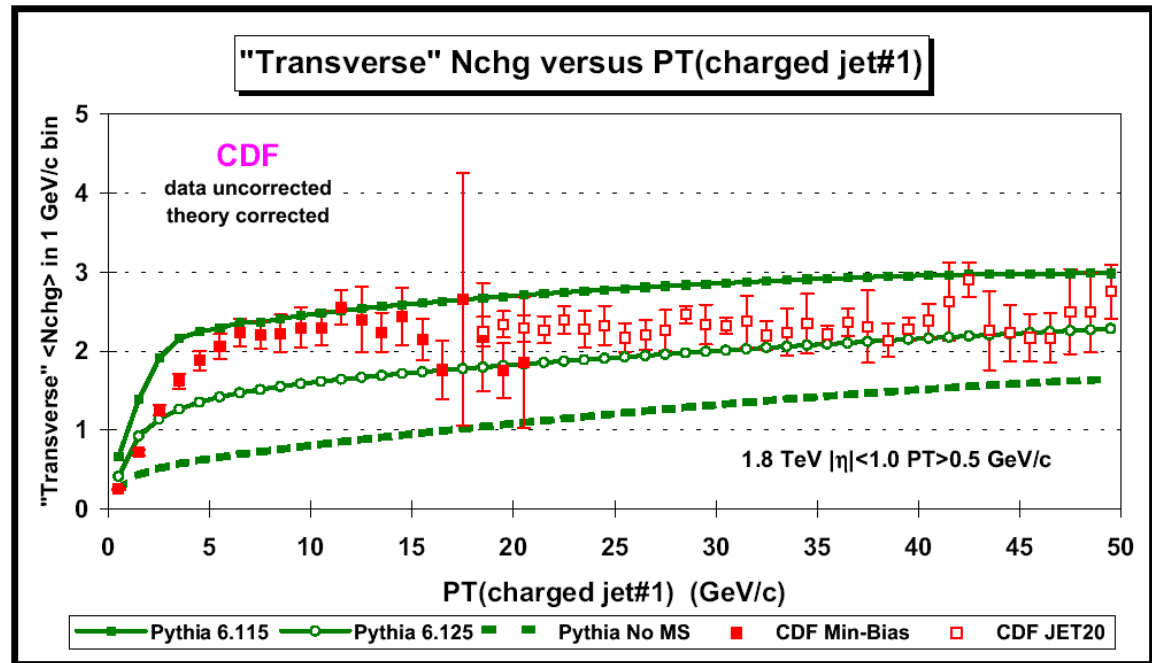
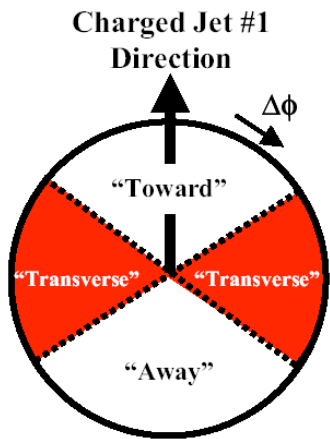
# Some Warnings

- Not everyone means same thing by “underlying event”
  - Remnant—remnant interaction
  - Everything except hard process final state
- Separation into components is model dependent
  - Operational definition (R Field): “transverse” regions



# Tuning PYTHIA to the Underlying Event

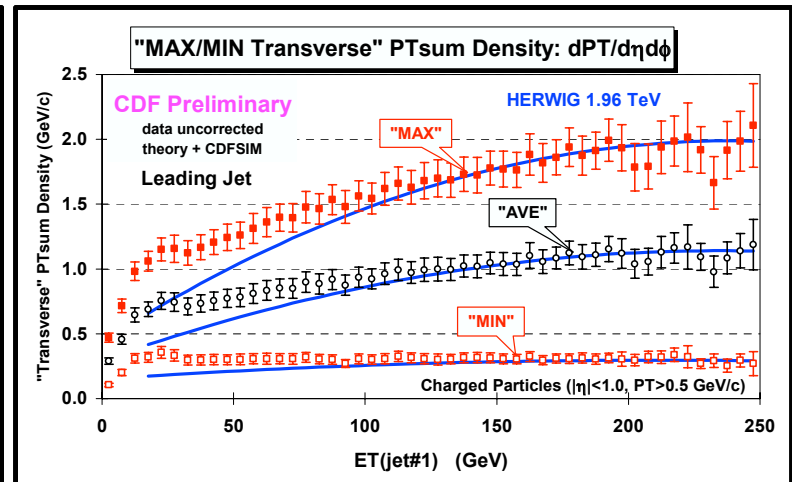
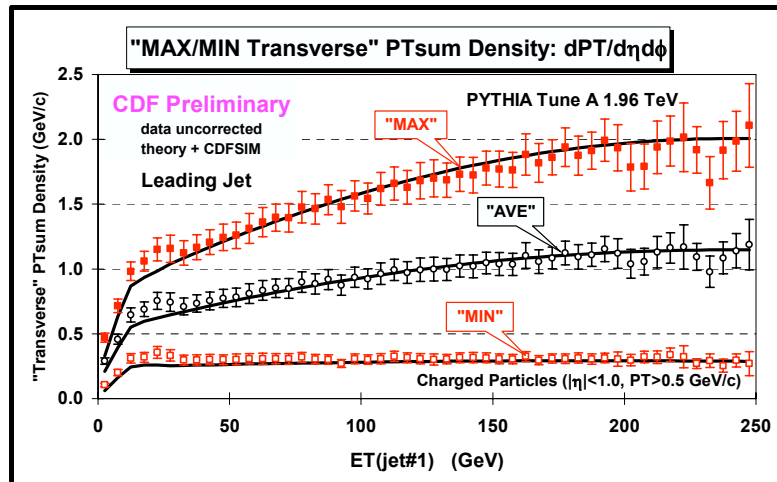
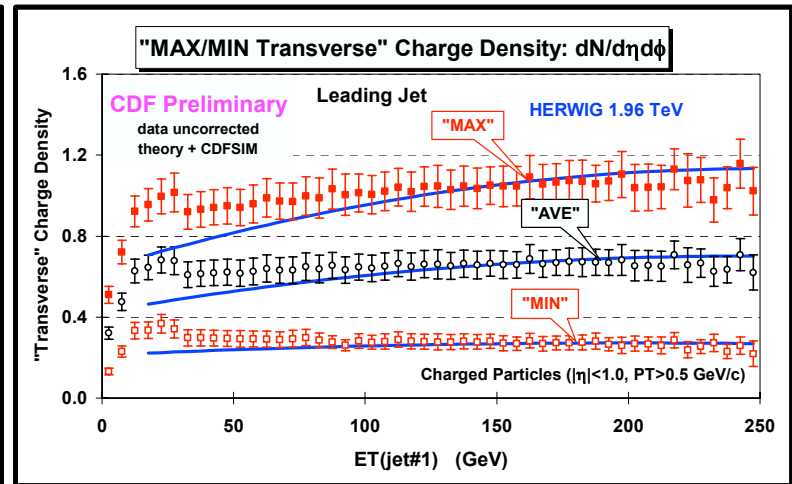
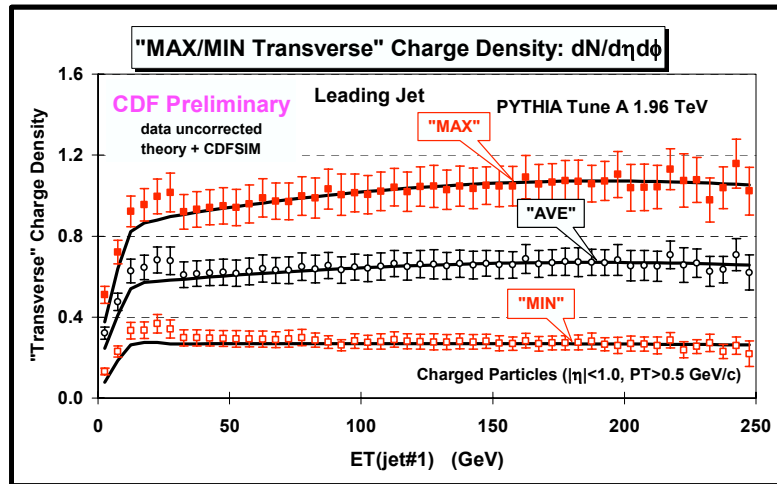
- Rick Field (CDF): keep all parameters that can be fixed by LEP or HERA at their default values. What's left?
- Underlying event. Big uncertainties at LHC...



# Leading Jet: "MAX & MIN Transverse" Densities

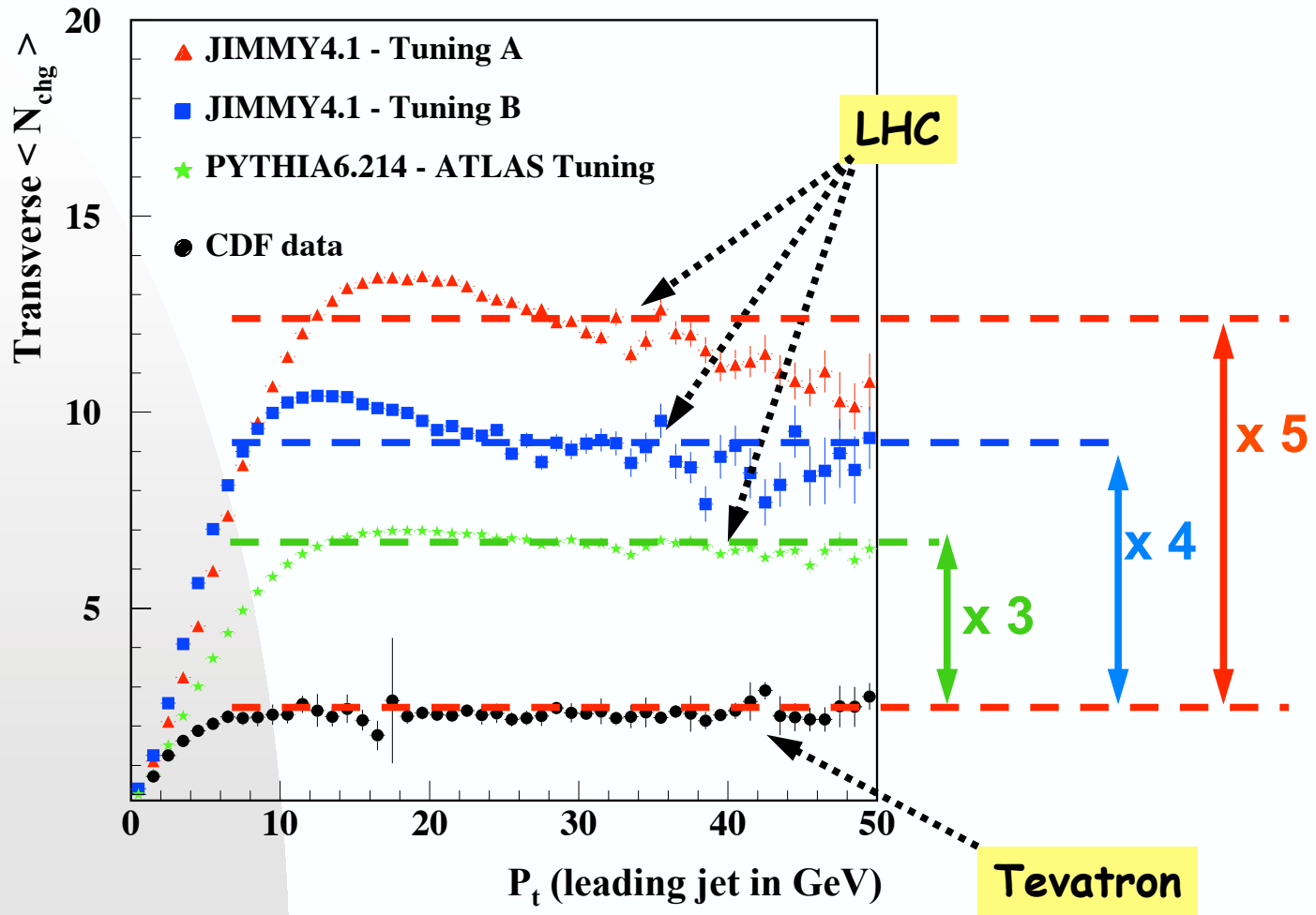
## PYTHIA Tune A

## HERWIG



Charged particle density and PTsum density for "leading jet" events versus  $E_T(\text{jet}\#1)$  for PYTHIA Tune A and HERWIG.

# LHC predictions: JIMMY4.1 Tunings A and B vs. PYTHIA6.214 – ATLAS Tuning (DC2)



A. M. Moraes

Minimum-bias and the Underlying Event at the LHC

5<sup>th</sup> November 2004

# Summary

- Hard Process is very well understood: firm perturbative basis
- Parton Shower is fairly well understood: perturbative basis, with various approximations
- Hadronization is less well understood: modelled, but well constrained by data. Extrapolation to LHC fairly reliable.
- Underlying event least understood: modelled and only weakly constrained by existing data. Extrapolation?
- Always ask “What physics is dominating my effect?”